

- $\hat{\rho}_{\text{KPP}}$  is the stationary KPP scaling function  $\langle u(x,t) u(0,0) \rangle \stackrel{\text{Burgers}}{\approx} (\Gamma t)^{-2/3} \hat{\rho}_{\text{KPP}}((\Gamma t)^{-2/3} x)$

$\hat{\rho}_{\text{KPP}}$  is the long time asymptotics of second class particle

$$\hat{\rho}_{\text{KPP}} > 0, \text{ even}, \int \hat{\rho}_{\text{KPP}}(x) dx = 1, \text{ tails } e^{-0.295 |x|^3}$$

- $\hat{\rho}_o(k) = e^{-|k|^{5/3}}$   $\propto$  -stale distribution [Baik-Rains]

#### 5.4 Euler equations

ASSUMPTION: no further conservation laws

known exceptions are harmonic  $V(x) = x^2$

$$V(x) = e^{-x}$$

Calogero-Moser  $\frac{1}{x^2}$  is not nearest neighbor, long range

3 conservation laws  $\Leftrightarrow$  infinite number

local equilibrium

$$\frac{1}{2} \exp \left[ - \sum_j \beta(\varepsilon_j) \frac{1}{2} (p_j - \alpha(\varepsilon_j))^2 + V(q_{ji} - q_{ij}) + P(\varepsilon_j) \tau_j \right]$$

as for kinetic limit  
hard to prove

maintained by flow by updating the parameters

space-time  $\varepsilon^{-1}$

$$\partial_t \langle \vec{q}_j(t) \rangle_{\tilde{\mu}} + \nabla \cdot \langle \vec{q}_j(t) \rangle_{\tilde{\mu}} = 0 \quad \tilde{\mu} = (P, v, \beta)$$

$$\partial_t \underbrace{\langle q_j \rangle_{\tilde{\mu}(t)}}_{\text{static expectation.}} + \nabla \cdot \underbrace{\langle \vec{q}_j \rangle_{\tilde{\mu}(t)}}_{\text{static expectation.}} = 0$$

macroscopic fields  $\ell, u, e_{\text{tot}}$

$$\partial_t \ell - \partial_x u = 0, \quad \partial_t u + \partial_x P(\ell, e_{\text{tot}} - \frac{1}{2} u^2) = 0, \quad \partial_t e_{\text{tot}} + \partial_x (u P(\ell, e_{\text{tot}} - \frac{1}{2} u^2)) = 0$$

$$\mathbb{P}, \beta \Leftrightarrow \ell, e_{\text{int.}}$$

$$\ell = \frac{1}{2} \int e^{-\beta(V(x) + Px)} x dx, \quad e_{\text{int.}} = \frac{1}{2\beta^2} + \frac{1}{2} \int e^{-\beta(V(x) + Px)} V(x) dx$$

linearization  $\lambda + u_1, 0 + u_2, e_{\text{tot}} + u_3$

$$\partial_t \vec{u} + \partial_x A \vec{u} = 0$$

$$A = \begin{pmatrix} 0 & -1 & 0 \\ \partial_x P & 0 & \partial_x P \\ 0 & P & 0 \end{pmatrix}$$

static correlator

eigenvalues  $0, \pm c$ ,  $c > 0$   
 $c$  speed of sound

$$C = \sum_j S(j, 0)$$

$$\text{sum rule } C = \sum_j S(j, t) \quad \text{conservation law}$$

Landau - Lifshitz

$$\partial_t u + \partial_x (Au - \partial_x Du + Bu) = 0$$

initial condition  $u(x, 0)$  white noise, mean zero

$$\langle u_\alpha(x, 0) u_{\alpha'}(x', 0) \rangle = \delta(x-x') C_{\alpha\alpha'}$$

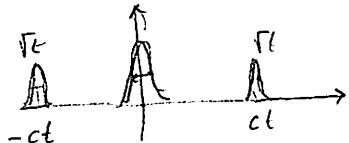
solution with  $D=B=0$

$$\langle u_\alpha(x, t) u_{\alpha'}(0, 0) \rangle = (e^{-t \partial_x A} C)_{\alpha\alpha'}(x, 0)$$



add  $D, B$

peaks broaden



empirically wrong

We have to do better!

## 5.5 Nonlinear Fluctuating Hydrodynamics

- expand currents to second order
- linear term is dominating ( $A\bar{u}$ ), switch to basis in which  $A$  is diagonal
- linear transformation (normal modes)

$$\vec{\phi} = R \vec{u}, \quad R A R^{-1} = \text{diag}(-c, 0, c), \quad R C R^T = \mathbb{1}$$

$$\partial_t \phi_\alpha + \partial_x \left( c_x \phi_\alpha + \langle \phi, G^\alpha \phi \rangle - \partial_x (D\phi)_\alpha + (B\tilde{\xi})_\alpha \right) = 0 \quad (*)$$

$$\vec{c} = (-c, 0, c), \quad G^\alpha \text{ transformed Hamiltonian,}$$

$$R^{-1} D R = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \frac{c}{2} & 0 \\ 0 & 0 & \frac{c}{2} \end{pmatrix}, \quad B B^T = 2 D$$

