

Limit shapes in integrable models in statistical mechanics

Lect 1. Local Models in
Statistical Mechanics
Dimer models

Lect 2. The Thermodynamic
Limit in Dimer models
and the Limit Shape
Phenomenon

Lect 3. Fluctuations around the
limit shape and correlation functions
in dimer models.

Lect 4 . The 6-vertex model.
Yang-Baxter equation, commutativity
and the spectrum of transfer-matrices

Lect 5. The thermodynamic limit.
Torus, cylinder, limit shapes.
Variational principle, its Hamiltonian
version. Infinitely many conservation
laws.

Lect 6. Markov properties of the
transfer-matrix and the stochastic
point. Limit shapes for the
stochastic 6-vertex model.
Fluctuations.

Lect 1.

"Lattice" models in statistical mechanics

$\dim = 2$ case

"Space time"

$$C = C_0 \sqcup C_1 \sqcup C_2 - \text{cell complex}$$

vertices edges faces
 (2-cells)

States:

States:

• State bundle $X \leftarrow \text{discrete}$ $X_c \leftarrow \left\{ \begin{array}{l} \text{space} \\ \text{of} \\ \text{states} \\ \text{on } c \end{array} \right.$

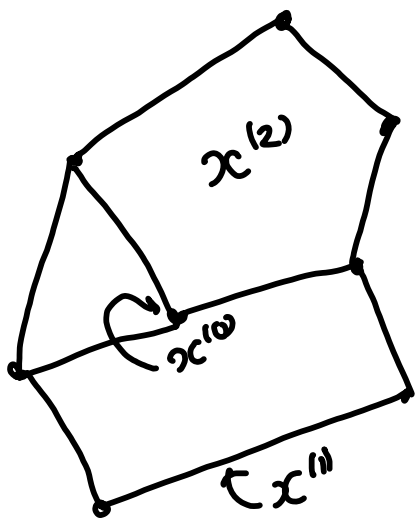
$\downarrow \pi$ $\downarrow c$

C c

• States on $C =$ sections of X :

$$x : c \mapsto x(c)$$

state on c



• Boltzmann weights

$$W: \Gamma(X, r) \rightarrow \mathbb{R}_{\geq 0}$$

$$x \xrightarrow{W} W(x) \geq 0$$

• Physical meaning:

$$W(x) = \exp\left(-\frac{E(x)}{kT}\right)$$

$E(x)$ = the energy of x

• Probability distribution:

