

# Large systems of diffusions interacting through their ranks

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# Outline

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# Setting

- Fix a natural number  $N \in \mathbb{N}$ , real numbers  $b_1, b_2, \dots, b_N$  and positive real number  $\sigma_1, \sigma_2, \dots, \sigma_N$ .
- Consider a system of **interacting diffusions (particles)** on  $\mathbb{R}$ :

$$dX_i(t) = \sum_{j=1}^N b_j \mathbf{1}_{\{X_i(t)=X_{(j)}(t)\}} dt + \sum_{j=1}^N \sigma_j \mathbf{1}_{\{X_i(t)=X_{(j)}(t)\}} dW_i(t),$$

where  $W_1, W_2, \dots, W_N$  are i.i.d. standard Brownian motions and some initial values  $X_1(0), X_2(0), \dots, X_N(0)$  are fixed.

## Some historic remarks

- Model appeared '87 in this form in paper by **Bass** and **Pardoux** in the context of filtering theory. They proved existence and uniqueness in law.
- Model reappeared in stochastic portfolio theory ('02 book by **Fernholz**, '06 survey by **Fernholz** and **Karatzas**): diffusions  $X_1, X_2, \dots, X_N$  represent **logarithmic capitalizations**:

$$\log(\text{stock price} \times \text{number of stocks}). \quad (1)$$

Model assumes that dynamics depends only on **ranks**. **True in the long run**: explains following picture.

# Capital distribution curves

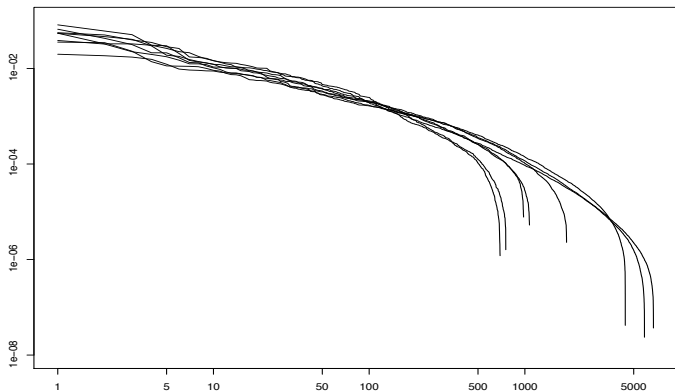


Figure: Capital distribution curves 1929-1989

## Concentration results

- **Chatterjee and Pal '07**: for particle system above, vector of **market weights**

$$\left( \frac{e^{X_i(t)}}{\sum_{j=1}^N e^{X_j(t)}}, i = 1, 2, \dots, N \right) \quad (2)$$

is a Markov process and its **invariant distribution concentrates** around curves of above type as  $N \rightarrow \infty$ . Moreover, the limit  $N \rightarrow \infty$  is given by a **Poisson-Dirichlet point process** of first kind.

- **Pal, S. '10** and **Ichiba, Pal, S. '11**: strong concentration of paths of market weights on any  $[0, t]$  as  $N \rightarrow \infty$  and fast mean-reversion as  $t \rightarrow \infty$  for any fixed  $N \in \mathbb{N}$ .

## Back to particle system

All previous results for **market weights**, which correspond to the **spacings process** in

$$dX_i(t) = \sum_{j=1}^N b_j \mathbf{1}_{\{X_i(t)=X_{(j)}(t)\}} dt + \sum_{j=1}^N \sigma_j \mathbf{1}_{\{X_i(t)=X_{(j)}(t)\}} dW_i(t),$$

What about the **particle system** itself? Can we understand **evolution of particle density**:

$$\varrho^N(t) = \frac{1}{N} \sum_{i=1}^N \delta_{X_i(t)}, \quad t \geq 0.$$

# Preliminaries

**Here:** will look at limit  $N \rightarrow \infty$ , which corresponds to a **hydrodynamic limit**.

**First question:** how to choose drift and diffusion coefficients for different  $N$  to have a meaningful limit?

**Crucial observation:** for **fixed**  $N$  the particle system can be written as

$$dX_i(t) = b(F_{\varrho^{N(t)}}(X_i(t))) dt + \sigma(F_{\varrho^{N(t)}}(X_i(t))) dW_i(t)$$

for **functions**  $b : [0, 1] \rightarrow \mathbb{R}$ ,  $\sigma : [0, 1] \rightarrow (0, \infty)$ .

$\Rightarrow$  particle system is of **mean-field** type.



