



Aalto University  
School of Science  
and Technology

# Conformal field theory on the lattice: from discrete complex analysis to Virasoro algebra

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joint work with

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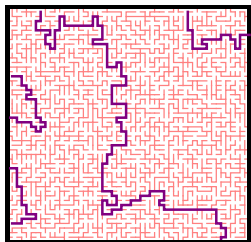
March 15, 2018 — Columbia University

# Outline

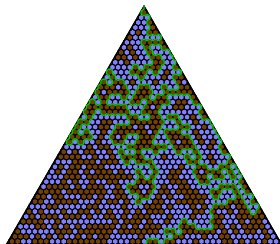
1. Introduction: Conformal Field Theory and Virasoro algebra
2. Main results: local fields of probabilistic lattice models form Virasoro representations
  - ▶ discrete Gaussian free field
  - ▶ Ising model
3. An algebraic theme and variations (Sugawara construction)
4. Proof steps (discrete complex analysis)

# 1. INTRODUCTION

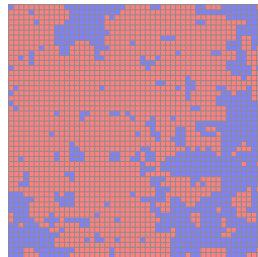
# Intro: Two-dimensional statistical physics



(uniform spanning tree)



(percolation)



(Ising model)

etc. etc.

# Intro: Conformally invariant scaling limits

**Conventional wisdom:** Any interesting scaling limit of any two-dimensional random lattice model is conformally invariant:

- ▶ *interfaces*  $\rightarrow$  SLE-type random curves
- ▶ *correlations*  $\rightarrow$  CFT correlation functions

## Remarks:

- ▶ SLE: Schramm-Loewner Evolution
  - \* is not today's topic
- ▶ CFT: Conformal Field Theory
  - \* powerful algebraic structures  
(Virasoro algebra, modular invariance, quantum groups, ...)
  - \* exact solvability (critical exponents, PDEs for correlation fns, ...)
  - \* mysteries — what is CFT, really?
- ▶ This talk: concrete probabilistic role for Virasoro algebra

# Intro: The role of Virasoro algebra

**Virasoro algebra:**  $\infty$ -dim. Lie algebra, basis  $L_n$  ( $n \in \mathbb{Z}$ ) and  $C$

$$[L_n, L_m] = (n - m)L_{n+m} + \frac{n^3 - n}{12} \delta_{n+m,0} C$$

$$[C, L_n] = 0$$

( $C$  a central element)

## Role of Virasoro algebra in CFT?

- ▶ stress tensor  $T$ : first order response to variation of metric  
(in particular “infinitesimal conformal transformations”)
- ▶ Laurent modes of stress tensor  $T(z) = \sum_{n \in \mathbb{Z}} L_n z^{-2-n}$
- ▶  $C$  acts as  $c \times \text{id}$ , with  $c \in \mathbb{R}$  the “central charge” of the CFT
- ▶ action on local fields (effect of variation of metric on correlations)
  - ▶ local fields form a Virasoro representation
  - ▶ highest weights of the representation  $\rightsquigarrow$  critical exponents
  - ▶ degenerate representations  $\rightsquigarrow$  PDEs for correlations

(exact solvability & classification)

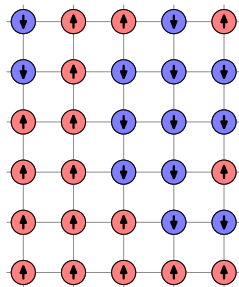
## II. LOCAL FIELDS IN LATTICE MODELS

# The critical Ising model

- ▶ domain  $\Omega \subsetneq \mathbb{C}$  open, 1-connected
- ▶  $\delta > 0$  small mesh size
- ▶ lattice approximation  $\Omega_\delta \subset \mathbb{C}_\delta := \delta\mathbb{Z}^2$

**Ising model:** random spin configuration

$$\sigma = (\sigma_z)_{z \in \mathbb{C}_\delta} \in \{+1, -1\}^{\mathbb{C}_\delta}$$



$$\sigma|_{\mathbb{C}_\delta \setminus \Omega_\delta} \equiv +1 \quad (\text{plus-boundary conditions})$$

$$P[\{\sigma\}] \propto \exp(-\beta E(\sigma)) \quad (\text{Boltzmann-Gibbs})$$

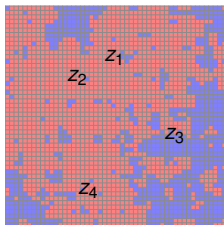
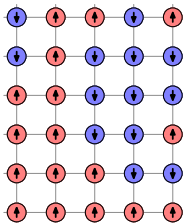
$$E(\sigma) = - \sum_{z \sim w} \sigma_z \sigma_w \quad (\text{energy})$$

$$\beta = \beta_c = \frac{1}{2} \log(\sqrt{2} + 1) \quad (\text{critical point})$$



# Celebrated scaling limits of Ising correlations

$\phi: \Omega \rightarrow \mathbb{H} = \{z \in \mathbb{C} \mid \Im m(z) > 0\}$   
conformal map



**Thm** [Chelkak & Hongler & Izuyurov, Ann. Math. 2015]

$$\lim_{\delta \rightarrow 0} \frac{1}{\delta^{k/8}} \mathbb{E} \left[ \prod_{j=1}^k \sigma_{z_j} \right]$$

$$= \prod_{j=1}^k |\phi'(z_j)|^{1/8} \times C_k(\phi(z_1), \dots, \phi(z_k))$$

**Thm** [Hongler & Smirnov, Acta Math. 2013] [Hongler, 2011]

$$\lim_{\delta \rightarrow 0} \frac{1}{\delta^m} \mathbb{E} \left[ \prod_{j=1}^m \left( -\sigma_{z_j} \sigma_{z_j + \delta} + \frac{1}{\sqrt{2}} \right) \right]$$

$$= \prod_{j=1}^m |\phi'(z_j)| \times \mathcal{E}_m(\phi(z_1), \dots, \phi(z_m))$$

# Local fields of the Ising model

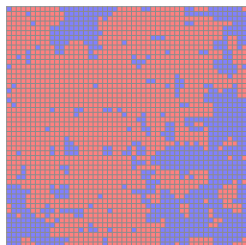
## Local fields $\mathfrak{F}(z)$ of Ising

- ▶  $V \subset \mathbb{Z}^2$  finite subset
- ▶  $P: \{+1, -1\}^V \rightarrow \mathbb{C}$  a function
- ▶  $\mathfrak{F}(z) = P((\sigma_{z+\delta x})_{x \in V})$   
(makes sense when  $\Omega_\delta$  is large enough)
- ▶  $\mathcal{F}$  space of local fields

Null fields: “zero inside correlations”

- ▶  $\mathfrak{F}(z)$  null field:  
 $\exists R > 0$  s.t.  $E \left[ \mathfrak{F}(z) \prod_{j=1}^n \sigma_{w_j} \right] = 0$   
whenever  $\|z - w_j\|_1 > R\delta \quad \forall j$
- ▶  $\mathcal{N} \subset \mathcal{F}$  space of null fields

$$\sigma = (\sigma_z)_{z \in \Omega_\delta} \quad \text{Ising}$$



Examples of local fields:

- ▶  $\mathfrak{F}(z) = \sigma_z$  (spin)
- ▶  $\mathfrak{F}(z) = -\sigma_z \sigma_{z+\delta}$  (energy)

$\mathcal{F}/\mathcal{N}$  equivalence classes of local fields (same correlations)

# Main result 1: Virasoro action on Ising local fields

Theorem (Hongler & K. & Viklund, 2017)

The space  $\mathcal{F}/\mathcal{N}$  of correlation equivalence classes of local fields of the Ising model forms a representation of the Virasoro algebra with central charge  $c = \frac{1}{2}$ .