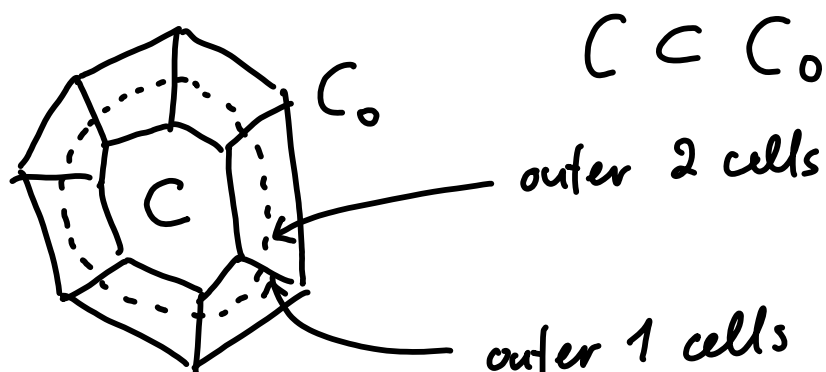
$$\text{Prob}(x) = \frac{1}{Z} W(x)$$

$$Z = \sum W(x)$$

all x with
given boundary
conditions

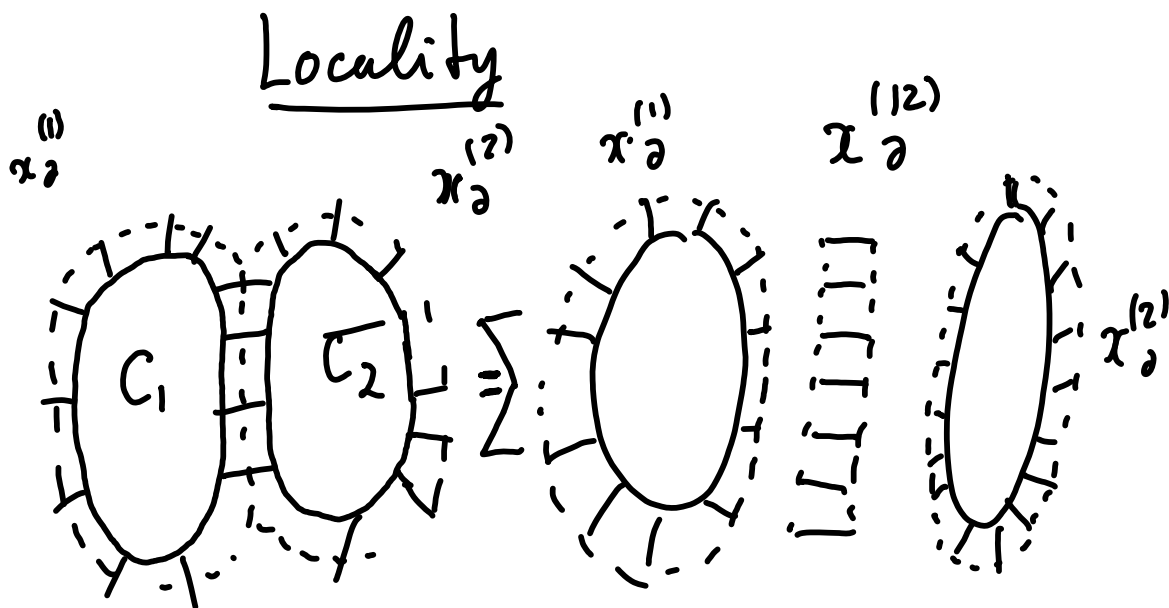
Boundary condition

Assume C_0 is planar (or a cell approximation of a surface).



Boundary states $(x_\partial) =$ states on
outer 1-cells and 2-cells

$$Z(x_\partial) = \sum_{\substack{x \text{ on } C \\ \text{states on } C}} w(x, x_\partial)$$



$$Z_{C_1 \cup C_2}(x_{\partial}^{(1)}, x_{\partial}^{(2)}) = \sum_{x_{\partial}^{(12)}} Z_{C_1}(x_{\partial}^{(1)}, x_{\partial}^{(12)}) \cdot w(x_{\partial}^{(12)}) Z_{C_2}(x_{\partial}^{(12)}, x_{\partial}^{(2)})$$

$$\cdot w(x_{\partial}^{(12)}) Z_{C_2}(x_{\partial}^{(12)}, x_{\partial}^{(2)})$$

"Minimal" (local) weights: $\left\{ \begin{array}{l} \text{depend on the} \\ \text{star of } c \end{array} \right\}$

$$w\left(\begin{array}{c} x_{e_1} \quad x_{e_2} \\ \diagup \quad \diagdown \\ x_o \end{array}\right), \quad w\left(\begin{array}{c} x_{c_+} \\ \text{---} x_e \text{---} \\ x_{e_-} \end{array}\right), \quad w\left(\begin{array}{c} x_{v_1} \quad x_{e_1} \\ \diagup \quad \diagdown \\ x_c \end{array}\right)$$

($\mathbb{R}^2$ is oriented)

Examples:

1) Dimer models:

- states on edges, occupied or non occupied
(by a dimer)



$$w(D) = \begin{cases} 0 & \text{if two dimers meet at a vertex} \\ 0 & \text{if } \exists \text{ a vertex not adjacent to a dimer edge} \\ \prod_{e \in D} w(e) & \text{otherwise} \end{cases}$$

$w: E(C) \rightarrow \mathbb{R}_{>0}$ given weights

Note: weights can be zero

2) Ising model

• States on vertices $(+, -)$

$$w(\sigma) = \exp\left(-\sum_{e \in C_1} J_e \sigma_{e+} \sigma_{e-}\right)$$

Note: all $w(\sigma) > 0$

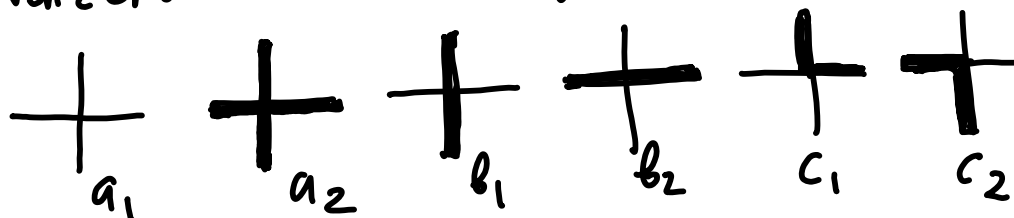
3) The 6-vertex model

C = square grid, oriented

• States on edges: occupied by a path or not occupied by a path



• Nonzero local weights:



