

Lecture 6 Hodge theorem & index theory.

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Let's start by finishing the proof of:

Hodge theorem (M, g) compact Riemannian. Then

$$\Omega^p = \mathcal{H}^p \oplus \text{Im } \Delta \quad (L^2\text{-orthogonal decomposition}).$$

$$\text{harmonic } p\text{-forms} = \text{ker } \Delta$$

$$\Rightarrow \Omega^p = d\Omega^{p-1} \oplus \mathcal{H}^p \oplus d^*\Omega^{p+1}.$$

Prop $d + d^* : \Omega^p(M) \rightarrow \Omega^p(M)$ is a first order elliptic

op ($\Rightarrow \Delta = (d + d^*)^2$ is second order elliptic).

Proof intrinsic def: $x \in M$, $\xi = d\varphi(x) \in T_x^*(M \setminus \{x\})$

with $\varphi(x) = 0$, $\alpha \in \Omega^p$. Want to show

$$(d + d^*)(\varphi\alpha)(x) \neq 0.$$

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To see this, compute:

$$d(\varphi \alpha)(x) = (d\varphi \wedge \alpha + \varphi d\alpha)(x) = \xi \wedge \alpha.$$

$$\rightarrow d^*(\varphi \alpha)(x) = *d(\varphi \cdot * \alpha)(x) = \dots = *(\xi \wedge * \alpha)$$

$$\Rightarrow d^*(\varphi \alpha)(x) = \xi \lrcorner \alpha \text{ (where } \xi \lrcorner \alpha \text{ is}$$

is $e \rightarrow \omega := (-1)^{n(p+1)} * (e \wedge * \omega)$, used contraction with e^i up to sign).

$$\Rightarrow \text{symbol is } \xi \wedge \alpha + \xi \lrcorner \alpha \neq 0.$$

Proof of Hodge thm write $\Omega^p = \mathcal{H}^p \oplus (\mathcal{H}^p)^\perp$,

want to prove $(\mathcal{H}^p)^\perp = \Delta(\Omega^p)$.

\Leftarrow is clear: if $\alpha = \Delta \beta$, $\omega \in \mathcal{H}^p$

$$\Rightarrow \langle \alpha, \omega \rangle_{L^2} = \langle \Delta \beta, \omega \rangle_{L^2} = \langle \beta, \Delta \omega \rangle_{L^2} = 0.$$

non trivial part is \subseteq !

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Rule we don't know yet $\mathcal{H}^p \cong H^p(\Omega; \mathbb{R})$!

If $\alpha \in (\mathcal{H}^p)^\perp$, want to show $\exists \beta \in \Omega^p$ s.t.

$\alpha = \Delta \beta$. In fact, by elliptic regularity,

we only need to find a weak solution, i.e.

$\beta \in L^2(\Omega^p)$ s.t.

$$\langle \alpha, \varphi \rangle_{L^2} = \langle \beta, \Delta \varphi \rangle_{L^2} \quad \forall \varphi \in \Omega^p.$$

But β exists by abstract nonsense as follows.

Define $l: \text{Im} \Delta \xrightarrow{\subseteq \Omega^p} \mathbb{R}$

$$\Delta \varphi \mapsto \langle \alpha, \varphi \rangle_{L^2}$$

•) well-defined: if $\Delta \varphi = \Delta \varphi' \rightarrow \varphi - \varphi' \in \mathcal{H}^p$

$$\langle \alpha, \varphi \rangle_{L^2} = \langle \alpha, \varphi' \rangle_{L^2}.$$

