

closed-oriented

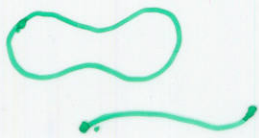
... so far



open-oriented

Homological Mirror Symmetry

$$\begin{array}{ccc}
 D^b(\text{Coh}M) & \leftrightarrow & \text{Fuk}(\tilde{M}) \\
 \vdots & & \vdots
 \end{array}$$



unoriented

Spaces with involutions

real geometry

\mathcal{C} with anti-involutions

Clifford algebra

⋮



heterotic

Spaces with G -bundles

"Flux"

Generalized Geometry

Open String



need

boundary conditions

||

D-branes

NLSM $\phi: \Sigma \rightarrow M$

- $N \subset M$ submanifold : b.c. $\phi: \partial\Sigma \rightarrow N$

- $(E, A) \rightarrow N$ hermitian vector bundle with connection :

path integral w.t. $e^{-S} \text{Pexp}\left(-i \int_{\partial\Sigma} \phi^* A\right)$ holonomy factor
 "Chan-Paton factor"

Supersymmetry :



boundary relates $\left\langle \begin{array}{l} \text{left} \\ \text{right} \end{array} \right.$ modes

$\rightarrow Q_{\pm}, \bar{Q}_{\pm}$ cannot independently conserved.

best cases : $\frac{1}{2}$ is preserved.

A-brane : $Q_A = \bar{Q}_+ + Q_-$, Q_A^{\dagger} conserved

B-brane : $Q_B = \bar{Q}_+ + \bar{Q}_-$, Q_B^{\dagger} conserved

mirror

NLSM on M Kähler $\left\{ \begin{array}{l} \omega \text{ symplectic str.} \\ J \text{ complex str.} \end{array} \right.$
 (N, E) is an

A-brane if $N \subset (M, \omega)$ Lagrangian
 E flat ($F_A = 0$)

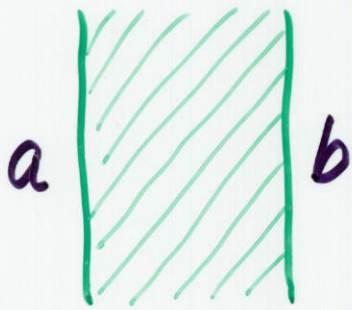
B-brane if $N \subset (M, J)$ complex submfd
 E holomorphic ($F_A \in \Omega^{2,0}$)

LG model $W: M \rightarrow \mathbb{C}$ superpotential

$$A \Rightarrow \text{Im } W|_N = \text{constant}$$

$$B \Rightarrow W|_N = \text{constant}$$

Open String States



Quantize on strip with

boundary conditions a and b
(both A or both B)

$\rightsquigarrow \mathcal{H}_{a,b} = \text{space of states}$
(∞ -dim)

On $\mathcal{H}_{a,b} \exists$ operators $Q, F = \begin{cases} Q_A, F_A & \text{A-branes} \\ Q_B, F_B & \text{B-branes} \end{cases}$
↑ optional

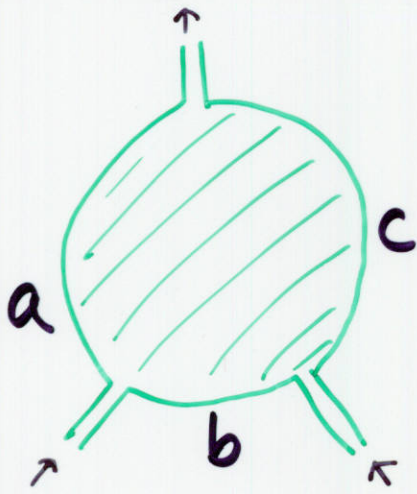
$$\{Q, Q^\dagger\} = 2H$$

$$Q^2 = 0 \quad (\text{in good cases: later})$$

$$[F, Q] = Q$$

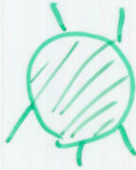
F values $\subset \mathbb{Z}$: $\rightarrow \mathcal{H}_{a,b}^q \xrightarrow{Q} \mathcal{H}_{a,b}^{q+1} \xrightarrow{Q} \dots$ Complex

$$\mathcal{H}_{a,b}^q(\text{sur}) \cong H^q(\mathcal{H}_{a,b}; \mathbb{Q})$$



$$\mathcal{H}_{a,b} \times \mathcal{H}_{b,c} \rightarrow \mathcal{H}_{a,c}$$

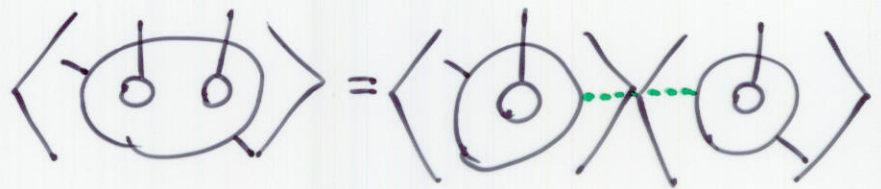
dg-Category \mathcal{D} $\begin{cases} \text{Obj: D-branes} \\ \text{mor: } \mathcal{H}_{a,b} \end{cases}$

- Coupling to topological gravity \Rightarrow  cyclic A_∞ -category

- $H(\mathcal{D}) \rightsquigarrow$ open/closed TFT (no top. grav.)