Note: other weights are zero

Van Hove theorem

$$\mathcal{Z} = \begin{array}{|c|} \hline \text{\tiny $\times$} \\ \hline \end{array}, \quad \varphi_\varepsilon(\mathcal{Z}) \hookrightarrow \mathbb{R}^{\varepsilon} \begin{array}{|c|} \hline \text{\tiny $\times$} \\ \hline \end{array}$$

$$C_\varepsilon = \varphi_\varepsilon(Z) \cap D, \quad D \subset \mathbb{R}^2 \text{ is fixed}$$

$|C_\xi| \rightarrow \infty$, means $\varepsilon \rightarrow 0$

Assume :

- ① $w(x)$ is also periodic and $w(x)$ do not change as $\varepsilon \rightarrow 0$
- ② $w(x) \neq 0$ for all x
- ③ X is compact

Thm(VH, see Ruelle)

$$f = - \lim_{|C| \rightarrow \infty} \frac{1}{|C|} \ln(Z_C(x_0))$$

exists and does not depend on boundary conditions (free energy per site)

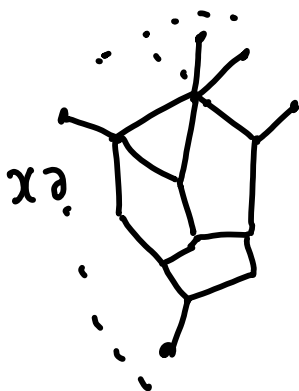
True for Ising model

Not true for dimers and 6-vertex :
the limit exists but depends
on boundary conditions.

Simplest example: Gaussian model

$$X = \mathbb{R}^{V(C)}, \quad \text{positive definite}$$

$$w(x) = \exp\left(-\frac{1}{2} \sum_{v,v'} E(v,v') x_v x_{v'}\right)$$



$$Z(x_2) = \int_{x_{int}} w(x) d x^{V_{int}(C)}$$

Gaussian integral

$$Z(x_\partial) = \frac{|V_{\text{inner}}| (2\pi)}{\sqrt{\det(E_{\text{Dirichlet}})}} \cdot \exp\left(-\frac{1}{2} \sum_{v, v' \in \partial C} x_v x_{v'} (E^{-1})_{vv'}\right)$$

if $\{E(v, v')\} = \text{discretized Laplacian on } C_\varepsilon$
 $\varepsilon \rightarrow 0$

$$\{E\} \rightarrow \Delta, \quad x_\partial \rightarrow \eta$$

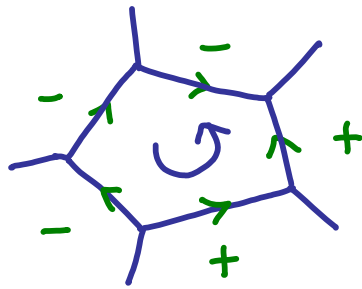
$$Z(x_\partial) = \exp\left(\frac{1}{\varepsilon^2} \int_{\partial D} \int_{\partial D} \eta(x) \eta(y) G(x, y) dx dy (1 + o(1))\right)$$

Depends on boundary conditions
 (local space of states is noncompact, $\mathbb{R}$)

Kasteleyn solution to planar dimer models

1. K-orientations of $\Gamma \subset \mathbb{R}^2$ (spin struct.)

(Kuperberg; Cimasoni, N.R.)

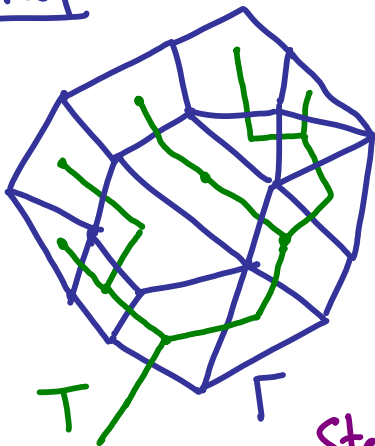


$$\prod_{e \in \partial f} \varepsilon(e) = -1$$

for each 2-cell

Prop. Such orientations exist.

