## Liouville Quantum Gravity

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## What is Liouville Quantum Gravity?

First introduced by Polyakov in 1981:

"Quantum Geometry of Bosonic Strings"

#### Brownian motion:

- Canonical random path
- Scaling limit of random walks

#### Liouville Quantum Gravity:

- Canonical two-dimensional geometry
- Conjectured limit of planar maps

## Conformal Field Theory

Goal of field theory: compute correlations of certain observables called fields:  $\langle \prod_{i \in I} \phi_i(z_i) \rangle$ 

Conformal Field Theory: conformal invariance in 2D ⇒ Imposes constraints on correlation functions

#### Examples of CFT's:

- Continuum limit of critical Ising model
- Liouville Quantum Gravity

Another point of view: SLE curves (see Miller/Sheffield course)

# Brownian motion seen as a path integral

Space of paths:  $\Sigma = \{ \sigma : [0,1] \to \mathbb{R}, \sigma(0) = 0 \}$ 

Action functional:  $S_{BM}(\sigma) = \frac{1}{2} \int_0^1 |\sigma'(r)|^2 dr$ 

$$\mathbb{E}[F((B_s)_{0 \leq s \leq 1})] = \frac{1}{Z} \int_{\Sigma} D\sigma F(\sigma) e^{-S_{BM}(\sigma)}$$

 $D\sigma$ : formal uniform measure on  $\Sigma$ 

#### Classical Theory / Quantum Theory

Minimum of  $S_{BM} \to \text{straight line} = \text{classical solution}$ Path integral  $\to$  Brownian motion = quantum correction

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#### Some definitions

Let M be a two-dimensional surface (sphere, torus,..).

Metric tensor g :  $M o S_2^+(\mathbb{R})$ 

Simple case: 
$$g(x) = \begin{pmatrix} e^{f(x)} & 0 \\ 0 & e^{f(x)} \end{pmatrix}$$

- Area of A:  $\int_A e^{f(x)} dx^2 = \int_A \lambda_g(dx)$
- Gradient squared:  $|\partial^g X|^2 = e^{-f} |\partial X|^2$
- Scalar curvature  $R_g = -e^{-f}\Delta f$ .

Spherical metric on 
$$\mathbb{R}^2$$
:  $g(x) = \frac{4}{(1+|x|^2)^2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$ ,  $R_g = 2$ .

## Classical Liouville Theory

For all maps  $X : M \to \mathbb{R}$ , we define:

$$S_L(X,g) = \frac{1}{4\pi} \int_M (|\partial^g X|^2 + QR_g X + 4\pi \mu e^{\gamma X}) \lambda_g$$

 $Q, \gamma, \mu > 0$  positive constants

#### Uniformization of (M, g)

Assume  $X_{min}$  to be the minimum of  $S_L$  and define  $g'=e^{\gamma X_{min}}g$ . Then  $R_{g'}=-2\pi\mu\gamma^2$  if we choose  $Q=\frac{2}{\gamma}$ .  $\Longrightarrow$  The minimum of  $S_L$  provides a metric of constant negative curvature.

## Defining Liouville Quantum Gravity

#### Formal definition

Random metric  $e^{\gamma\phi}g$  where the law of  $\phi$  is given by:

$$\mathbb{E}[F(\phi)] = \frac{1}{7} \int F(X) e^{-S_L(X,g)} DX$$

First goal: give a meaning to  $\phi$  for different M.

- M = Riemann sphere: David-Kupiainen-Rhodes-Vargas
- ullet M = Torus or higher genus: David-Guillarmou-Rhodes-Vargas
- $\bullet$  M =Unit disk: Huang-Rhodes-Vargas
- *M* = Annulus: Remy

 $\phi = \text{Liouville field}$ 



## Why the Liouville action?

- $|\partial^{g}X|^{2}$ : analogue of the  $|\sigma'|^{2}$  for Brownian motion  $\frac{1}{Z}\int F(X)e^{-\frac{1}{4\pi}\int_{M}|\partial^{g}X|^{2}\lambda_{g}}DX$ : formally defines the law of the Gaussian Free Field (GFF)
- $QR_gX$ : curvature term
- $\int_M e^{\gamma X} \lambda_g$  = area of M in the metric  $g' = e^{\gamma X} g$   $\Rightarrow$  penalizes large areas  $\Rightarrow$  required to have a well defined Liouville field

