

# Pauli Operators in Quantum Computing

Pauli operators form the backbone of quantum operations on single qubits. These operators, denoted as  $X$ ,  $Y$ , and  $Z$ , correspond to the physical concepts of bit flips, phase flips, and combined bit and phase flips, respectively. They are represented by the matrices:

- **X (Bit-flip operator):**  $\sigma_x = X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$ . This operator flips the state of a qubit from 0 to 1 and vice versa, analogous to the classical NOT gate.
- **Y (Bit-phase-flip operator):**  $\sigma_y = Y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$ . The Y operator introduces both a bit flip and a phase flip, thus affecting both the computational basis state and its phase.
- **Z (Phase-flip operator):**  $\sigma_z = Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$ . The Z operator changes the phase of the 1 state without altering the bit value.

These operators are unitary, ensuring reversibility of quantum operations, and Hermitian, guaranteeing real eigenvalues which correspond to measurable quantities. They obey the anticommutation relation,  $XY = -YX$ ,  $XZ = -ZX$ , and  $YZ = -ZY$ , highlighting the non-commutative nature of quantum mechanics. Additionally, squaring any of these operators yields the identity matrix, reflecting their involutory nature.

These operators are commonly referred to as sigma operators, particularly when denoted by  $\sigma_x$ ,  $\sigma_y$ , and  $\sigma_z$ , or as Pauli spin matrices when expressed as matrices in the standard basis, as presented here. Given their widespread prevalence in quantum physics, it is essential to commit them to memory.

## From Bit-flips to Phase-flips, and Back Again

The Hadamard gate ( $H$ ) plays a pivotal role in quantum computing by converting between

$$H = \frac{(X + Z)}{\sqrt{2}} \rightarrow H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

bit-flips and phase-flips. Represented as

superpositions when applied to basis states. The ability of the Hadamard gate to transform  $Z$  operations into  $X$  operations (and vice versa) is fundamental for algorithms that rely on quantum interference and entanglement.

## Any Unitary Operation on a Single Qubit

Every single-qubit operation can be represented as a unitary matrix, which, due to the constraints of unitarity, can be decomposed into rotations around the Bloch sphere. This geometric interpretation allows for any unitary operation to be understood as a combination of rotations, which can be parameterized by three angles, offering a comprehensive way to visualize and implement quantum operations.

A unitary matrix of size  $n \times n$  has  $2n^2$  real parameters initially. The unitarity condition reduces this to  $n^2$  parameters because specifying  $n^2$  of them determines the rest through the unitarity equation.

A complex number  $z$  is defined by two real parameters, either as  $z = x + iy$  or  $z = re^{i\phi}$ , indicating that  $\mathbb{C}$  is a two-dimensional vector space over  $\mathbb{R}$ .

For a  $(2 \times 2)$  unitary matrix, four real parameters are needed, but ignoring global phase, three suffice. It's possible to construct any single qubit unitary operation using a circuit with two Hadamard gates and three phase gates.

The universal circuit for a  $(2 \times 2)$  unitary matrix is determined by three real parameters, up to a global phase. The matrix from this circuit is:

$$U(\alpha, \beta, \phi) = \begin{bmatrix} e^{-i(\frac{\alpha+\beta}{2})} \cos \frac{\phi}{2} & -ie^{i(\frac{\alpha-\beta}{2})} \sin \frac{\phi}{2} \\ -ie^{-i(\frac{\alpha-\beta}{2})} \sin \frac{\phi}{2} & e^{i(\frac{\alpha+\beta}{2})} \cos \frac{\phi}{2} \end{bmatrix}.$$

This shows a  $(2 \times 2)$  unitary matrix can be expressed with three parameters  $\alpha$ ,  $\beta$ , and  $\phi$ , within  $[0, 2\pi]$ , linking single-qubit unitaries to three-dimensional rotations.

## The Bloch Sphere

The Bloch sphere is a spherical representation of qubit state space, where every point on the sphere corresponds to a possible state of a qubit. The north and south poles represent the basis states  $|0\rangle$  and  $|1\rangle$ , respectively, while points on the surface represent superpositions. This model is instrumental in visualizing the effects of quantum operations as rotations on the sphere, providing intuitive insights into the behavior of qubits under various transformations.

Unitary operations on a single qubit form a non-abelian group under matrix multiplication, called  $U(2)$ . This group's behavior is analogous to rotations in three-dimensional space, leading to the relation  $U(2)/U(1) \cong SO(3)$ . Essentially,  $2 \times 2$  unitaries, minus their global phase, are isomorphic to  $SO(3)$ , the group of 3D rotations. This isomorphism aids in visualizing single-qubit gate effects.  $U(1)$  corresponds to  $\mathbb{C}^\times$ , the multiplicative group of non-zero complex numbers. The action of  $U(2)$  on single-qubit states is paralleled by  $SO(3)$  acting on the unit sphere  $S^2 \subset \mathbb{R}^3$ . A single-qubit state  $|\psi\rangle$ , normalized such that  $|\alpha|^2 + |\beta|^2 = 1$ , can be parameterized as:  $|\psi\rangle = \cos(\theta/2)e^{i\phi_0}|0\rangle + \sin(\theta/2)e^{i\phi_1}|1\rangle$ , with global phase ignored, as it doesn't affect the physical state.

The Bloch sphere visualizes quantum states with the Bloch vector, defined by  $\theta$  and  $\phi$ , indicating the state's position. Unitary operations rotate this vector, with orthogonal vectors on the sphere representing orthogonal quantum states.

Rotations by unitaries like the phase gate  $P_\alpha$ , and the Pauli operators  $Z$ ,  $X$ , and  $Y$ , correspond to specific axes rotations on the Bloch sphere. The Hadamard gate's effect, swapping the  $x$  and  $z$  axes, represents a  $180^\circ$  rotation around the  $(x + z)$ -axis.

The transformation of a Bloch vector by a unitary  $U$  is given by:

$$U(s_x X + s_y Y + s_z Z)U^\dagger = s'_x X + s'_y Y + s'_z Z,$$

showcasing how quantum operations are equivalent to geometrical rotations on the Bloch sphere.

## Drawing Points on the Bloch Sphere

The position of a qubit state on the Bloch sphere can be determined by its state vector,  $\psi = \alpha|0\rangle + \beta|1\rangle$ . The angles  $\theta$  and  $\varphi$  in spherical coordinates relate to the probability amplitudes  $\alpha$  and  $\beta$ , enabling the mapping of quantum states to points on the sphere. This visual representation aids in understanding the geometric interpretation of qubit states and the impact of quantum gates as rotations around the sphere.

## Composition of Rotations

The ability to decompose any rotation on the Bloch sphere into a sequence of rotations around the axes simplifies the implementation of unitary operations. This decomposition is analogous to Euler angles in classical mechanics, where any orientation can be achieved through three consecutive rotations about principal axes. In quantum computing, this principle allows for the construction of any single-qubit gate using a finite set of universal gates, reducing the complexity of quantum algorithms.