## Aside on mean-field models

Systems of the form

$$dX_i(t) = \hat{b}(\varrho^N(t), X_i(t)) dt + \hat{\sigma}(\varrho^N(t), X_i(t)) dW_i(t) \quad (3)$$

appeared in **statistical physics**.

See: **McKean '69, Funaki '84, Oelschläger '84, Nagasawa, Tanaka '85, '87, Sznitman '86, Leonard '86, Dawson, Gärtner '87, Gärtner '88.**

In Gärtner '88, limit  $\lim_{N \rightarrow \infty} \varrho^N(t)$  is obtained under **two assumptions**.

## Previous results I

**Theorem.** (Gärtner '88) Fix  $T > 0$  and suppose  $\hat{b}$ ,  $\hat{\sigma}$  **continuous (!)**,  $\hat{\sigma}$  **strictly positive**.

Let  $Q^N$  be the law of  $\varrho^N(t)$ ,  $t \in [0, T]$  on  $C([0, T], M_1(\mathbb{R}))$ . Then the sequence  $Q^N$ ,  $N \in \mathbb{N}$  is tight. Moreover, under any limit point  $Q^\infty$ :

$$\forall f : (\varrho(t), f) - (\varrho(0), f) = \int_0^t (\varrho(s), L_{\varrho(s)} f) ds \quad (4)$$

for all  $t \in [0, T]$  almost surely. Hereby:

$$\begin{aligned} (\varrho(t), f) &= \int_{\mathbb{R}} f d\varrho(t) \\ L_{\varrho(s)} f &= \hat{b}(\varrho(s), \cdot) f' + \frac{1}{2} \hat{\sigma}(\varrho(s), \cdot)^2 f''. \end{aligned}$$

## Previous results II

In particular, if

$$\forall f : (\varrho(t), f) - (\varrho(0), f) = \int_0^t (\varrho(s), L_{\varrho(s)} f) ds \quad (5)$$

has a **unique** solution  $\varrho^\infty$  in  $C([0, T], M_1(\mathbb{R}))$ , then it must hold

$$\varrho^N \rightarrow \varrho^\infty, \quad N \rightarrow \infty \quad \text{in probability.} \quad (6)$$

This is **not known** in general (some conditions in work of **Sznitman** '86)!

# Diffusions interacting through their mean-field, intuition

How can one **guess** the **limiting dynamics**?

- Suppose we already knew  $\varrho^N \rightarrow \varrho^\infty$  with  $\varrho^\infty$  **deterministic**
- Then for **large**  $N$  the system of diffusions should behave as

$$dX_i(t) = \hat{b}(\varrho^\infty(t), X_i(t)) dt + \hat{\sigma}(\varrho^\infty(t), X_i(t)) dW_i(t) \quad (7)$$

- Thus, the empirical measure converges to the law of

$$dX(t) = \hat{b}(\varrho^\infty(t), X(t)) dt + \hat{\sigma}(\varrho^\infty(t), X(t)) dW(t) \quad (8)$$

- **Ito's formula** for  $f(X(t))$  and  $\mathcal{L}(X(t)) = \varrho^\infty(t)$  imply:

$$(\varrho^\infty(t), f) - (\varrho^\infty(0), f) = \int_0^t (\varrho^\infty(s), L_{\varrho^\infty(s)} f) ds$$

## Previous results III

**Theorem.** (Dawson, Gärtner '87) Fix  $T > 0$  and suppose  $\hat{b}$  **continuous**,  $\hat{\sigma} \equiv 1$  (!)

Then the sequence  $(\varrho^N(t), t \in [0, T])$ ,  $N \in \mathbb{N}$  satisfies a large deviations principle on  $C([0, T], M_1(\mathbb{R}))$  with the good rate function

$$I(\gamma) = \sup_{g \in \bar{S}} \left[ (\gamma(T), g) - (\gamma(0), g) - \int_0^T (\gamma(t), \mathcal{R}_t^\gamma g + \frac{1}{2}(g_x)^2) dt \right]$$

and scale  $N$ . Hereby:

$$\mathcal{R}_t^\gamma g = g_t + \hat{b}(\varrho(s), \cdot) g_x + \frac{1}{2} g_{xx}.$$

## Relation to Burger's equation

### Remarks:

- **Goodness** of rate function + **LDP** imply:  $\varrho^N$  will concentrate around the set  $\{\gamma : I(\gamma) = 0\}$ .
- If we apply this result with  $\hat{b}(\varrho(s), \cdot) = -F_{\varrho(s)}(\cdot)$  (**discontinuity!**), then only point with  $I(\gamma) = 0$  is the one, whose path of cdfs  $R(t, \cdot) = F_{\gamma(t)}(\cdot)$  is the unique weak solution of **viscous Burger's equation**:  $R_t = -\frac{1}{2}(R^2)_x + \frac{1}{2}R_{xx}$ .