Proof

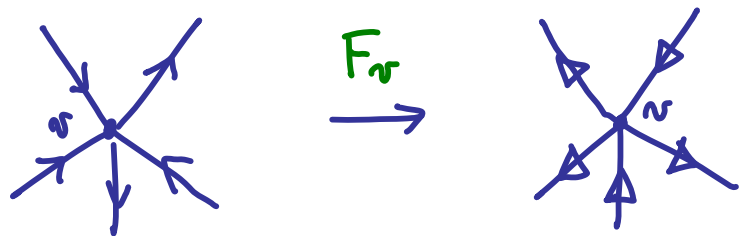


$\Gamma \subset \mathbb{R}^2$ graph, $T =$
= a spanning tree
on Γ^v

Step 1 Orient all edges
non intersecting T arbitrary

Step 2 Going from leaves to
the root orient remaining edges s.t. each
face is K-oriented. $\square$

$K \sim K'$ if K' obtained from K
by a sequence F_r :



Theorem. For each $\Gamma \subset \mathbb{R}^2 \exists!$
equivalence class of K -orientations.

2. K -matrix Given $\{w(e) \geq 0\}_{e \in E(\Gamma)}$

$$A^K: \mathbb{R}^{V(\Gamma)} \rightarrow \mathbb{R}^{V(\Gamma)}$$

(Dirac
operator)

$$A^K_{vv'} = w(v, v') \varepsilon^K_{vv'},$$

$$\varepsilon^K_{vv'} = \begin{cases} 1, & v \longrightarrow v' \\ -1, & v \longleftarrow v' \\ 0, & v \not\longleftrightarrow v' \end{cases}$$

3. Theorem (Kasteleyn 1961, Fisher 1961)

$$Z_\Gamma = \sum_{\mathcal{D} \subset \Gamma} \prod_{e \in \mathcal{D}} w(e) = |\text{Pf}(A^K)| \quad (\text{Kast})$$

Recall: $A_{ij} = -A_{ji}$, $i, j = 1, \dots, N$

$$\text{Pf}(A) = \frac{1}{2^{\frac{N}{2}} (\frac{N}{2})!} \sum_{\sigma \in S_N} (-1)^\sigma A_{\sigma_1 \sigma_2} A_{\sigma_3 \sigma_4} \cdots A_{\sigma_{N-1} \sigma_N}$$

$$= \sum_{\substack{\text{perfect matchings} \\ \text{on } (1, 2, 3, \dots, N)}} (-1)^{\sigma_m} A_{i_1 j_1} A_{i_2 j_2} \cdots A_{i_{\frac{N}{2}} j_{\frac{N}{2}}}$$

$\underbrace{(i_1 j_1)}_{w(e_{ij_1}) \in K_{ij_1}}$

$$\sigma_m = \begin{pmatrix} 1 & 2 & 3 & 4 & \cdots & N-1 & N \\ i_1 j_1 & i_2 j_2 & \cdots & i_{\frac{N}{2}} j_{\frac{N}{2}} \end{pmatrix}, \text{ represent. } m$$

$$m = \{ \sigma \in S_N \} / (i_a \leftrightarrow j_a, (i_a j_a) \leftrightarrow (i_b j_b))$$

Recall: $\sigma_e(\mathcal{D}) = \begin{cases} 1, & e \in \mathcal{D} \\ 0, & e \notin \mathcal{D} \end{cases}$

$$\langle \sigma_{e_1} \dots \sigma_{e_k} \rangle = \frac{\sum_{\mathcal{D}} \prod_{e \in \mathcal{D}} w(e) \sigma_{e_1}(\mathcal{D}) \dots \sigma_{e_k}(\mathcal{D})}{Z}$$

$$= w(e_1) \dots w(e_k) \frac{\partial}{\partial w(e_1)} \dots \frac{\partial}{\partial w(e_k)} \log Z$$

Corollary of (Kast)

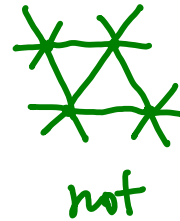
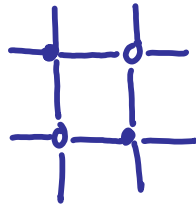
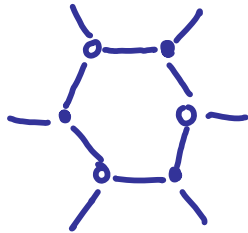
$$\langle \sigma_{i_1 j_1} \sigma_{i_2 j_2} \dots \sigma_{i_k j_k} \rangle = \text{Pf} \left(\left((A^K)^{-1} \right)_{a,b} \right)$$

$$a, b \in \{i_1, j_1, i_2, j_2, \dots, i_k, j_k\}$$

Combinatorics $\rightarrow$ Linear algebra !