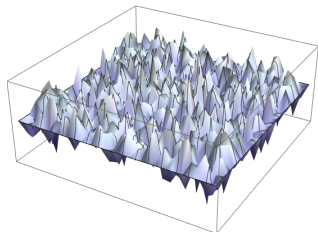
# Discrete Gaussian Free Field

Discrete Gaussian Free Field (dGFF):

$$\Phi = (\Phi(z))_{z \in \Omega_\delta}$$

Domain and discretization:

- ▶  $\Omega \subsetneq \mathbb{C}$  open, simply connected
- ▶ lattice approximation:  $\Omega_\delta \subset \mathbb{C}_\delta := \delta\mathbb{Z}^2$



- ▶ centered Gaussian field on vertices of discrete domain  $\Omega_\delta$
- ▶ probability density  $p(\phi) \propto \exp(-\frac{1}{2}E(\phi))$ 
  - ▶  $E(\phi) = \sum_{z \sim w} (\phi(z) - \phi(w))^2$

“Dirichlet energy”

# Local fields of the dGFF

## Local fields $\mathfrak{F}(z)$ of dGFF

- ▶  $V \subset \mathbb{Z}^2$  finite subset
- ▶  $P: \mathbb{R}^V \rightarrow \mathbb{C}$  polynomial function
- ▶  $\mathfrak{F}(z) = P((\Phi(z + \delta x))_{x \in V})$   
(makes sense when  $\Omega_\delta$  is large enough)
- ▶  $\mathcal{F}$  space of local fields

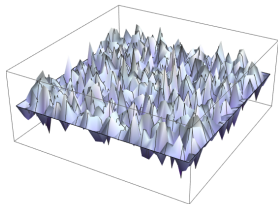
$$\Phi = (\Phi(z))_{z \in \Omega_\delta} \quad \text{dGFF}$$

## Examples of local fields:

- ▶  $\mathfrak{F}(z) = \Phi(z)$
- ▶  $\mathfrak{F}(z) = \frac{1}{2} \Phi(z + \delta) - \frac{1}{2} \Phi(z - \delta)$
- ▶  $\mathfrak{F}(z) = 361 \Phi(z)^3$

## Null fields: “zero inside correlations”

- ▶  $\mathfrak{F}(z)$  null field:  
 $\exists R > 0$  s.t.  $E \left[ \mathfrak{F}(z) \prod_{j=1}^n \sigma_{w_j} \right] = 0$   
whenever  $\|z - w_j\|_1 > R\delta \quad \forall j$
- ▶  $\mathcal{N} \subset \mathcal{F}$  space of null fields



$\mathcal{F}/\mathcal{N}$  equivalence classes of local fields (same correlations)

# Main result 2: Virasoro action on dGFF local fields

Theorem (Hongler & K. & Viklund, 2017)

The space  $\mathcal{F}/\mathcal{N}$  of correlation equivalence classes of local fields of the dGFF forms a representation of the Virasoro algebra with central charge  $c = 1$ .

# III. AN ALGEBRAIC THEME AND VARIATIONS (SUGAWARA CONSTRUCTION)

# Bosonic Sugawara construction

commutator  $[A, B] := A \circ B - B \circ A$

## Proposition (bosonic Sugawara construction)

Suppose:

- ▶  $V$  vector space and  $\alpha_j: V \rightarrow V$  linear for each  $j \in \mathbb{Z}$
- ▶  $\forall v \in V \exists N \in \mathbb{Z} : j \geq N \implies \alpha_j v = 0$
- ▶  $[\alpha_i, \alpha_j] = i \delta_{i+j,0} \text{id}_V$

Define:

$$L_n := \frac{1}{2} \sum_{j < 0} \alpha_j \circ \alpha_{n-j} + \frac{1}{2} \sum_{j \geq 0} \alpha_{n-j} \circ \alpha_j \quad \text{for } n \in \mathbb{Z}$$

Then:

- ▶  $L_n: V \rightarrow V$  is well defined
- ▶  $[L_n, L_m] = (n - m) L_{n+m} + \frac{n^3 - n}{12} \delta_{n+m,0} \text{id}_V$

$\therefore V$  Virasoro representation, central charge  $c = 1$



# Fermionic Sugawara construction 1

commutator  $[A, B] := A \circ B - B \circ A$

anticommutator  $[A, B]_+ := A \circ B + B \circ A$

## Proposition (fermionic Sugawara, Neveu-Schwarz sector)

Suppose:

- ▶  $V$  vector space,  $b_k: V \rightarrow V$  linear for each  $k \in \mathbb{Z} + \frac{1}{2}$
- ▶  $\forall v \in V \exists N \in \mathbb{Z} : k \geq N \implies b_k v = 0$
- ▶  $[b_k, b_\ell]_+ = \delta_{k+\ell, 0} \text{id}_V$

Def.:  $L_n := \frac{1}{2} \sum_{k>0} \left(\frac{1}{2} + k\right) b_{n-k} b_k - \frac{1}{2} \sum_{k<0} \left(\frac{1}{2} + k\right) b_k b_{n-k} \quad (n \in \mathbb{Z})$

Then:

- ▶  $L_n: V \rightarrow V$  is well defined
- ▶  $[L_n, L_m] = (n - m) L_{n+m} + \frac{n^3 - n}{24} \delta_{n+m, 0} \text{id}_V$

$\therefore V$  Virasoro representation, central charge  $c = \frac{1}{2}$

# Fermionic Sugawara construction 2

commutator  $[A, B] := A \circ B - B \circ A$

anticommutator  $[A, B]_+ := A \circ B + B \circ A$

## Proposition (fermionic Sugawara, Ramond sector)

Suppose:

- ▶  $V$  vector space,  $b_j: V \rightarrow V$  linear for each  $j \in \mathbb{Z}$
- ▶  $\forall v \in V \exists N \in \mathbb{Z} : j \geq N \implies b_j v = 0$
- ▶  $[b_i, b_j]_+ = \delta_{i+j,0} \text{id}_V$

Def.:

$$L_n := \frac{1}{2} \sum_{j \geq 0} \left(\frac{1}{2} + j\right) b_{n-j} b_j - \frac{1}{2} \sum_{j < 0} \left(\frac{1}{2} + j\right) b_j b_{n-j} \quad (n \in \mathbb{Z} \setminus \{0\})$$

$$L_0 := \frac{1}{2} \sum_{j > 0} j b_{-j} b_j + \frac{1}{16} \text{id}_V$$

Then:

- ▶  $L_n: V \rightarrow V$  is well defined
- ▶  $[L_n, L_m] = (n - m) L_{n+m} + \frac{n^3 - n}{24} \delta_{n+m,0} \text{id}_V$

$\therefore V$  Virasoro representation, central charge  $c = \frac{1}{2}$

# IV. PROOF STEPS (DISCRETE COMPLEX ANALYSIS)

# Outline / steps

For the Ising model and discrete GFF:

- ✓ 0.) Define model and local fields
- 1.) Suitable discrete contour integrals and residue calculus
- 2.) Introduce discrete holomorphic observable
- 3.) Define Laurent modes of the observable
- 4.) Commutation relations of Laurent modes
- 5.) Virasoro action through Sugawara construction

