

Lect 2

Note Title

12/1/2016

The thermodynamic limit

" $C \rightarrow \infty$ "

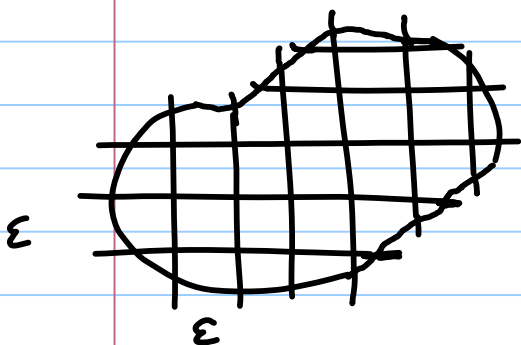
- the asymptotic of Z_{C_ε} ?
- the asymptotic of local correlation funcs.?
- the asymptotic of the probability measure?

* * *

- $\mathcal{L}$ = lattice (double periodic graph)

$$\cdot \varphi_\varepsilon: \mathcal{L} \hookrightarrow \mathbb{R}^2$$

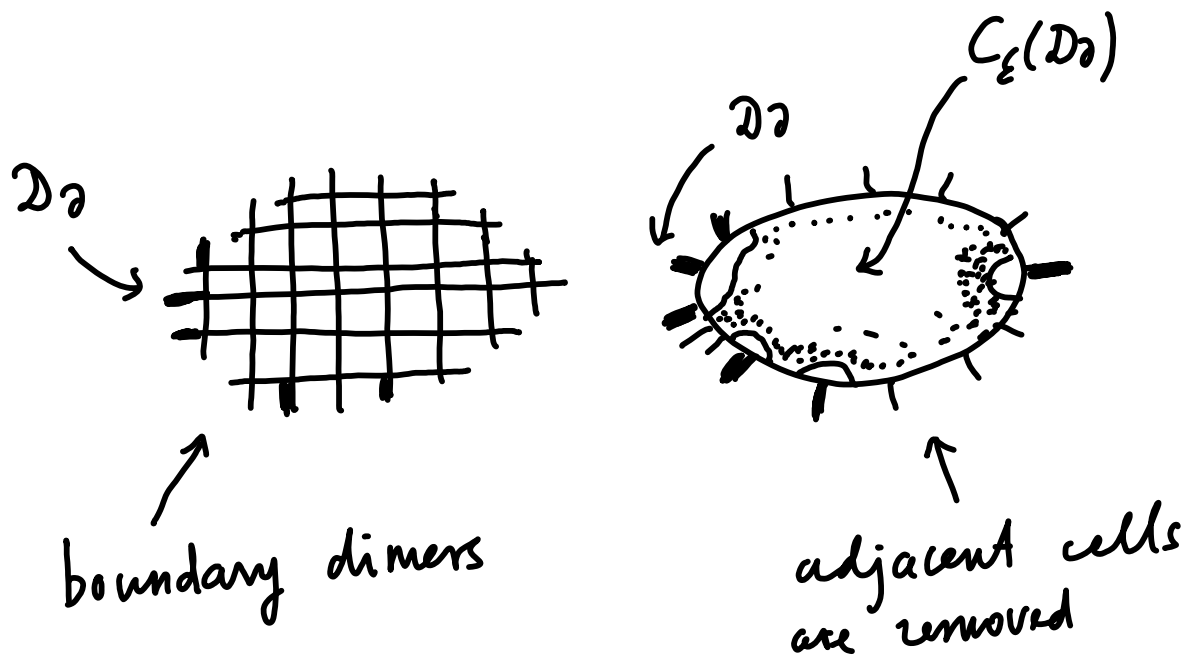
• $D \subset \mathbb{R}^2$ domain
connected, simply connected



$$\cdot C_\varepsilon = \varphi_\varepsilon(\mathcal{L}) \cap D$$

$C_\varepsilon \rightarrow \infty$ means $\{\varepsilon_n\}, \varepsilon_n \rightarrow 0$
 $n \rightarrow \infty$

Boundary conditions



Dimer exclusion rule:

$$\left(\begin{array}{c} \text{Dimers on outer} \\ \text{edges} \end{array} \right) \approx \left(\begin{array}{c} \text{adjacent } 0, 1, 2 \\ \text{cells are removed} \end{array} \right)$$

$$Z_{C_\varepsilon}(D) = Z_{C_\varepsilon(D)} = \sum_{D \subset C_\varepsilon(D)} \prod_{e \in D} w(e)$$

$$\text{Prob}(D \subset C_\varepsilon(D)) = \frac{\prod_{e \in D} w(e)}{Z_{C_\varepsilon(D)}}$$

$$\text{Prob}(h \in \mathcal{H}(C_\varepsilon(D))) = \frac{\prod_f q_f^{h_D(f)}}{\sum_{h \in \mathcal{H}(C_\varepsilon(D))} \prod_f q_f^{h(f)}}$$

$$\text{Prob}(D) = \text{Prob}(h_D)$$

$\varepsilon_n \rightarrow 0$, • stabilizing sequence of boundary conditions
 • periodic or scaling weights