I.e.: **particle system approximation** of  $R$ . Same result for a particle system with local time interactions in **Sznitman '86**.

## Our results I

**Theorem 1.** (Dembo, Krylov, S., Varadhan, Zeitouni '12) Fix  $T > 0$  and suppose  $\hat{b}(\varrho(t), x) = b(F_{\varrho(t)}(x))$ ,  $\hat{\sigma} = \sigma(F_{\varrho(t)}(x))$ ;  $b$  and  $A := \frac{1}{2}\sigma^2$  nice. Then,  $(\varrho^N(t), t \in [0, T])$ ,  $N \in \mathbb{N}$  satisfies an LDP on  $C([0, T], M_1(\mathbb{R}))$  with scale  $N$  and good rate function  $J$  defined by

$$J(R) = \frac{1}{2} \left\| \frac{\sigma(R)}{2} \frac{R_t - (A(R)R_x)_x}{A(R)(R_x)^{1/2}} - \frac{b(R)}{\sigma(R)} (R_x)^{1/2} \right\|_{L^2(\mathbb{R}_T)}^2$$

for all  $R \in C_b(\mathbb{R}_T)$  with  $R_t, R_x, R_{xx} \in L^{3/2}(\mathbb{R}_T)$ ,  $R_x \in L^3(\mathbb{R}_T)$ ,  $R_x$  having finite  $(1 + \varepsilon)$  moment,  $t \mapsto (R_x(t, \cdot), g(t, \cdot))$  abs. cont.;  $J = \infty$  otherwise.

Hereby,  $R = F_{\gamma(\cdot)}(\cdot)$ .

## Our results II

### Consequences:

- **Goodness of rate function** and **LDP** imply that  $\varrho^N$  concentrates around the set  $\{\gamma : J(\gamma) = 0\}$ .
- The only path  $\gamma$  with  $J(\gamma) = 0$  corresponds to the unique weak solution of the **generalized porous medium equation with convection**:  $R_t = \Sigma(R)_{xx} + \Theta(R)_x$ .
- Hence, we found a **particle system approximation** for the solution of the latter, which converges exponentially fast.



## Our results III

In the course of the proof show the following **regularity result** in nonlinear PDEs:

**Theorem 2.** Consider a weak solution of the Cauchy problem for the **tilted generalized porous medium equation**:

$$R_t - (A(R)R_x)_x = h A(R) R_x, \quad R(0, \cdot) = R_0.$$

such that  $R \in C_b(\mathbb{R}_T)$  and  $R_x(t, \cdot) dx$  is a probability measure for every  $t$ .

If  $\int_{\mathbb{R}_T} h^2 R_x dm < \infty$  and  $R_x$  has finite  $(1 + \varepsilon)$  moment, then  $R_t, R_x, R_{xx}$  exist as elements of  $L^{3/2}(\mathbb{R}_T)$ ,  $R_x \in L^3(\mathbb{R}_T)$  and

$$\int_{\mathbb{R}_T} \frac{R_{xx}^2}{R_x} dm < \infty, \quad \int_{\mathbb{R}_T} \frac{R_t^2}{R_x} dm < \infty.$$

## Proof sketch, general principles I

**Localization:** LDP holds, if we can show **weak/local LDP**:

$$\limsup_{N \rightarrow \infty} \frac{1}{N} \log \mathbb{P}(\varrho^N \in C) \leq - \inf_{\gamma \in C} J(\gamma) \text{ for all compacts } C,$$

$$\liminf_{N \rightarrow \infty} \frac{1}{N} \log \mathbb{P}(\varrho^N \in U) \geq - \inf_{\gamma \in U} J(\gamma) \text{ for all open sets } U$$

and **exponential tightness**:

$$\forall K > 0 \exists C_K \text{ compact} : \limsup_{N \rightarrow \infty} \frac{1}{N} \log \mathbb{P}(\varrho^N \notin C_K) \leq -K.$$

## Proof sketch, general principles II

Alternative characterization of **weak/local LDP**:

$$\forall \gamma : \lim_{\delta \downarrow 0} \limsup_{N \rightarrow \infty} \frac{1}{N} \log \mathbb{P}(\varrho^N \in B(\gamma, \delta)) \leq -J(\gamma),$$

$$\forall \gamma : \lim_{\delta \downarrow 0} \liminf_{N \rightarrow \infty} \frac{1}{N} \log \mathbb{P}(\varrho^N \in B(\gamma, \delta)) \geq -J(\gamma)$$

**What we prove:**

- Local upper bound holds with a **Dawson-Gärtner type rate function  $I$**
- Local lower bound holds with the **desired rate function  $J$**
- $J \leq I$ .