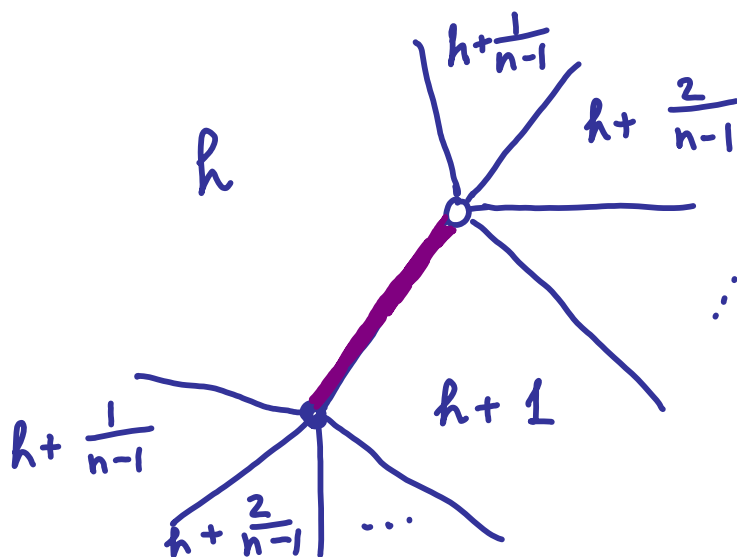
Dimers and random surfaces

$\Gamma \subset \mathbb{R}^2$, 2-cell complex, bipartite



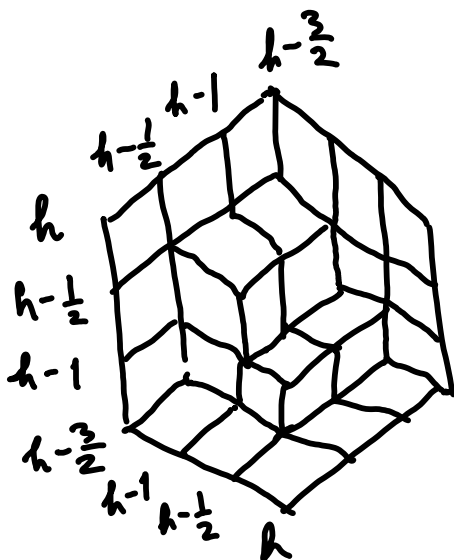
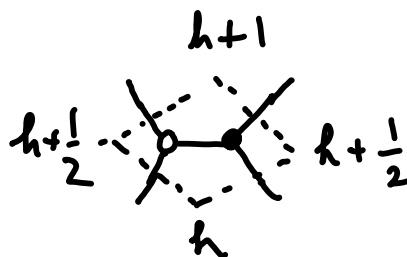
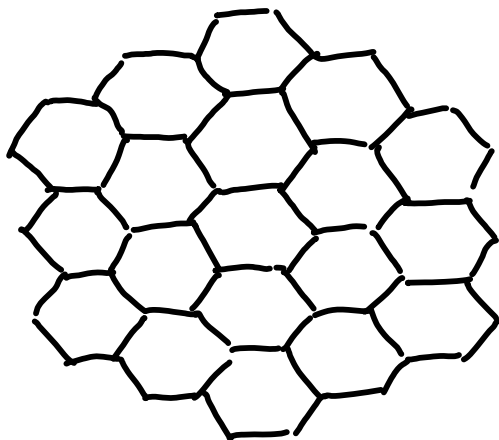
The height function