$\mathcal{H}(C_\varepsilon(D\partial)) =$ the space of height functions on $C_\varepsilon(D\partial)$

$h|_{\partial C_\varepsilon(Dc)} = \text{fixed by } (C_\varepsilon, D\partial)$

Normalized height function on C_ε

$(n, m) \in \mathcal{Z} \quad \varphi_\varepsilon(n, m) = (\varepsilon n, \varepsilon m)$
 $\uparrow$ coord. of the fundamental domain

$$\varphi_\varepsilon(\varepsilon n, \varepsilon m) = \varepsilon h(n, m)$$

$$|\varphi_\varepsilon(x+\varepsilon, y) - \varphi_\varepsilon(x, y)| < \varepsilon$$

$$|\varphi_\varepsilon(x, y+\varepsilon) - \varphi_\varepsilon(x, y)| < \varepsilon$$

We are interested in the limit $\varepsilon \rightarrow 0$
 when $\varepsilon h|_{\partial C_\varepsilon(D_0)} \rightarrow \eta$ on ∂D

Continuous counterpart of $\mathcal{H}(C_\varepsilon)$

the space of height functions on D :

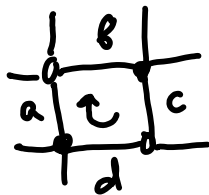
$$H(D) = \left\{ \varphi: D \rightarrow \mathbb{R} \mid \begin{array}{l} \text{Lipschitz;} \\ | \overset{\eta}{\varphi}(x+t, y) - \varphi(x, y) | \leq t, | \varphi(x, y+t) - \varphi(x, y) | \leq t \\ \varphi|_{\partial D} = \eta \end{array} \right\}$$

- The thermodynamic limit $\varepsilon_n \rightarrow 0$
 weights $w(\varepsilon)$ are double periodic,
 do not change as $\varepsilon_n \rightarrow 0$.

- The thermodynamic limit and scaling, when $\{w(e)\}$ or $\{q_f\}$ change as $\epsilon_n \rightarrow 0$

Two examples:

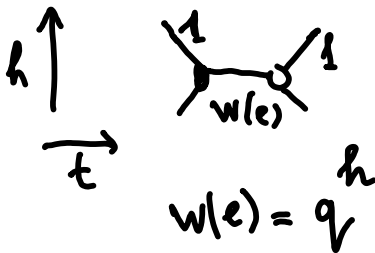
- square grid, minimal fundamental domain



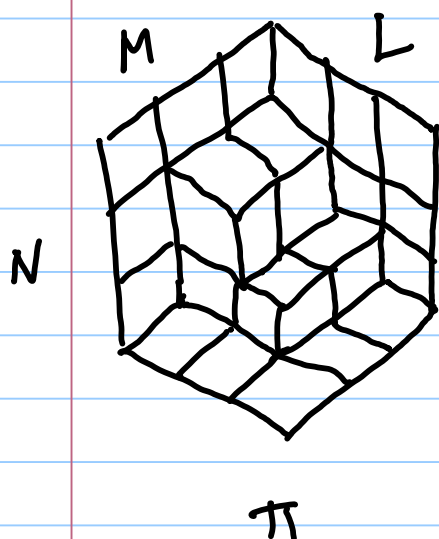
$$\begin{array}{|c|c|} \hline \circ & \circ \\ \hline \circ & \circ \\ \hline \end{array} \rightarrow a^{-1} b^{-1} c^{-1} d^{-1} = q$$

$$\begin{array}{|c|c|} \hline \circ & \circ \\ \hline \circ & \circ \\ \hline \end{array} \rightarrow \bar{a}^{-1} \bar{b}^{-1} \bar{c}^{-1} \bar{d} = \bar{q}^{-1}$$

- hexagonal lattice, minimal fund. dom.



$$\begin{array}{c} h+1 \\ \circ \\ \uparrow q \\ \circ \\ \uparrow q \\ \circ \\ \uparrow q \\ \circ \\ \uparrow q \\ \circ \\ \uparrow q \\ \circ \end{array} \rightarrow q^{-h} q^{h+1} = q$$



$$\text{Prob}(\pi) \propto q^{|\pi|}$$

Boundary conditions for $h(x,t) =$
 $=$ the shape of the lattice domain
near the boundary

Thermodynamic limit:

$$N:M:L = A:B:C, \quad N \rightarrow \infty$$

fixed ($\varepsilon \rightarrow 0$)