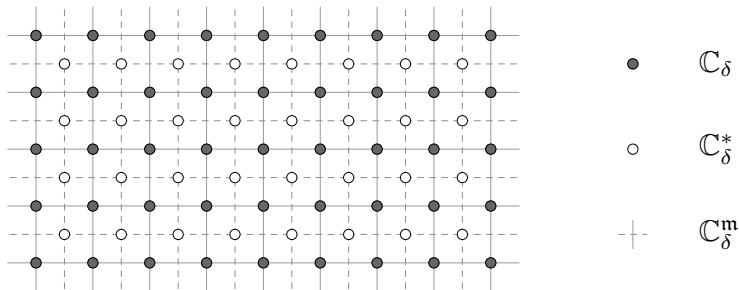
# Outline / steps

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- ✓ 0.) Define model and local fields
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# Lattices (square lattice and related lattices)

- ▶ fix small mesh size  $\delta > 0$

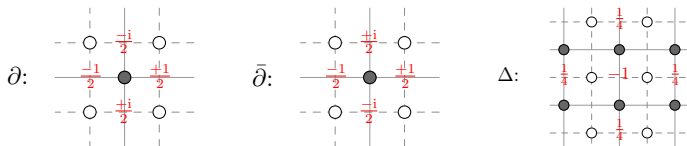


- ▶ square lattice  $\mathbb{C}_\delta$
- ▶ dual lattice  $\mathbb{C}_\delta^*$

$$\mathbb{C}_\delta = \delta\mathbb{Z}^2$$

- ▶ medial lattice  $\mathbb{C}_\delta^m$
- ▶ diamond lattice  $\mathbb{C}_\delta^\diamond$
- ▶ corner lattice  $\mathbb{C}_\delta^c$

# Lattices (discretization of differential operators)



► Discrete  $\partial$  and  $\bar{\partial}$ :

$$\partial_{\delta} f(z) = \frac{1}{2} \left( f\left(z + \frac{\delta}{2}\right) - f\left(z - \frac{\delta}{2}\right) \right) - \frac{i}{2} \left( f\left(z + \frac{i\delta}{2}\right) - f\left(z - \frac{i\delta}{2}\right) \right)$$

$$\bar{\partial}_{\delta} f(z) = \frac{1}{2} \left( f\left(z + \frac{\delta}{2}\right) - f\left(z - \frac{\delta}{2}\right) \right) + \frac{i}{2} \left( f\left(z + \frac{i\delta}{2}\right) - f\left(z - \frac{i\delta}{2}\right) \right)$$

$$f: \mathbb{C}_{\delta}^m \rightarrow \mathbb{C} \Rightarrow \partial_{\delta} f, \bar{\partial}_{\delta} f: \mathbb{C}_{\delta}^{\circ} \rightarrow \mathbb{C}$$

$$f: \mathbb{C}_{\delta}^{\circ} \rightarrow \mathbb{C} \Rightarrow \partial_{\delta} f, \bar{\partial}_{\delta} f: \mathbb{C}_{\delta}^m \rightarrow \mathbb{C}$$

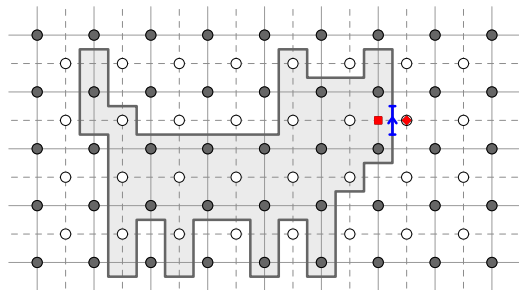
► Discrete Laplacian:




$$(\Delta_{\delta} = \bar{\partial}_{\delta} \partial_{\delta} = \partial_{\delta} \bar{\partial}_{\delta})$$

$$\Delta_{\delta} f(z) = -f(z) + \frac{1}{4} f(z + \delta) + \frac{1}{4} f(z - \delta) + \frac{1}{4} f(z + i\delta) + \frac{1}{4} f(z - i\delta)$$

# Discrete residue calculus (contour integral)

- ▶ two functions  $f: \mathbb{C}_\delta^m \rightarrow \mathbb{C}$  and  $g: \mathbb{C}_\delta^\diamond \rightarrow \mathbb{C}$
- ▶  $\gamma$  path on the corner lattice  $\mathbb{C}_\delta^c$



-  oriented edge of  $\mathbb{C}_\delta^c$
-   $f$  defined on  $\mathbb{C}_\delta^m$
-   $g$  defined on  $\mathbb{C}_\delta^\diamond$

$$\oint_{[\gamma]} f(z_m) g(z_\diamond) [dz]_\delta := \sum_{\vec{e} \in \gamma} f(e_m) g(e_\diamond) \cdot \vec{e}$$



# Discrete residue calculus (properties)

$$\oint_{[\gamma]} f(z_m) g(z_\diamond) [dz]_\delta := \sum_{\vec{e} \in \gamma} f(e_m) g(e_\diamond) \cdot \vec{e}$$

## Proposition (properties of discrete contour integral)

- ▶ **Green's formula** (sum over  $w_m \in \mathbb{C}_\delta^m \cap \text{int}(\gamma)$  and  $w_\diamond \in \mathbb{C}_\delta^\diamond \cap \text{int}(\gamma)$ )

$$\oint_{[\gamma]} f(z_m) g(z_\diamond) [dz]_\delta = i \sum_{w_m} f(w_m) (\bar{\partial}_\delta g)(w_m) + i \sum_{w_\diamond} (\bar{\partial}_\delta f)(w_\diamond) g(w_\diamond)$$

- ▶ **contour deformation**

$\gamma_1, \gamma_2$  two counterclockwise closed contours on  $\mathbb{C}_\delta^\xi$   
 $\bar{\partial}_\delta f \equiv 0$  and  $\bar{\partial}_\delta g \equiv 0$  on  $\text{symm. diff. int}(\gamma_1) \oplus \text{int}(\gamma_2)$

$$\oint_{[\gamma_1]} f(z_m) g(z_\diamond) [dz]_\delta = \oint_{[\gamma_2]} f(z_m) g(z_\diamond) [dz]_\delta$$

- ▶ **integration by parts**

$\gamma$  counterclockwise closed contour on  $\mathbb{C}_\delta^\xi$   
 $\bar{\partial}_\delta f \equiv 0$  and  $\bar{\partial}_\delta g \equiv 0$  on neighbors of  $\gamma$

$$\oint_{[\gamma]} (\partial_\delta f)(z_m) g(z_\diamond) [dz]_\delta = - \oint_{[\gamma]} f(z_m) (\partial_\delta g)(z_\diamond) [dz]_\delta$$

# Discrete monomial functions (defining properties)

## Proposition (discrete monomial functions)

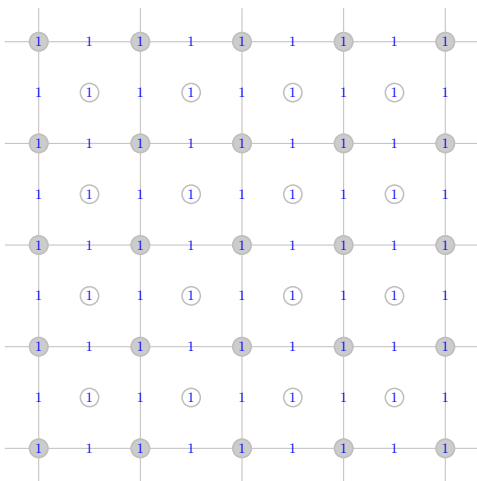
$\exists$  functions  $z \mapsto z^{[\rho]}$ ,  $\rho \in \mathbb{Z}$ , defined on  $\mathbb{C}_\delta^\diamond \cup \mathbb{C}_\delta^m$ , such that