$h_D(f)$

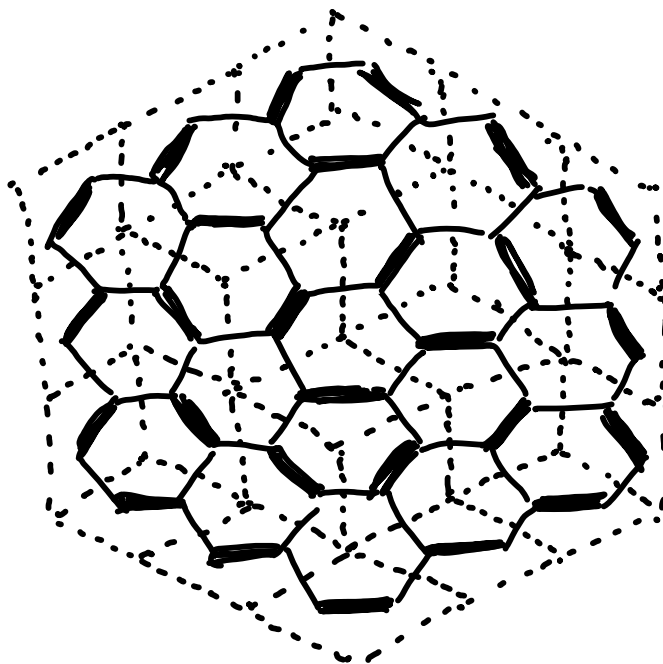


height function = discrete surface
over Γ

The space of height functions (hexagonal lattice)

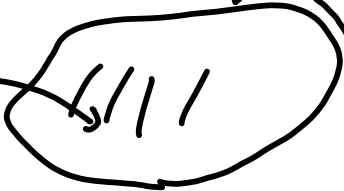


(1,1,1) projection



$$-1 + \frac{1}{2} - 1 + \frac{1}{2} - 1 + \frac{1}{2} = -\frac{3}{2}$$

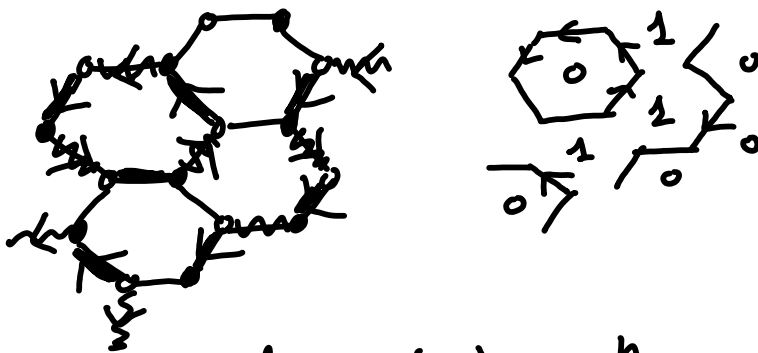
Thm. $h_D|_{\partial C}$ does not depend on D .

h_D  boundary value is determined by C

The space of height functions on C

$$H_C = \left\{ \text{fncs on faces of } C \mid h \nearrow h+1 \text{ or } h \searrow h+\frac{1}{n} \right\}$$

Remark. Relative height function



$$\bullet \quad h_{D_1, D_2}(f) = \frac{n}{n+1} (h_{D_1}(f) - h_{D_2}(f))$$

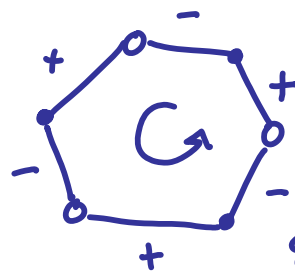
$$\bullet \quad h_{D_1, D_2}(f)|_{\partial C} = 0$$

Prob. measure on dimers $\leftrightarrow$ Prob. measure on height fncs

$$\left\{ \begin{array}{l} \text{Edge weights} \\ w(e) \\ e \in E(\Gamma) \end{array} \right\} \xrightarrow{\text{gauge}} \left\{ \begin{array}{l} \text{face weights} \\ q_f \\ f \in F(\Gamma') \end{array} \right\}$$

$$\left\{ w(e) \mapsto S(e_+) w(e) S(e_-) \right\}$$

$$q_f = \prod_{e \in \partial f} w(e)^{\varepsilon(e, f)},$$



$$\varepsilon(e, f) =$$

q_f = "gauge invariant"