## Proof sketch, upper bound around a path $\gamma$

Appropriate **variational problem**:

$$I(\gamma) = \sup_{g \in \mathcal{S}} \left[ (\gamma(T), g) - (\gamma(0), g) - \int_0^T (\gamma(t), \mathcal{R}_t^\gamma g + \frac{1}{2} A(R)(g_x)^2) dt \right],$$

where  $\mathcal{R}_t^\gamma g = g_t + b(R)g_x + A(R)g_{xx}$ ,  $R = F_{\gamma(\cdot)}(\cdot)$ .

Why **appropriate**? On the event  $\varrho^N \in B(\gamma, \delta)$ , our particle system is **close** to solution of

$$dY_i(t) = b(F_{\gamma(t)}(Y_i(t))) dt + \sigma(F_{\gamma(t)}(Y_i(t))) dW_i(t), \quad i = 1, 2, \dots, N.$$

on **exponential scale**.

## Proof sketch, upper bound around a path $\gamma$

Pick test function  $g$ , apply **Itô's formula**:

$$\begin{aligned} dg(t, Y_i(t)) &= (g_t + b(R)g_x + A(R)g_{xx})(t, Y_i(t)) dt \\ &\quad + g_x(t, Y_i(t))\sigma(F_{\gamma(t)}(Y_i(t))) dW_i(t). \end{aligned}$$

Hence,

$$\begin{aligned} d(\varrho_Y^N(t), g(t, \cdot)) &= (\varrho_Y^N(t), g_t + b(R)g_x + A(R)g_{xx}) dt \\ &\quad + \frac{1}{N} \sum_{i=1}^N g_x(t, Y_i(t))\sigma(F_{\gamma(t)}(Y_i(t))) dW_i(t). \end{aligned}$$

Note: **martingale part of order  $\frac{1}{N}$** . **Freidlin-Wentzell** type problem!

## Proof sketch, upper bound around a path $\gamma$

Thus, for **fixed**  $g$ , rate is given by

$$I^g(f) = \frac{1}{2} \int_0^T \frac{|\dot{f}(u) - (\gamma(u), g_t + b(R)g_x + A(R)g_{xx})|^2}{(\gamma(u), \sigma(R)^2(g_x)^2)} du.$$

We are interested in  $f(t) = (\gamma(t), g(t, \cdot))$ . Plug it in, integrate by parts, take sup: **upper bound** with

$$I(\gamma) = \sup_{g \in \bar{\mathcal{S}}} I^g((\gamma(t), g(t, \cdot))).$$

Done with **local UBD!**

## Proof sketch, lower bound around a path $\gamma$

Consider a  $\gamma$  such that  $J(\gamma) < \infty$ . Then, view  $R = F_{\gamma(\cdot)}(\cdot)$  as solution of

$$R_t - (A(R)R_x)_x = h A(R) R_x.$$

That is, set  $h = \frac{R_t - (A(R)R_x)_x}{A(R)R_x}$ . This form allows for a **tilting argument**:

**Main idea**: apply **Girsanov's Theorem** to change particle system to:

$$dX_i(t) = -h(t, X_i(t))A(F_{\varrho^N(t)}(X_i(t))) dt + \sigma(F_{\varrho^N(t)}(X_i(t))) d\tilde{W}_i(t),$$

then show  $\frac{dP^N}{d\tilde{P}^N} \approx e^{-NJ(\gamma)}$  on  $\{\varrho^N \in B(\gamma, \delta)\}$  and **LLN** under  $\tilde{P}^N$ :

$$\lim_{N \rightarrow \infty} \varrho^N = \gamma.$$

## Proof sketch, lower bound around a path $\gamma$

The proof of **LLN** in the usual way:

- First, show **tightness** of  $\varrho^N$ ,  $N \in \mathbb{N}$ .
- Then, show every limit point  $\tilde{\gamma}$  corresponds to a weak solution of

$$\tilde{R}_t - (A(\tilde{R})\tilde{R}_x)_x = h A(\tilde{R}) \tilde{R}_x$$

via  $\tilde{R} = F_{\tilde{\gamma}(\cdot)}(\cdot)$ .