- ▶  $\bar{\partial}_\delta z^{[\rho]} = 0$  whenever ... "discrete holomorphicity"
  - ▶  $\rho \geq 0$  and  $z \in \mathbb{C}_\delta^\diamond \cup \mathbb{C}_\delta^m$
  - ▶  $\rho < 0$  and  $z \in \mathbb{C}_\delta^\diamond \cup \mathbb{C}_\delta^m$ ,  $\|z\|_1 > R_\rho \delta$
- ▶  $z^{[0]} \equiv 1$  for all  $z \in \mathbb{C}_\delta^\diamond \cup \mathbb{C}_\delta^m$  "constant function"
- ▶  $\bar{\partial}_\delta z^{[-1]} = 2\pi \delta_{z,0} + \frac{\pi}{2} \sum_{x \in \{\pm \frac{\delta}{2}, \pm i \frac{\delta}{2}\}} \delta_{z,x}$  " $\bar{\partial}$  Green's function"
- ▶  $\partial_\delta z^{[\rho]} = \rho z^{[\rho-1]}$  "derivatives"
- ▶  $z^{[\rho]}$  has the same  $90^\circ$  rotation symmetry as  $z^\rho$  "symmetry"
- ▶ for  $\rho < 0$  we have  $z^{[\rho]} \rightarrow 0$  as  $\|z\| \rightarrow \infty$  "decay"
- ▶ for any  $z$  there exists  $D_z$  such that  $z^{[\rho]} = 0$  for  $\rho \geq D_z$  "truncation"

For  $\gamma$  large enough counterclockwise closed contour surrounding the origin...

- ▶  $\oint_{[\gamma]} z_m^{[\rho]} z_\diamond^{[q]} [dz]_\delta = 2\pi i \delta_{\rho+q,-1}$  "residue calculus"
- ▶  $\oint_{[\gamma]} z_m^{\{\rho\}} z_\diamond^{[q]} [dz]_\delta = 2\pi i \delta_{\rho+q,-1}$  where  $z_m^{\{\rho\}} = \frac{1}{4} \sum_{x \in \{\pm \frac{\delta}{2}, \pm i \frac{\delta}{2}\}} (z_m - x)^{[\rho]}$

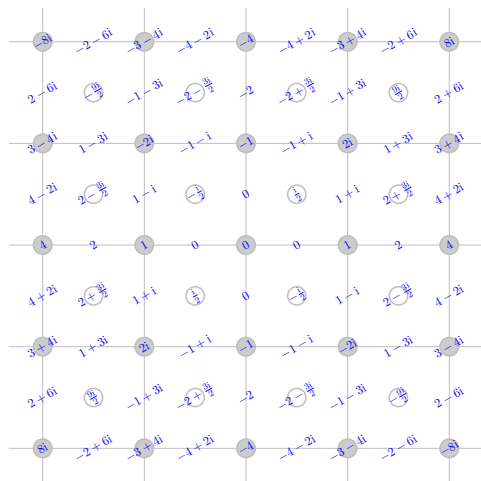
# Discrete monomial functions (example 0)



values of  $z^{[0]}$

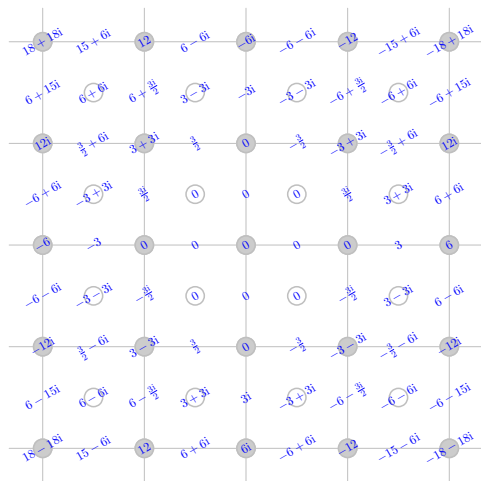


# Discrete monomial functions (example 2)



values of  $z^2$

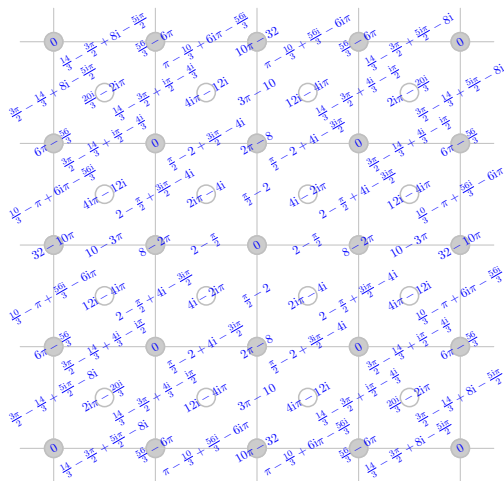
# Discrete monomial functions (example 3)



values of  $z^{[3]}$



# Discrete monomial functions (example -2)



values of  $z^{-2}$



# Outline / steps

For the Ising model and discrete GFF:

- ✓ 0.) Define model and local fields
- ✓ 1.) Suitable discrete contour integrals and residue calculus
- 2.) Introduce discrete holomorphic observable
- 3.) Define Laurent modes of the observable
- 4.) Commutation relations of Laurent modes
- 5.) Virasoro action through Sugawara construction

# Discrete Gaussian Free Field (definition again)

Discrete Gaussian Free Field (dGFF):

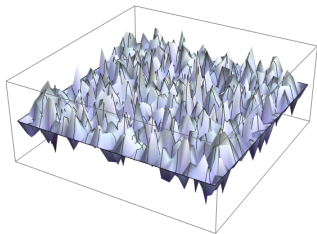
$$\Phi = (\Phi(z))_{z \in \Omega_\delta}$$

Domain and discretization:

- ▶  $\Omega \subsetneq \mathbb{C}$  open, simply connected

- ▶ lattice approximation

$$\Omega_\delta \subset \mathbb{C}_\delta, \quad \Omega_\delta^\circ \subset \mathbb{C}_\delta^\circ, \quad \Omega_\delta^m \subset \mathbb{C}_\delta^m, \quad \Omega_\delta^c \subset \mathbb{C}_\delta^c$$



- ▶ centered Gaussian field on vertices of discrete domain  $\Omega_\delta$

- ▶ probability density  $p(\phi) \propto \exp\left(-\frac{1}{2}E(\phi)\right)$

- ▶  $E(\phi) = \sum_{z \sim w} (\phi(z) - \phi(w))^2$

“Dirichlet energy”

- ▶ covariance  $E[\Phi(z)\Phi(w)] = G_{\Omega_\delta}(z, w)$

- ▶  $G_{\Omega_\delta}(z, w) =$  expected time at  $w$  for random walk from  $z$  before exiting  $\Omega_\delta$

- ▶  $\Delta_\delta G(\cdot, w) = -\delta_w(\cdot)$

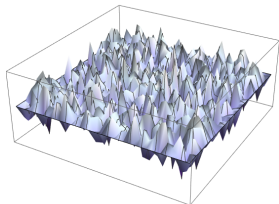
“ $\Delta$  Green’s function”

# Discrete holomorphic current (definition)

Discrete Gaussian Free Field (dGFF):

$$\Phi = (\Phi(z))_{z \in \mathbb{C}_\delta}$$

- ▶ originally defined on  $\Omega_\delta \subset \mathbb{C}_\delta$
- ▶ extend as zero to  $\mathbb{C}_\delta \setminus \Omega_\delta$  and  $\mathbb{C}_\delta^*$
- ↪ centered Gaussian field on  $\mathbb{C}_\delta^\circ$
- ▶ covariance  $E[\Phi(z)\Phi(w)] = G_{\Omega_\delta}(z, w)$