= relative orientation

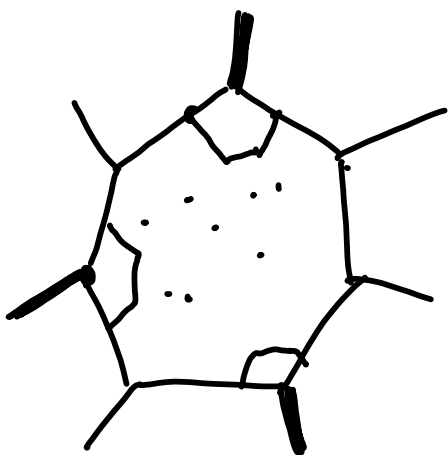
Lemma

$$\frac{\prod_{e \in \mathcal{D}} w(e)}{\sum_{\mathcal{D}'} \prod_{e \in \mathcal{D}'} w(e)} = \frac{\prod_f q_f h_{\mathcal{D}}(f)^{\frac{n+1}{n}}}{\sum_h \prod_f q_f h_{\mathcal{D}}(f)^{\frac{n+1}{n}}},$$

Proof. $\prod_{e \in D_0} w(e)^{-1} \prod_{e \in D} w(e) = \prod_C \prod_{e \in C} w(e)^{\xi(e, C)}$

$$= \prod_f q_f^{h_D D_0(f)} = \prod_f q_f^{-\frac{n+1}{n} h_{D_0}(f)} \prod_f q_f^{\frac{n+1}{n} h_D(f)}$$

Boundary conditions



dimers on outer
edges = dents
around corresponding
boundary vertices of C

$$\text{Prob}(D, D_0) = \frac{\prod_{e \in D} w(e)}{\sum_{D \text{ on } C \setminus \{D_0\text{-dents}\}} \prod_e w(e)},$$

$$\text{Prob}(h; \partial h) = \frac{\prod_{f \in C \setminus \{\text{dents}\}} q_f^{h(f)}}{\sum_{\substack{\text{height} \\ \text{functions on } C \setminus \{\text{dents}\}}} \prod_f q_f^{h(f)}}$$

↑
bijectively correspond to ∂h

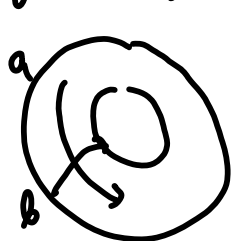
Exact solution (integrability):

- the partition function (the denominator)
 - correlation functions
- are given by Kasteleyn Pfaffians.

Remark. Surface graphs $C = (\Gamma \subset \Sigma_g)$.

Assume Γ is bipartite.

Local rules for h_D define a function only on Σ with branch cuts along cycles generating $H_1(\Sigma)$. Extra



weights

$$z_c = \exp(H_c)$$

$$Z_C = \sum_{D \subset \Gamma} \prod_{e \in D} w(e) \prod_{c \subset \Gamma} z_c^{\Delta h_D(c)}$$

$\Delta h_D(c)$ = the change of h_D along c

Thm (Kasteleyn 1963, Tesler, 2000, Galuccio, Loebl 1999, Cimasoni, R. 2006)

$$Z_C = \frac{1}{2^g} \sum_{[K]} \text{Arf}(q_{D_0}^K) \text{Pf}(A^K)$$

A^K = Kasteleyn matrix (twisted by $\{z_c\}$)

K = Kasteleyn orientation, $[K]$ its eq.cl.

D_0 = reference dimer configuration

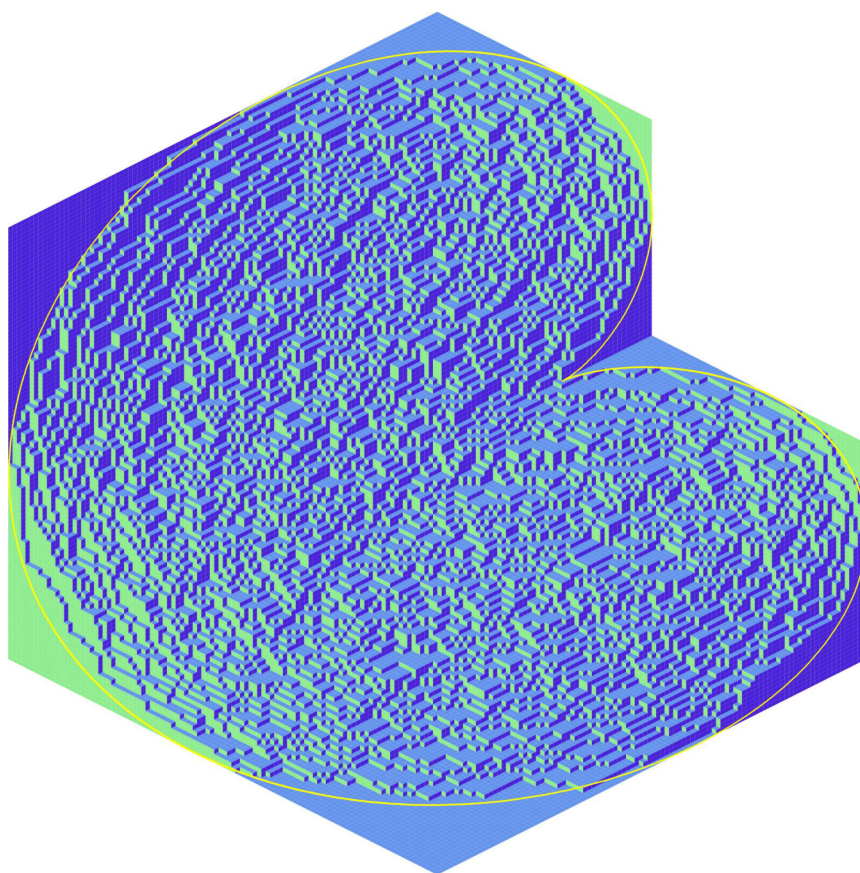
$[K, D_0]$ defines spin structure on Σ_g
(Kuperberg)