(\*) NFH; multi-component Burgers

•  $A$  non-degenerate,  $A = 0$  Ertas, Kadar 1993 not so much explored

• No Cole-Hopf, Haiver scheme has been worked out

Perczowski and Kupriainov, Marzocchi 2016

• stationary measures Tanaki (2016)

Gaussian white noise iff cyclicity

$$\| G_{\beta\gamma}^\alpha = G_{\beta\gamma}^\kappa = G_{\alpha\beta}^\kappa \|$$

never satisfied in physical models

## 5.6 Decoupling and mode-coupling

Decoupling hypothesis (A nondegenerate). If  $G_{\alpha\alpha}^{\times} \neq 0$ , then mode  $\alpha$  is governed by the stochastic Burgers equation as  $t, x \rightarrow \infty$ .

Implies

$$\langle \phi_{\alpha}(x, t) \phi_{\alpha}(0, 0) \rangle \approx (T_s t)^{-2/3} f_{\alpha\alpha} \left( (P_s t)^{-3/2} (x - c_{\alpha} t) \right)$$

$$\text{and } T_s = 2\sqrt{2} |G_{\alpha\alpha}^{\times}|$$

anharmonic chains

$$\text{generically } G_{\alpha\alpha}^{\times} \neq 0, \quad \alpha = \pm 1$$

BUT

$$G_{\alpha\alpha}^0 = 0 \text{ always}$$

One-loop: approximate equation for  $S$  (3x3 matrix)

$$\begin{aligned} \partial_t S(x, t) &= (-A \partial_x + D \partial_x^2) S(x, t) \\ &+ \partial_x^2 \int_0^t ds \underbrace{\int dy M(y, s) S(x-y, t-s)}_{\text{memory kernel}} \end{aligned}$$

$$M_{\alpha\alpha}(x, t) = 2 \operatorname{tr}((S G^{\times})^T (S G^{\times})) (x, t)$$

$$\Rightarrow \text{heat mode: diagonal } S_{\alpha\beta} = \delta_{\alpha\beta} f_{\alpha}$$

sound modes scale as  $k P_z$  determines the memory kernel

linear equation for  $f_0 \rightsquigarrow$

$$\hat{f}_0(k) = e^{-|k|^{5/3}}$$

input	$e^{i c k t} \hat{g}(t k^8)$	$c \neq 0$
output	$e^{-i k t^{1+\frac{1}{8}}}$	
$k P_z$	$t^{2/3} k = (t k^{3/2})^{2/3}$	$\frac{1}{8} = \frac{2}{3} \quad 1 + \frac{1}{8} = \frac{5}{3}$

## 5.7 Phase diagram

based on decoupling and mode-coupling

general discussion      Schmidt, Schütz 2016

special case       $(-c, 0, c)$        $G_{\alpha\alpha}^0 = 0$

relevant couplings are the diagonals (normal mode)

$\text{diag } G^{\sigma}$        $\sigma = \pm 1, 0$

symmetries

$$\text{diag } G^1 = (a_1, a_2, a_3), \quad \text{diag } G^{-1} = - (a_3, a_2, a_1)$$

$$\text{diag } G^0 = (-a_4, 0, a_4) \quad \text{and} \quad a_4 > 0$$

Phase 1 :       $a_1 \neq 0$       KPZ + Levy  $5/3$

Phase 2 :       $a_1 = 0$        $a_2 = 0, a_3 = 0$       diffusive + Levy  $\frac{3}{2}$  ( $= 1 + \frac{1}{2}$ )

see picture: this is the only proved case!

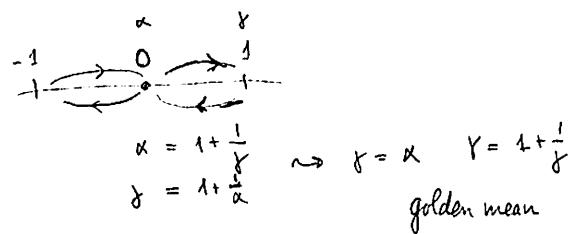
Jara, Komorowski, Olla 2014  
Bernardin, Goncalves, Jara

standard example       $V(x) = V(-x), \quad \beta = 0$

out of the blue Lee-Drossel 2015       $\beta = 1, \quad \beta = 0.5\beta, \quad V = -\frac{2}{3}x^3 + \frac{1}{4}x^4$

Phase 3 :       $a_3 = 0$ , but not  $a_1 = 0 = a_2$

all three Levy golden mean  $\frac{1}{2}(1 + \sqrt{5})$



this cannot be the case power law

towards  $|x| \rightarrow \infty$



unphysical

correct result is sound: maximally asymmetric Levy  $\frac{1}{2}(1+\sqrt{5})$

heat: symmetric Levy  $\frac{1}{2}(1+\sqrt{5})$

• only studied example C. Mendl, A. Dhar (numerics)

$$V(x) = \frac{1}{2}x^2 + \cos(\pi(x - \frac{1}{3})) + \frac{1}{8}x^4$$

$$\beta = 1, \Gamma = 2.214$$