Discrete holomorphic current  $\mathfrak{J} = (\mathfrak{J}(z))_{z \in \mathbb{C}_\delta^m}$

$$\mathfrak{J}(z) := \partial_\delta \Phi(z)$$

$$= \frac{1}{2} \underbrace{\left( \Phi\left(z + \frac{\delta}{2}\right) - \Phi\left(z - \frac{\delta}{2}\right) \right)}_{\text{vanishes if } z \text{ on vertical edge}} - \frac{i}{2} \underbrace{\left( \Phi\left(z + \frac{i\delta}{2}\right) - \Phi\left(z - \frac{i\delta}{2}\right) \right)}_{\text{vanishes if } z \text{ on horizontal edge}}$$

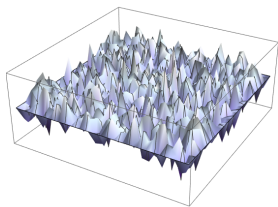
- ▶ centered complex Gaussian field (... and a local field of dGFF!)
- ▶ purely real on horizontal edges, imaginary on vertical edges
- ▶ covariance  $E[\mathfrak{J}(z)\mathfrak{J}(w)] = \partial_\delta^{(z)} \partial_\delta^{(w)} G_{\Omega_\delta}(z, w)$

# Discrete holomorphic current (correlations)

$$\Phi = (\Phi(z))_{z \in \mathbb{C}_\delta}$$

Wick's formula for centered Gaussians:

$$\mathbb{E} \left[ \prod_{j=1}^n \Phi(z_j) \right] = \sum_{\substack{P \text{ pairing} \\ \text{of } \{1, \dots, n\}}} \prod_{\{k, l\} \in P} \underbrace{\mathbb{E}[\Phi(z_k)\Phi(z_l)]}_{G_{\Omega_\delta}(z_k, z_l)}$$



Discrete holomorphic current  $\mathfrak{J} = (\mathfrak{J}(z))_{z \in \mathbb{C}_\delta^m}$ ,  $\mathfrak{J}(z) = \partial_\delta \Phi(z)$

Proposition (harmonicity of  $\Phi$ , holomorphicity of  $\mathfrak{J}$ )

- ▶  $\mathbb{E}[(\Delta_\delta \Phi)(z) \prod_{j=1}^n \Phi(w_j)] = 0$  when  $\|z - w_j\|_1 > \delta$  for all  $j$
- ▶  $\mathbb{E}[(\bar{\partial}_\delta \mathfrak{J})(z) \prod_{j=1}^n \Phi(w_j)] = 0$  when  $\|z - w_j\|_1 > \delta$  for all  $j$

$\therefore \bar{\partial}_\delta \mathfrak{J} = \Delta_\delta \Phi$  is a null field

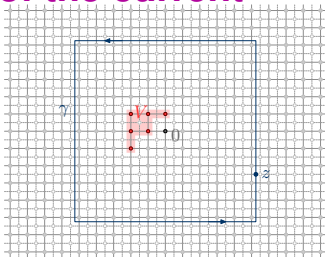
# Outline / steps

For the Ising model and discrete GFF:

- ✓ 0.) Define model and local fields
- ✓ 1.) Suitable discrete contour integrals and residue calculus
- ✓ 2.) Introduce discrete holomorphic observable
  - 3.) Define Laurent modes of the observable
  - 4.) Commutation relations of Laurent modes
  - 5.) Virasoro action through Sugawara construction

# Discrete Laurent modes of the current

- ▶  $\mathfrak{F}(w) = F[(\Phi(w+x\delta))_{x \in V}]$   
local field of dGFF
- ▶  $\gamma$  sufficiently large  
counterclockwise  
closed path on  $\mathbb{C}_\delta^c$   
surrounding origin and  $V\delta$



For  $j \in \mathbb{Z}$  define a new local field of dGFF  $(\mathfrak{J}_j \mathfrak{F})(w)$  by

$$(\mathfrak{J}_j \mathfrak{F})(0) := \frac{1}{\sqrt{2\pi}} \oint_{[\gamma]} \mathfrak{J}(z_m) z_\diamond^{[j]} \mathfrak{F}(0) [dz]_\delta$$

## Lemma (discrete current modes)

$\mathfrak{J}_j: \mathcal{F}/\mathcal{N} \rightarrow \mathcal{F}/\mathcal{N}$  is well-defined

independent of choice of ...

- ▶ add null field to  $\mathfrak{F}(0) \rightsquigarrow$  add null field to  $(\mathfrak{J}_j \mathfrak{F})(0)$  ... representative
- ▶ change  $\gamma \rightsquigarrow$  add null fields  $\bar{\partial}_\delta \mathfrak{J}(z) \times (\dots)$  to  $(\mathfrak{J}_j \mathfrak{F})(0)$  ... contour

# Outline / steps

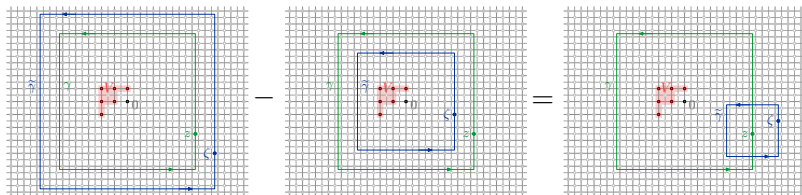
For the Ising model and discrete GFF:

- ✓ 0.) Define model and local fields
- ✓ 1.) Suitable discrete contour integrals and residue calculus
- ✓ 2.) Introduce discrete holomorphic observable
- ✓ 3.) Define Laurent modes of the observable
- 4.) **Commutation relations of Laurent modes**
- 5.) Virasoro action through Sugawara construction

# Commutation of modes of the discrete current

Proposition (commutation of discrete current modes)

$$[\mathfrak{J}_i, \mathfrak{J}_j] = i \delta_{i+j,0} \text{id}_{\mathcal{F}/\mathcal{N}}$$



$$\begin{aligned} & \mathbb{E} \left[ \left( \mathfrak{J}_i \mathfrak{J}_j \mathfrak{F}(0) - \mathfrak{J}_j \mathfrak{J}_i \mathfrak{F}(0) \right) \cdots \right] \\ &= i \delta_{i+j,0} \mathbb{E} \left[ \mathfrak{F}(0) \cdots \right] \end{aligned}$$

(residue calculus)



# Outline / steps

For the Ising model and discrete GFF:

- ✓ 0.) Define model and local fields
- ✓ 1.) Suitable discrete contour integrals and residue calculus
- ✓ 2.) Introduce discrete holomorphic observable
- ✓ 3.) Define Laurent modes of the observable
- ✓ 4.) Commutation relations of Laurent modes
- 5.) Virasoro action through Sugawara construction

# Sugawara construction with the dGFF current

Verify assumptions:

- ▶  $V$  vector space  
space  $\mathcal{F}/\mathcal{N}$  of local fields modulo null fields
- ▶  $\alpha_j: V \rightarrow V$  linear for each  $j \in \mathbb{Z}$   
discrete current Laurent mode  $\tilde{\mathfrak{J}}_j: \mathcal{F}/\mathcal{N} \rightarrow \mathcal{F}/\mathcal{N}$   
 $(\tilde{\mathfrak{J}}_j \tilde{\mathfrak{F}})(0) := \frac{1}{\sqrt{2\pi}} \oint_{[\gamma]} \tilde{\mathfrak{J}}(z_m) z_\diamond^{[j]} \tilde{\mathfrak{F}}(0) [dz]$
- ▶  $\forall v \in V \exists N \in \mathbb{Z} : j \geq N \implies \alpha_j v = 0$   
monomial truncation:  $\forall z_\diamond \in \mathbb{C}_\diamond^\circ \exists D : j \geq D \implies z_\diamond^{[j]} = 0$
- ▶  $[\alpha_i, \alpha_j] = i \delta_{i+j,0} \text{id}_V$   
Laurent mode commutation  $[\tilde{\mathfrak{J}}_i, \tilde{\mathfrak{J}}_j] = i \delta_{i+j,0} \text{id}_{\mathcal{F}/\mathcal{N}}$

## Theorem (Virasoro action for dGFF)

$$\mathfrak{L}_n := \frac{1}{2} \sum_{j < 0} \tilde{\mathfrak{J}}_j \circ \tilde{\mathfrak{J}}_{n-j} + \frac{1}{2} \sum_{j \geq 0} \tilde{\mathfrak{J}}_{n-j} \circ \tilde{\mathfrak{J}}_j$$

defines Virasoro representation with  $c = 1$  on the space  $\mathcal{F}/\mathcal{N}$  of correlation equivalence classes of local fields of the dGFF.