#### Insertion points

- ullet For  $M=\mathbb{S}^2$ , Gauss-Bonnet:  $\int_{\mathbb{S}^2} R_g \lambda_g = 8\pi > 0$
- No metric of constant negative curvature  $\Rightarrow S_L$  has no minimum  $\Rightarrow \frac{1}{7} \int F(X)e^{-S_L(X,g)}DX$  not defined
- Instead we consider:

$$\frac{1}{Z}\int F(X)e^{\sum_{i=1}^{n}\alpha_{i}X(z_{i})}e^{-S_{L}(X,g)}DX=\langle \prod_{i=1}^{n}e^{\alpha_{i}X(z_{i})}\rangle$$

- = correlation function of the fields  $e^{\alpha_i \phi(z_i)}$
- $(z_i, \alpha_i)$ : insertion points = singularities of the metric
- For S<sup>2</sup>: at least 3 insertions required



## Computing the partition function

Consider  $M = \mathbb{S}^2$ 

Main objective: give a mathematical meaning to

$$\Pi_{\mu,\gamma}^{(z_i,\alpha_i)}(F,g) = \int DX \prod_i e^{\alpha_i X(z_i)} e^{-S_L(X,g)} F(X)$$

where again:

$$S_L(X,g) = \frac{1}{4\pi} \int_M (|\partial^g X|^2 + QR_g X + 4\pi \mu e^{\gamma X}) \lambda_g$$

Remark:  $Q = \frac{2}{\gamma} + \frac{\gamma}{2}$ 

## Step 1: the squared gradient term

Goal: give meaning to  $\frac{1}{Z}\int \widetilde{F}(X)e^{-\frac{1}{4\pi}\int_{M}|\partial^{g}X|^{2}\lambda_{g}}DX$ 

$$e^{-\frac{1}{4\pi}\int_M |\partial^g X|^2 d\lambda_g} = e^{-\frac{1}{2}\int_M X(-\frac{\Delta_g}{2\pi})Xd\lambda_g}$$

Density of an infinite dimensional Gaussian vector of covariance function  $(-\frac{\Delta}{2\pi})^{-1}=$  Green function  $\Rightarrow$  defines a GFF

Gradient term: defines X up to a constant c c = average value of the field We integrate over <math>c with the Lebesgue measure.

## Step 1: the squared gradient term

$$rac{1}{Z}\int\widetilde{F}(X)e^{-rac{1}{4\pi}\int_{M}|\partial^{g}X|^{2}\lambda_{g}}DX=\int_{\mathbb{R}}dc\mathbb{E}[\widetilde{F}(X+c)]$$

- On the l.h.s: formal functional integral
- On the r.h.s: X has the law of a GFF

The partition function  $\Pi_{\mu,\gamma}^{(z_i,\alpha_i)}(F,g)$  becomes:

$$\int_{\mathbb{R}} dc \mathbb{E}[F(X+c) \prod_{i} e^{\alpha_{i}(X(z_{i})+c)} e^{-\frac{1}{4\pi} \int_{M} (QR_{g}(X+c)+4\pi \mu e^{\gamma(X+c)}) \lambda_{g}}]$$

## Step 2 : Gaussian multiplicative chaos

X is a random distribution,  $e^{\gamma X}$  is ill defined

Regularization procedure: circle average  $X_{\epsilon}$ 

$$X_{\epsilon}(z) = \frac{1}{2\pi} \int_{\theta=0}^{2\pi} X(z + \epsilon e^{i\theta}) d\theta$$

#### Gaussian multiplicative chaos

The following limit exists in probability in the sense of weak convergence of measures for  $\gamma \in [0, 2)$ :

$$\lim_{\epsilon \to 0} e^{\gamma X_{\epsilon}(z) - \frac{\gamma^2}{2} \mathbb{E}[X_{\epsilon}(z)^2]} d\lambda_g(z) = e^{\gamma X(z) - \frac{\gamma^2}{2} \mathbb{E}[X(z)^2]} d\lambda_g(z)$$

Remark:  $\mathbb{E}[X_{\epsilon}(z)^2] + \ln \epsilon$  remains bounded as  $\epsilon \to 0$ .