The limit shape phenomenon in dimer models

1) Torus

$$Z_{NM}^{\text{torus}}(H, V) = \sum_D \prod_{e \in D} w(e) e^{\frac{H \Delta h(a) + V \Delta h(b)}{e}} =$$

$$= \frac{1}{2} \left(\pm \text{Pf}(A_{K_1}) \pm \text{Pf}(A_{K_2}) \pm \text{Pf}(A_{K_3}) \pm \text{Pf}(A_{K_4}) \right)$$

$N \rightarrow \infty$, $N:M$ fixed

Matrices A_{K_i} : diagonalization by
Fourier transform.

$$Z_{NM}^{\text{torus}}(H, V) = \exp(-NM f(H, V)(1 + o(1)))$$

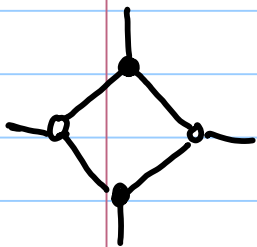
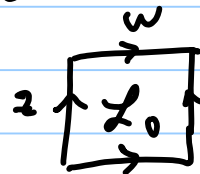
$$f(H, V) = \frac{1}{(2\pi i)^2} \oint \oint \ln |P(z, w)| \frac{dz}{z} \frac{dw}{w}$$

$$|z| = e^H$$

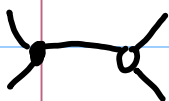
$$|w| = e^V$$

$P(z, w)$ = "spectral polynomial"
determined by the fundamental domain
of $\mathcal{L}$

$$P(z, w) = \det(A_K |_{\mathcal{L}_0}(z, w))$$



$$P(z, w) = z + \bar{z}^{-1} + w + \bar{w}^{-1} + 1$$



$$P(z, w) = 1 + z + w$$

$$Z_{NM}^{\text{torus}}[n, m] = \sum_{h \in \mathcal{H}_{\text{torus}}(n, m)} \prod_f q_f^{h(f)}$$

On a torus h is not a function:
it may have monodromies along
noncontractible cycles.

$$\mathcal{H}_{\text{torus}}[n, m] = \{h \mid \Delta_a h = n, \Delta_b h = m\}$$

Two partition functions are related

$$Z_{NM}^{\text{torus}}(H, V) = \sum_{n, m \geq 0} Z_{NM}^{\text{torus}}[n, m] e^{nH + mV} \quad (*)$$

$$\text{As } N \rightarrow \infty, N: M, s = \frac{n}{N}, t = \frac{m}{M}$$

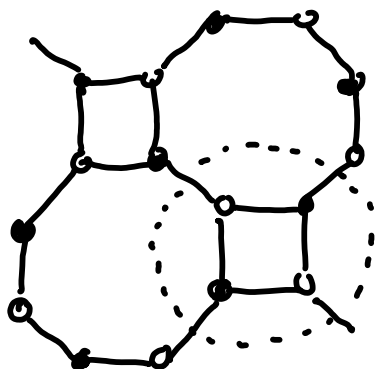
$$Z_{N,M}^{\text{torus}}[n,m] = \exp(NM \sigma(s,t)(1+o(1)))$$

(*) implies

$$f(H,V) = \max_{s,t} (Hs + Vt - \sigma(s,t))$$

inverse Legendre transform:

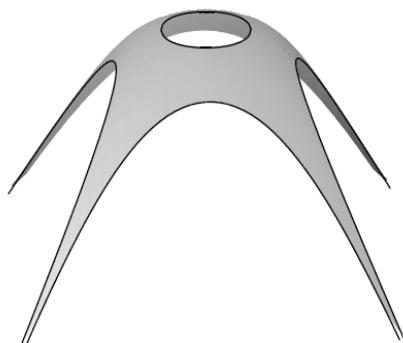
$$\sigma(s,t) = \max_{H,V} (Hs + Vt - f(H,V))$$



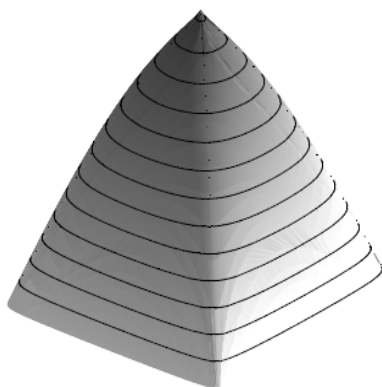
$$P(z, w) = z + \bar{z}' + w + \bar{w}' + 5$$