# Outline / steps

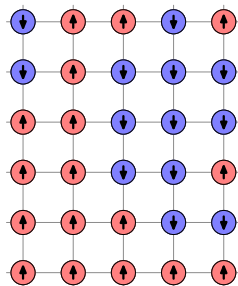
For the **Ising model** and ~~discrete GFF~~:

- ✓ 0.) Define model and local fields
- ✓ 1.) Suitable discrete contour integrals and residue calculus
- 2.) **Introduce discrete holomorphic observable**
- 3.) Define Laurent modes of the observable
- 4.) Commutation relations of Laurent modes
- 5.) Virasoro action through Sugawara construction

# The critical Ising model

- ▶  $\Omega \subsetneq \mathbb{C}$  open, 1-connected
- ▶ lattice approximation  $\Omega_\delta \subset \mathbb{C}_\delta$ ,  $\Omega_\delta^\circ \subset \mathbb{C}_\delta^\circ$ ,  $\Omega_\delta^m \subset \mathbb{C}_\delta^m$ ,  $\Omega_\delta^c \subset \mathbb{C}_\delta^c$

## Ising model: random spin configuration



$$\sigma = (\sigma_z)_{z \in \mathbb{C}_\delta} \in \{+1, -1\}^{\mathbb{C}_\delta}$$

- ▶  $\sigma|_{\mathbb{C}_\delta \setminus \Omega_\delta} \equiv +1$  (plus-boundary conditions)

$$P[\{\sigma\}] \propto \exp(-\beta E(\sigma)) \quad (\text{Boltzmann-Gibbs})$$

$$E(\sigma) = - \sum_{z \sim w} \sigma_z \sigma_w \quad (\text{energy})$$

$$\beta = \beta_c = \frac{1}{2} \log(\sqrt{2} + 1) \quad (\text{critical point})$$

# Local fields of the Ising model

Local fields  $\mathfrak{F}(z)$  of Ising

- ▶  $V \subset \mathbb{Z}^2$  finite subset
- ▶  $P: \{+1, -1\}^V \rightarrow \mathbb{C}$  a function

$$\text{parity} \begin{cases} P(-\sigma) = P(\sigma) & \text{even} \\ P(-\sigma) = -P(\sigma) & \text{odd} \end{cases}$$

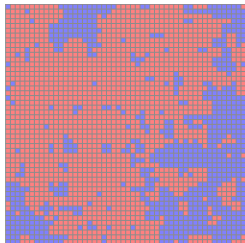
- ▶  $\mathfrak{F}(z) = P((\sigma_{z+\delta x})_{x \in V})$

- ▶  $\mathcal{F}$  space of local fields

$$\mathcal{F} = \mathcal{F}^+ \oplus \mathcal{F}^- \text{ by parity} \begin{cases} \mathcal{F}^+ & \text{even} \\ \mathcal{F}^- & \text{odd} \end{cases}$$

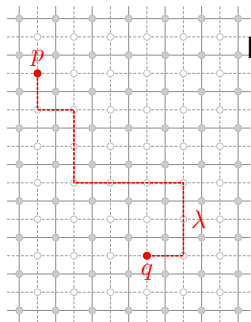
- ▶  $\mathcal{N} \subset \mathcal{F}$  space of null fields  
("zero in correlations")

$\sigma = (\sigma_z)_{z \in \Omega_\delta}$  Ising



$\mathcal{F}/\mathcal{N}$  equivalence classes of local fields (same correlations)

# Disorder operators in Ising model



Disorder operator pair:

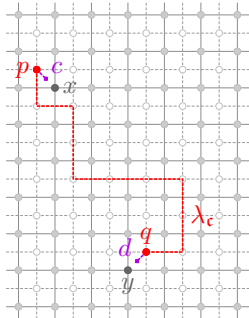
$$(\mu_p \mu_q)_\lambda := \exp \left( -2\beta \sum_{\langle z, w \rangle^* \in \lambda} \sigma_z \sigma_w \right)$$

- ▶  $p, q \in \mathbb{C}_\delta^*$  dual vertices
- ▶  $\lambda$  path between  $p$  and  $q$  on  $\mathbb{C}_\delta^*$  “disorder line”

*Remark:*

- ▶ a single disorder operator is **NOT** a local field
- ▶ a disorder operator pair is a local field (with fixed disorder line  $\lambda$ )

# Corner fermions in Ising model



- ▶  $c, d \in \mathbb{C}_\delta^c$  corners
- ▶  $x, y \in \mathbb{C}_\delta$  adjacent to  $c, d$ , respectively
- ▶  $p, q \in \mathbb{C}_\delta^*$  adjacent to  $c, d$ , respectively
- ▶  $\nu(c) := \frac{x-p}{|x-p|}$  phase factor
- ▶  $\lambda_c$  path between  $c$  and  $d$  "on  $\mathbb{C}_\delta^*$ "
- ▶  $\mathcal{W}(\lambda_c : c \rightsquigarrow d)$  cumulative angle of turning of  $\lambda_c$

Corner fermion pair:

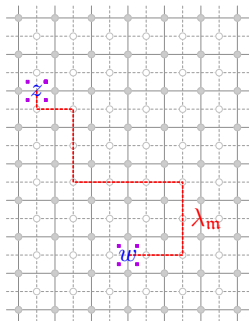
$$(\Psi_c^c \Psi_d^c)_{\lambda_c} := -\overline{\nu(c)} \exp\left(-\frac{i}{2} \mathcal{W}(\lambda_c : c \rightsquigarrow d)\right) (\mu_p \mu_q)_\lambda \sigma_x \sigma_y$$

*Remark:*

- ▶ one corner fermion is NOT a local field
- ▶ a corner fermion pair is a local field

(with fixed disorder line)

# Discrete holomorphic fermions in Ising model



- ▶  $z, w \in \mathbb{C}_\delta^m$  midpoints of edges
- ▶  $\lambda_m$  path between  $z$  and  $w$  “on  $\mathbb{C}_\delta^{**}$ ”
- ▶  $c, d \in \mathbb{C}_\delta^c$  adjacent to  $z, w$ , respectively
- ▶  $\lambda_c^{c,d}$  path between  $c$  and  $d$  on  $\mathbb{C}_\delta^*$  obtained by local modification of  $\lambda_m$

Holomorphic fermion pair:

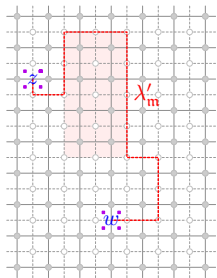
$$(\Psi(z)\Psi(w))_{\lambda_m} := \frac{1}{8\sqrt{2}} \sum_{c,d} (\Psi_c^c \Psi_d^c)_{\lambda_c^{c,d}}$$

*Remark: (as before)*

- ▶ one holomorphic fermion is NOT a local field
- ▶ a holomorphic fermion pair is a local field (with fixed disorder line)



# Properties of the fermion pairs



## Lemma (disorder line independence mod $\pm$ )

If  $\lambda_m, \lambda'_m$  are disorder lines between  $z, \zeta \in \mathbb{C}_\delta^m$  then

$$E\left[(\Psi(\zeta)\Psi(z))_{\lambda_m} \prod_{j=1}^n \sigma_{w_j}\right] = (-1)^{\mathcal{N}} \times E\left[(\Psi(\zeta)\Psi(z))_{\lambda'_m} \prod_{j=1}^n \sigma_{w_j}\right]$$

where  $\mathcal{N}$  is the number of points  $w_j$  in the area enclosed by  $\lambda_m$  and  $\lambda'_m$ .