•) bounded, see above if

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$$\varphi = \varphi^{\text{harm}} + \varphi^\perp \in \mathcal{H}^P \oplus (\mathcal{H}^P)^\perp,$$

$$|\ell(\Delta\varphi)| = |\ell(\Delta\varphi^\perp)| = |\langle \alpha, \varphi^\perp \rangle_{L^2}|$$

$$\leq \|\alpha\|_{L^2} \|\varphi^\perp\|_{L^2} \quad (\text{by Cauchy-Schwarz})$$

$$\leq \|\alpha\|_{L^2} \|\varphi^\perp\|_{L^2}^2$$

$$\leq C \|\alpha\|_{L^2} \|\Delta\varphi^\perp\|_{L^2} \quad (\text{Gårding inequality}).$$

$$= C \|\alpha\|_{L^2} \|\Delta\varphi\|_{L^2}.$$

By Hahn-Banach, ℓ extends to bounded

$$\ell: L^2(\Omega^P) \rightarrow \mathbb{R}. \quad (\text{i.e. } \ell \in L^2(\Omega^P)^*).$$

\hookrightarrow Hilbert space

By Riesz representation, $\ell(-) = \langle \beta, - \rangle_{L^2}$

for $\beta \in L^2(\Omega^P)$, which is weak solution

$$\text{of } \Delta\beta = \alpha$$

□.

Elliptic operators have very interesting
topological properties!

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Basic observation $T: \mathbb{R}^n \rightarrow \mathbb{R}^m$ linear map.

By rank-nullity,
 $\text{ind}(T) := \dim \ker T - \dim \text{coker } T = n - m$
depends on T . ($= \mathbb{R}^m / \text{Im } T$)

independent
of T

Def X, Y Hilbert spaces (or Banach),

$T: X \rightarrow Y$ bounded op is Fredholm if

- a) $\dim \ker T < \infty$
- b) $\dim \text{coker } T < \infty$
- c) $\text{Im } T$ closed.

The index of T is

$$\text{ind}(T) := \dim \ker T - \dim \text{coker } T \in \mathbb{Z}.$$

Rank in this setup, $2) \Rightarrow 3)$ by
open mapping thm.

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Ex $\rightarrow R: \ell^2(\mathbb{N}) \xrightarrow{\circlearrowright}$ right shift

$$(a_0, a_1, \dots) \mapsto (0, a_0, a_1, \dots)$$

$$\text{has } \text{ind } R = 0 - 1 = -1.$$

$\cdot) L: \ell^2(\mathbb{N}) \xrightarrow{\circlearrowleft}$ left shift

$$(a_0, a_1, \dots) \mapsto (a_1, \dots)$$

$$\text{has } \text{ind } L = 1 - 0 = 1$$

Thm T order l elliptic op on M cpt, ker

$$\forall k \quad T: L^2_{k+l} \rightarrow L^2_k \text{ is Fredholm,}$$

and $\text{ind}(T)$ is independent of k

(inject, $\text{coker } T \cong \text{ker } T^*$
 \hookrightarrow formally adjoint
 diff op)

$\xrightarrow{\text{hot word!}}$
Thm $\text{Fred}(X, Y) \subseteq \mathcal{B}(X, Y)$ is open, ⑧
 \uparrow \uparrow
 Fredholm bounded

and $\text{ind}: \text{Fred}(X, Y) \rightarrow \mathbb{Z}$ is locally constant.

Cor $T_t: \mathcal{P}(E) \rightarrow \mathcal{P}(F)$ 1-parameter family of elliptic ops on M cpt $\Rightarrow \text{ind}(T_t) \in \mathbb{Z}$ is constant.

Ex if T is self adjoint (e.g. $d+d^*: \Omega^k \rightarrow \Omega^k$)
 $\text{coker } T = \ker T^* = \ker T \Rightarrow \text{ind } T = 0$.

More interestingly, consider

$$D = d + d^*: \Omega^{\text{even}} \rightarrow \Omega^{\text{odd}} \quad (\text{elliptic!})$$

$$\Rightarrow \ker = \bigoplus \mathcal{H}^{\text{even}} = \bigoplus H^{\text{even}}(M; \mathbb{R})$$

$$\text{coker} = \ker(d + d^*: \Omega^{\text{odd}} \rightarrow \Omega^{\text{even}})$$

$$= \bigoplus H^{\text{odd}}(M; \mathbb{R})$$

$$\Rightarrow \text{ind}(D) = \chi(M)$$

Now, choose $\alpha \in \mathcal{R}^{\text{odd}}(M)$. Then

$$D + t(\alpha \cdot -) : \mathcal{R}^{\text{even}} \rightarrow \mathcal{R}^{\text{odd}} \text{ is}$$

1-parameter family of elliptic ops

$$\Rightarrow \text{ind}(D + t(\alpha \cdot -)) = \chi(M)!$$

So if $\chi(M) > 0$ (e.g. $M = S^{2n}$)

$$\ker(D + t(\alpha \cdot -)) \neq \emptyset$$

\leadsto topology forces stability of elliptic PDEs.