Kenyon
Okounkov
Sheffield

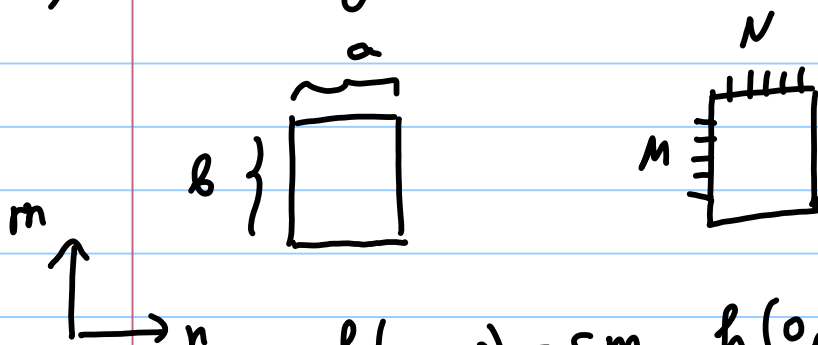
f



σ



2) Rectangular domain



$$h(m, 0) = sm, \quad h(0, n) = tn$$

$$h(m, N) = tN + sm, \quad h(M, n) = sM + tn$$

$$N = \frac{b}{\varepsilon}, \quad M = \frac{a}{\varepsilon}, \quad \text{with } s, t \text{ fixed}$$

$$\varepsilon \rightarrow 0$$

Lemma (Cohn, Kenyon, Propp) as $\varepsilon \rightarrow 0$

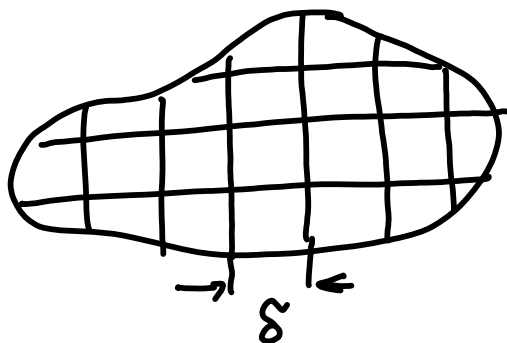
$$Z_{N,M}(n, m) = \exp\left(\frac{ab}{\varepsilon^2} \sigma(s, t)(1 + o(1))\right)$$

where

$$\sigma(s, t) = \lim_{\varepsilon \rightarrow 0} \varepsilon^2 \ln \left(Z_{N,M}^{\text{torus}}(n, m) \right)$$

3) The limit shape phenomenon

Heuristic arguments:



Partition D
into small
regions of
linear size δ

$$Z_{C_\varepsilon}(D_\varepsilon) = \sum_{\substack{\text{over} \\ \text{boundary height} \\ \text{functions (on cuts)}}} \prod_{\substack{\text{over} \\ \text{small regions}}} Z_{C_\varepsilon}(d)(h_d)$$

$$\simeq \sum_{\substack{\text{over} \\ \text{boundary height} \\ \text{functions (on cuts)}}} \exp \left(\frac{1}{\varepsilon^2} \sum_{\substack{\text{over} \\ \text{small regions}}} \Delta a \Delta b \, \sigma \left(\frac{\Delta a h}{\Delta a}, \frac{\Delta b h}{\Delta b} \right) \right)$$

$$\approx \int_{\varphi|_{\partial D} = \varphi_0} \exp\left(\frac{1}{\varepsilon^2} \iint_D \sigma(\partial_x \varphi, \partial_y \varphi) dx dy\right)$$

$$= \exp\left(\frac{1}{\varepsilon^2} \iint_D \sigma(\partial_x \varphi_0, \partial_y \varphi_0) dx dy\right)$$

φ_0 = the minimizer of

$$S[\varphi] = \iint_D \sigma(\partial_x \varphi, \partial_y \varphi) dx dy$$

on the space $\mathcal{H}(D)$ (Lipshitz)

with boundary condition $\varphi|_{\partial D} = \varphi_0$

$$\varphi_0 = \lim_{\varepsilon \rightarrow 0} \varepsilon h|_{\partial C_\varepsilon(D_0)},$$

$$\lim_{\varepsilon \rightarrow 0} \varepsilon^2 \ln Z_{C_\varepsilon(\partial D)} = \iint_D \sigma(\partial_x \varphi_0, \partial_y \varphi_0) dx dy$$

Thm (Cohn, Kenyon, Propp) This is true
and

$$\varepsilon h(\varepsilon n, \varepsilon m) \longrightarrow \varphi_0(x, y)$$

as $\varepsilon \rightarrow 0$ (in probability).

[For hexagonal lattice, uniform
distribution]

Proof. More precise version
of the above arguments.

Stronger version of this theorem:

$$\mathbb{E}[\varepsilon h(\varepsilon n_1, \varepsilon m_1) \dots \varepsilon h(\varepsilon n_k, \varepsilon m_k)] \rightarrow \\ \rightarrow \varphi_0(x_1, y_1) \dots \varphi_0(x_k, y_k)$$

For the proofs in various cases see
Borodin Ferrari [2009], Petrov [2015]
Bufetov Knizel [2016]

Fluctuations around limit
shapes. Next lecture