## Lemma (antisymmetry of fermions)

$$(\Psi(\zeta)\Psi(z))_{\lambda_m} = -(\Psi(z)\Psi(\zeta))_{\lambda_m}$$

## Lemma (holomorphicity and singularity of fermion)

$$E\left[(\bar{\partial}_\delta \Psi(\zeta_\diamond)\Psi(z_m)) \prod_{j=1}^n \sigma_{w_j}\right] = \frac{-1}{4} \sum_{x \sim z_m} \delta_{\zeta_\diamond, x} \times E\left[\prod_{j=1}^n \sigma_{w_j}\right]$$

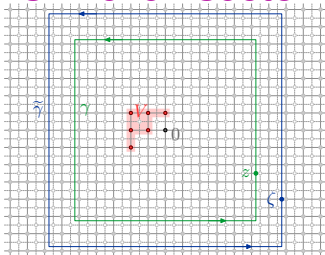
# Outline / steps

For the **Ising model** and ~~discrete GFF~~:

- ✓ 0.) Define model and local fields
- ✓ 1.) Suitable discrete contour integrals and residue calculus
- ✓ 2.) Introduce discrete holomorphic observable
  - 3.) **Define Laurent modes of the observable**
  - 4.) Commutation relations of Laurent modes
  - 5.) Virasoro action through Sugawara construction

# Laurent modes of fermions in even sector

- ▶  $\mathfrak{F}(w) = P[(\sigma_{w+x\delta})_{x \in V}]$   
even local field of Ising
- ▶  $\gamma, \tilde{\gamma}$  large nested counterclockwise closed paths on  $\mathbb{C}_\delta^c$



For  $k, \ell \in \mathbb{Z} + \frac{1}{2}$  define a new local field  $((\Psi_k \Psi_\ell) \mathfrak{F})(w)$  by

$$((\Psi_k \Psi_\ell) \mathfrak{F})(0) := \frac{1}{2\pi} \oint_{[\tilde{\gamma}]} \oint_{[\gamma]} \zeta_\diamond^{[k-\frac{1}{2}]} z_\diamond^{[\ell-\frac{1}{2}]} (\Psi(\zeta_m) \Psi(z_m)) \mathfrak{F}(0) [dz]_\delta [d\zeta]_\delta$$

## Lemma (discrete fermion mode pairs)

$(\Psi_k \Psi_\ell): \mathcal{F}^+ / \mathcal{N}^+ \rightarrow \mathcal{F}^+ / \mathcal{N}^+$  is well-defined

*Remark: (as before)*

- ▶ one fermion Laurent mode **is NOT defined**
- ▶ a fermion Laurent mode pair **is defined**, and acts on (even) local fields

# Outline / steps

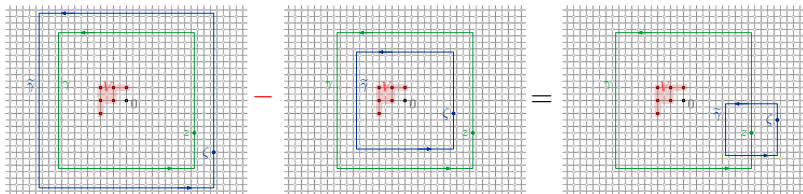
For the **Ising model** and ~~discrete GFF~~:

- ✓ 0.) Define model and local fields
- ✓ 1.) Suitable discrete contour integrals and residue calculus
- ✓ 2.) Introduce discrete holomorphic observable
- ✓ 3.) Define Laurent modes of the observable on *even* sector
  - 4.) **Anticommutation relations of Laurent modes**
  - 5.) Virasoro action through Sugawara construction

# Anticommutation of fermion modes in even sector

## Proposition (anticommutation of fermion modes)

$$(\Psi_k \Psi_\ell) + (\Psi_\ell \Psi_k) = \delta_{k+\ell,0} \text{id}_{\mathcal{F}^+/\mathcal{N}^+}$$



$$\begin{aligned} & \mathbb{E} \left[ \left( (\Psi_k \Psi_\ell) \mathfrak{F}(0) + (\Psi_\ell \Psi_k) \mathfrak{F}(0) \right) \cdots \right] \\ &= \delta_{k+\ell,0} \mathbb{E} \left[ \mathfrak{F}(0) \cdots \right] \end{aligned}$$

(residue calculus)

# Outline / steps

For the **Ising model** and ~~discrete GFF~~:

- ✓ 0.) Define model and local fields
- ✓ 1.) Suitable discrete contour integrals and residue calculus
- ✓ 2.) Introduce discrete holomorphic observable
- ✓ 3.) Define Laurent modes of the observable
- ✓ 4.) *Anti*commutation relations of Laurent modes
- 5.) **Virasoro action through Sugawara construction**

# Sugawara construction for Ising even local fields

- ▶  $V$  vector space

space  $\mathcal{F}^+/\mathcal{N}^+$  of even local fields modulo null fields

- ▶  $b_k: V \rightarrow V$  linear for each  $k \in \mathbb{Z} + \frac{1}{2}$

fermion Laurent mode pairs  $(\Psi_k \Psi_\ell): \mathcal{F}^+/\mathcal{N}^+ \rightarrow \mathcal{F}^+/\mathcal{N}^+$

$$\frac{1}{2\pi} \oint_{[\gamma]} \oint_{[\gamma]} \zeta_\diamond^{[k-\frac{1}{2}]} z_\diamond^{[\ell-\frac{1}{2}]} (\Psi(\zeta_m) \Psi(z_m)) (\cdots) [dz]_\delta [d\zeta]_\delta$$

- ▶  $\forall v \in V \exists N \in \mathbb{Z} : \ell \geq N \implies b_\ell v = 0$

monomial truncation:  $\forall z_\diamond \in \mathbb{C}_\diamond^\circ \exists D : \ell \geq D \implies z_\diamond^{[\ell-\frac{1}{2}]} = 0$

- ▶  $[b_k, b_\ell]_+ = \delta_{k+\ell,0} \text{id}_V$

anticommutation  $(\Psi_k \Psi_\ell) + (\Psi_\ell \Psi_k) = \delta_{k+\ell,0} \text{id}_{\mathcal{F}^+/\mathcal{N}^+}$

and  $(\Psi_p \Psi_k)(\Psi_\ell \Psi_q) + (\Psi_p \Psi_\ell)(\Psi_k \Psi_q) = \delta_{k+\ell,0} (\Psi_p \Psi_q)$

## Theorem (Virasoro action for Ising even sector)

$$\mathcal{L}_n := \frac{1}{2} \sum_{k>0} \left(\frac{1}{2} + k\right) (\Psi_{n-k} \Psi_k) - \frac{1}{2} \sum_{k<0} \left(\frac{1}{2} + k\right) (\Psi_k \Psi_{n-k})$$

defines Virasoro repr. with  $c = \frac{1}{2}$  on the space  $\mathcal{F}^+/\mathcal{N}^+$  of correlation equivalence classes of Ising even local fields.