# Step 3: regularization of the partition function

Define 
$$\Pi_{\mu,\gamma}^{(z_i,\alpha_i)}(F,g,\epsilon) = \int_{\mathbb{R}} dc \mathbb{E}[F(X_{\epsilon}+c)\prod_i \epsilon^{\frac{\alpha_i^2}{2}} e^{\alpha_i(X_{\epsilon}(z_i)+c)} e^{-\frac{1}{4\pi}\int_M (QR_g(X_{\epsilon}+c)+4\pi\mu\epsilon^{\frac{\gamma^2}{2}} e^{\gamma(X_{\epsilon}+c)})\lambda_g}]$$

When does the limit  $\lim_{\epsilon \to 0} \Pi_{\mu,\gamma}^{(z_i,\alpha_i)}(F,g,\epsilon)$  exist?

#### Main result

Consider  $M = \mathbb{S}^2$ 

#### Non-triviality of the partition function

Assume  $\gamma \in [0,2)$  and  $\mu > 0$ , then  $\Pi_{\mu,\gamma}^{(z_i,\alpha_i)}(F,g) = \lim_{\epsilon \to 0} \Pi_{\mu,\gamma}^{(z_i,\alpha_i)}(F,g,\epsilon)$  exists and is finite and non zero

$$\iff \sum_i \alpha_i > 2Q$$
 and  $\forall i, \ \alpha_i < Q$ 

Definition of the law of the Liouville field  $\phi$ :

$$\mathbb{E}[F(\phi)] = rac{\Pi_{\mu,\gamma}^{(z_i,lpha_i)}(F,g)}{\Pi_{\mu,\gamma}^{(z_i,lpha_i)}(1,g)}$$

#### Liouville measure

- $\phi$  is a random distribution  $\Rightarrow$  difficult to define  $e^{\gamma\phi}$
- Well defined Liouville measure  $Z(A)=\int_A e^{\gamma\phi}\lambda_g$  Conjectured limit of uniform planar maps for  $\gamma=\sqrt{\frac{8}{3}}$

Conjectured limit of planar maps with an Ising model for  $\gamma = \sqrt{3}$ 

## Surfaces with boundary

Must add boundary terms to  $S_L$ :

$$\tilde{S}_L(X,g) = S_L(X,g) + \frac{1}{2\pi} \int_{\partial M} (QK_gX + 2\pi\mu_\partial e^{\frac{\gamma}{2}X}) \lambda_{\partial g}$$

Example: M = unit disk

- Bulk insertion points  $(z_i, \alpha_i)$
- Boundary insertion points  $(s_j, \beta_j)$

#### Non-triviality of the partition function

Assume  $\gamma \in [0,2)$ ,  $\mu_{\partial} > 0$ , and  $\mu > 0$ , then

$$\Pi_{\mu,\gamma}^{(z_i,\alpha_i)}(F,g) = \lim_{\epsilon \to 0} \Pi_{\mu,\gamma}^{(z_i,\alpha_i)}(F,g,\epsilon)$$
 exists and is finite and non zero

$$\iff \sum_i \alpha_i + \sum_j rac{eta_j}{2} > Q$$
,  $\forall i, \ lpha_i < Q$  and  $\forall j, \ eta_j < Q$ 

# Surfaces of higher genus

Higher genus: non trivial moduli space

- Torus
- Annulus

$$A(a, b) =$$
 annulus in the plane with radii  $a < b$ 

Then 
$$A(a,b) \sim A(a',b') \Longleftrightarrow \frac{a}{b} = \frac{a'}{b'}$$

 $\tau = \frac{a}{b} = \text{modular parameter, important in physics.}$ 

## Thank you for listening!

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