Atiyah-Singer index thm (1963). For

$T: \Gamma(E) \rightarrow \Gamma(F)$ elliptic on M cpt,

purely topological (e.g. char classes of π^1, E, F)

explicit formula for $\text{ind}(T) \in \mathbb{Z}$.

Back to gauge theory

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The (linear part) of the equations we study.

is never elliptic because of gauge invariance.

Consider for example

$$F_A^+ = 0 \in \Omega^2(\mathfrak{g}_E) \quad \text{ASD (or instanton) eqn.}$$

Locally, if $\nabla_A^z = d + A^z$ $A^z \in \Omega^1(\mathfrak{g})$

$$F_A^+ = (dA^z + A^z \wedge A^z)^+$$

So the highest order term is

$$\Omega^1(\mathfrak{g}) \xrightarrow{d} \Omega^2(\mathfrak{g}) \xrightarrow{\pi^+} \Omega^+(\mathfrak{g})$$

$\xrightarrow{d^+}$

and of course $d^+(\underbrace{df}_{\text{on any line used}}) = 0 \Rightarrow$ not elliptic.

(comes from $F_{v \cdot A}^+ = 0 \quad \forall v \in \mathfrak{g}$).

This issue is resolved by suitably

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Gauge fixing. Let's do the U(1) case

(Weyl-abelian case highly non-trivial,

cf Uhlenbeck's gauge fixing thm (1982))

$L \rightarrow (M, g)$ cpt, fix base connection A_0 .

$\mathfrak{g}_0 = \{v = e^{\sharp}, f \in i\mathbb{R}^{\circ}\}$ id. component of \mathfrak{g} .

$$\text{Then } e^{\sharp} \cdot B = A - dg$$

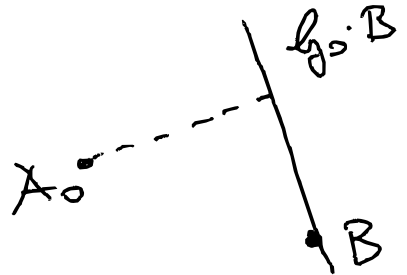
It's natural to consider

\tilde{B} that minimizes $\|v \cdot B - A_0\|_{L^2}$.

This solves (usual computation)

$$\boxed{d^*(B - A_0) = 0}$$

Coulomb gauge
condition



Hodge Lemma $\Rightarrow \forall B, \exists!$ f

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up to overall $\text{const} \rightarrow \mathbb{R}^+$ s.t. $e^f \cdot B$ is
is Coulomb gauge.

Now, the first order part of ASD eqn

+ Coulomb gauge fixing is

$$d^* + d^{\dagger} : i\Omega^1(M) \rightarrow i\Omega^0(M) \oplus i\Omega^2(M)$$

which is elliptic! (dimensional check: $4 = 1 + 3$)

One also checks $\ker = \mathcal{H}^1$, and

$$\text{coker} = \underbrace{\mathbb{R}}_{\text{constants}} \oplus \mathcal{H}^2, \text{ self-dual harmonic } 2\text{-forms},$$

(+1) eigenspace of $\star : \mathcal{H}^2 \xrightarrow{\cong}$, where \dim is the

$b^{\dagger}(x) = \text{positivity of quadratic form}$

$$Q_x : H^2(x; \mathbb{R}) \rightarrow \mathbb{R} \quad \alpha \mapsto (\alpha \cup \alpha)[X].$$

$\Rightarrow \text{ind} = b_1 - b^{\dagger} - 1$. (top quantity!).