# Outline / steps

For the **Ising model** and ~~discrete~~ GFF:

- ✓ 0.) Define model and local fields
- ✓ 1.) Suitable discrete contour integrals and residue calculus
- ✓ 2.) Introduce discrete holomorphic observable
- 3.) Define Laurent modes of the observable
- 4.) *Anti*commutation relations of Laurent modes
- 5.) Apply Sugawara construction to define Virasoro action on **odd** local fields



# Odd sector: Discrete half-integer monomials

## Proposition (discrete half-integer monomial functions)

$\exists$  functions  $z \mapsto z^{[\rho]}$ ,  $\rho \in \mathbb{Z} + \frac{1}{2}$ , defined on the double cover  $[\mathbb{C}_\delta^\diamond; 0] \cup [\mathbb{C}_\delta^m; 0]$  ramified at the origin, such that

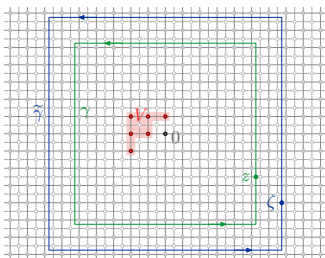
- ▶  $\bar{\partial}_\delta z^{[\rho]} = 0$  whenever ... “discrete holomorphicity”
  - ▶  $\rho > 0$  and  $z \in [\mathbb{C}_\delta^\diamond; 0] \cup [\mathbb{C}_\delta^m; 0]$
  - ▶  $\rho < 0$  and  $z \in [\mathbb{C}_\delta^\diamond; 0] \cup [\mathbb{C}_\delta^m; 0]$ ,  $\|z\|_1 > R_\rho \delta$
- ▶  $\partial_\delta z^{[\rho]} = \rho z^{[\rho-1]}$  “derivatives”
- ▶  $z^{[\rho]}$  has the same  $90^\circ$  rotation symmetry as  $z^\rho$  “symmetry”
- ▶ for  $\rho < 0$  we have  $z^{[\rho]} \rightarrow 0$  as  $\|z\| \rightarrow \infty$  “decay”
- ▶ for any  $z$  there exists  $D_z$  such that  $z^{[\rho]} = 0$  for  $\rho \geq D_z$  “truncation”

For  $\gamma$  large enough counterclockwise closed contour surrounding the origin...

- ▶  $\oint_{[\gamma]} z_m^{[\rho]} z_\diamond^{[q]} [dz]_\delta = 2\pi i \delta_{\rho+q, -1}$  “residue calculus”
- ▶  $\oint_{[\gamma]} z_m^{\{\rho\}} z_\diamond^{[q]} [dz]_\delta = 2\pi i \delta_{\rho+q, -1}$  where  $z_m^{\{\rho\}} = \frac{1}{4} \sum_{x \in \{\pm \frac{\delta}{2}, \pm i \frac{\delta}{2}\}} (z_m - x)^{[\rho]}$

# Odd sector: Laurent modes of fermions

- ▶  $\mathfrak{F}(w) = P[(\sigma_{w+x\delta})_{x \in V}]$   
odd local field of Ising
- ▶  $\gamma, \tilde{\gamma}$  large nested counterclockwise closed paths on  $\mathbb{C}_\delta^c$



For  $i, j \in \mathbb{Z}$  define a new local field  $((\Psi_i \Psi_j) \mathfrak{F})(w)$  by

$$((\Psi_i \Psi_j) \mathfrak{F})(0) := \frac{1}{2\pi} \oint_{[\tilde{\gamma}]} \oint_{[\gamma]} \zeta_\diamond^{[i-\frac{1}{2}]} z_\diamond^{[j-\frac{1}{2}]} (\Psi(\zeta_m) \Psi(z_m)) \mathfrak{F}(0) [dz]_\delta [d\zeta]_\delta$$

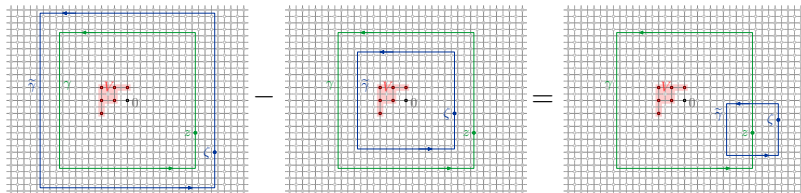
**Lemma (discrete fermion mode pairs)**

$(\Psi_i \Psi_j): \mathcal{F}^- / \mathcal{N}^- \rightarrow \mathcal{F}^- / \mathcal{N}^-$  is well-defined

# Odd sector: Anticommutation of fermion modes

Proposition (anticommutation of fermion modes)

$$(\Psi_i \Psi_j) + (\Psi_j \Psi_i) = \delta_{i+j,0} \text{id}_{\mathcal{F}^-/\mathcal{N}^-}$$



# Odd sector: Fermionic Sugawara construction

## Proposition (fermionic Sugawara, Ramond sector)

- Suppose:
- ▶  $V$  vector space,  $b_j: V \rightarrow V$  linear for each  $j \in \mathbb{Z}$
  - ▶  $\forall v \in V \exists N \in \mathbb{Z} : j \geq N \implies b_j v = 0$
  - ▶  $[b_i, b_j]_+ = \delta_{i+j,0} \text{id}_V$

Def.:

$$L_n := \frac{1}{2} \sum_{j \geq 0} \left(\frac{1}{2} + j\right) b_{n-j} b_j - \frac{1}{2} \sum_{j < 0} \left(\frac{1}{2} + j\right) b_j b_{n-j} \quad (n \in \mathbb{Z} \setminus \{0\})$$

$$L_0 := \frac{1}{2} \sum_{j > 0} j b_{-j} b_j + \frac{1}{16} \text{id}_V$$

Then:

- ▶  $L_n: V \rightarrow V$  is well defined
- ▶  $[L_n, L_m] = (n - m) L_{n+m} + \frac{n^3 - n}{24} \delta_{n+m,0} \text{id}_V$

## Theorem (Virasoro action on Ising odd local fields)

The space of odd Ising local fields modulo null fields becomes Virasoro representation with central charge  $c = \frac{1}{2}$ .

# Conclusions and outlook

✓ Lattice model fields of finite patterns form Virasoro repr.

- ▶ discrete Gaussian free field:  $\mathcal{L}_n$  on  $\mathcal{F}/\mathcal{N}$  by bosonic Sugawara
- ▶ Ising model:  $\mathcal{L}_n$  on  $\underbrace{\mathcal{F}^+/\mathcal{N}^+ \oplus \mathcal{F}^-/\mathcal{N}^-}_{\text{"Neveu-Schwarz } \oplus \text{ Ramond"}}$  by fermionic Sugawara

TODO Many CFT ideas rely on variants of Sugawara construction

- ▶ Wess-Zumino-Witten models
- ▶ symplectic fermions
- ▶ coset conformal field theories  $\rightsquigarrow$  CFT minimal models
- ▶ Coulomb gas formalism

TODO CFT fields  $\longleftrightarrow$  lattice model fields of finite patterns

- ▶ 1-1 correspondence via the Virasoro action on lattice model fields?
- ▶ correlations of lattice model fields with appropriate renormalization converge in scaling limit to CFT correlations?
- ▶ conceptual derivation of PDEs for limit correlations via singular vectors?

# THANK YOU!