- Finally, show that **weak solution** of PDE **unique**, thus:  $\gamma = \tilde{\gamma}$ .

**Technical point:** for Girsanov, tightness, passing to the limit, uniqueness need:  $h \in C_b(\mathbb{R}_T)$ , Lipschitz.



## Proof sketch, lower bound around a path $\gamma$

What do we mean by  $\frac{dP^N}{d\tilde{P}^N} \approx e^{-NJ(\gamma)}$  on  $\{\varrho^N \in B(\gamma, \delta)\}$ ?

- By **Girsanov's Theorem**:

$$P(\varrho^N \in B(\gamma, \delta)) = \mathbb{E}^{\tilde{P}} \left[ e^{M(T) - \langle M \rangle(T)/2} \mathbf{1}_{\{\varrho^N \in B(\gamma, \delta)\}} \right]$$

- Apply **Hölder's inequality** to lower bound  $P(\varrho^N \in B(\gamma, \delta))$  by:

$$\mathbb{E}^{\tilde{P}} \left[ e^{-\frac{q}{p}M(T) + \frac{q}{p}\langle M \rangle(T)/2} \right]^{-p/q} \tilde{P}(\varrho^N \in B(\gamma, \delta))^p.$$

## Proof sketch, lower bound around a path $\gamma$

- Next, **complete the martingale**:

$$\mathbb{E}^{\tilde{P}} \left[ e^{-\frac{q}{p}M(T) + \frac{q}{2p}\langle M \rangle(T)} \right] = \mathbb{E}^{\tilde{P}} \left[ e^{-\frac{q}{p}M(T) - \frac{q^2}{2p^2}\langle M \rangle(T) + \left(\frac{q}{2p} + \frac{q^2}{2p^2}\right)\langle M \rangle(T)} \right].$$

- Finally,

$$N(J(\gamma) - \varepsilon) \leq \langle M \rangle(T) \leq N(J(\gamma) + \varepsilon),$$

since we work under  $\tilde{P}$  now! Remains to take limits  $N \rightarrow \infty$ ,  $\delta \downarrow 0$ ,

$p \uparrow \infty$ ,  $q \downarrow 1$ ,  $\varepsilon \downarrow 0$ . Done with **local LBD**!

# Proof sketch, comparison of rate functions I

**What did we prove? local UBD with  $I$ , local LBD with  $J$ .**

**Need to show:  $J \leq I$ :**

- Fix  $\gamma$ . Can assume  $I(\gamma) < \infty$ .
- Use  $I(\gamma) < \infty$  to deduce regularity of  $R = F_{\gamma(\cdot)}(\cdot)$ :  
 $R_t, R_x, R_{xx} \in L^{3/2}(\mathbb{R}_T)$ ,  $R_x \in L^3(\mathbb{R}_T)$ ,  $\int_{\mathbb{R}_T} \frac{R_{xx}^2}{R_x} dm < \infty$ ,  
 $\int_{\mathbb{R}_T} \frac{R_t^2}{R_x} dm < \infty$ .

- Recall  $I$  defined as supremum over  $g \in \bar{\mathcal{S}}$  of

$$\left[ (\gamma(T), g) - (\gamma(0), g) - \int_0^T (\gamma(t), g_t + A(R)g_{xx} + \frac{1}{2}A(R)(g_x)^2) dt \right],$$

would like to take  $g_x = \frac{R_t - (A(R)R_x)_x}{A(R)R_x} = h$ .

## Proof sketch, comparison of rate functions II

- This is OK due to **regularity**:  $h \in L^2(\mathbb{R}_T, R_x)$  and denseness of  $\overline{\mathcal{S}}$  in the latter.
- Done, since  $J(\gamma) = \frac{1}{4} \int_{\mathbb{R}_T} h^2 R_x \, dm$ .
- Now, redo this **with drift** and end up with

$$J(R) = \frac{1}{2} \left\| \frac{\sigma(R)}{2} \frac{R_t - (A(R)R_x)_x}{A(R)(R_x)^{1/2}} - \frac{b(R)}{\sigma(R)} (R_x)^{1/2} \right\|_{L^2(\mathbb{R}_T)}^2$$

as desired.

## Under the rug

- **Regularity results** from  $I(\gamma) < \infty$ .
- **Getting rid of the atoms:** used continuity of  $R$  at various places, e.g. uniqueness of solutions to tilted generalized porous medium equation.
- **Uniqueness of weak solutions to tilted PME.**
- **Exponential tightness.**

THANK YOU  
FOR YOUR ATTENTION!