$q_{D_0}^K$ = the quadratic form corresponding
to this spin structure

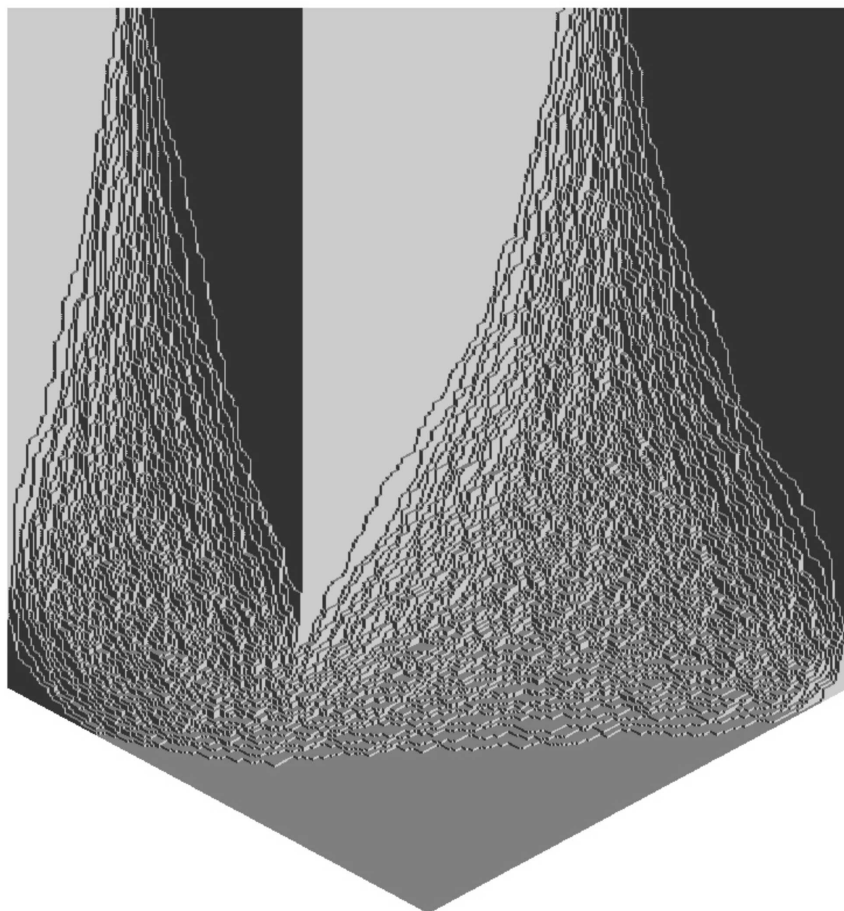
Corollary

$$Z^{\text{torus}} = \frac{1}{2} \left(\pm \text{Pf}(A_1) \pm \text{Pf}(A_2) \pm \text{Pf}(A_3) \pm \right. \\ \left. \pm \text{Pf}(A_4) \right)$$

Next two lectures:
 $|C_\varepsilon| \rightarrow \infty$ limit



uniform , $w(c) = 1$



$$q^{|\pi|}, \quad q = e^{-\varepsilon}, \quad \varepsilon > 0$$