Moore-Seegal
Lazarovici

- $1 \in H(a,a)$
- $\langle a \circ b \rangle : H(a,b) \times H(b,a) \rightarrow \mathbb{C}$ non-degenerate
- pinching axiom



e.g.

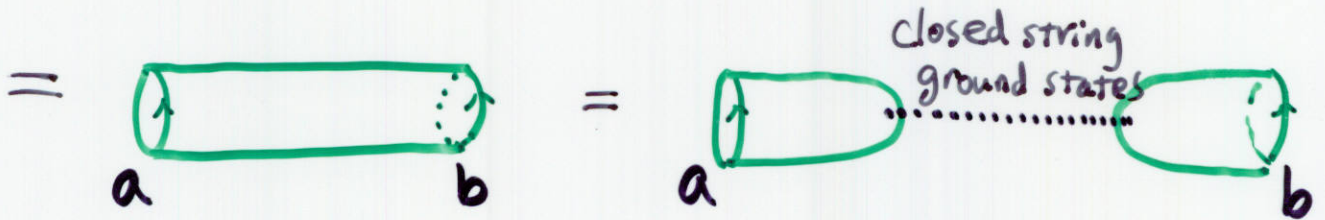


reconstructed by $\langle \text{pair of pants} \rangle$ $\langle \text{pair of pants} \rangle$ $\langle \text{pair of pants} \rangle$ obeying

associativity, Cardy, $\langle \text{pair of pants} \rangle = \langle \text{pair of pants} \rangle$, $\langle \text{pair of pants} \rangle = \langle \text{pair of pants} \rangle$

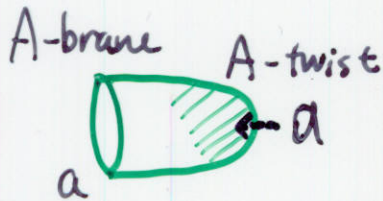
Witten index

$$\text{Tr}_{\mathcal{H}_{ab}} (-1)^F = \sum_p (-1)^p \dim H^p(a,b) = \chi(Q)$$



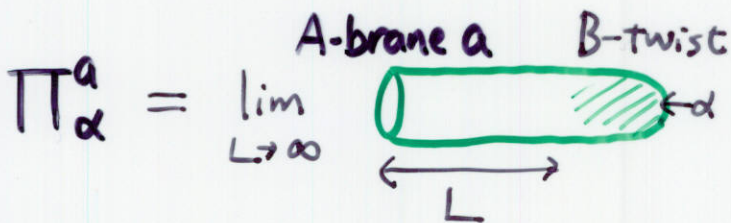
Riemann-Roch

Two possibilities



... topological disc amplitudes

- depend on M_f (Kähler)
- independent on M_c (complex str)



: Ramond-Ramond charge
of the brane **a**

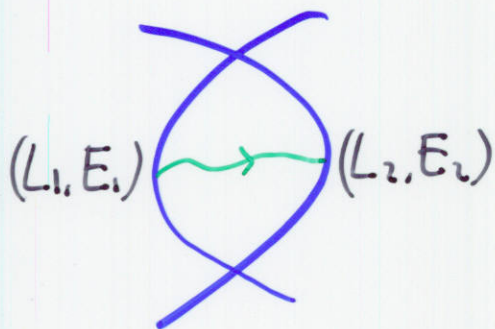
- independent on M_f (Kähler)
- dependence on M_c (complex)

$$\nabla_\alpha \Pi_\beta^a = \underbrace{C_{\alpha\beta}^r}_{R_B} \Pi_\gamma^a$$

R_B - structure constant.

A-branes in NLSM on M

$L \subset M$ Lagrangian submfld $E \rightarrow L$ flat



finite model for $\mathcal{H}_{1,2}, \mathcal{Q}$

$$\mathcal{H}_{1,2}^i = \bigoplus_{\substack{p \in L_1 \cap L_2 \\ \mu(p) = i}} \mathbb{C}[p] \otimes \text{Hom}(E_{1p}, E_{2p})$$

$$Q[p] = \sum_{\mu(p) = \mu(p) + 1} [q] \cdot \# \left(\begin{array}{c} \text{hol} \\ \text{disc} \end{array} \right) e^{-A} \underbrace{\Phi(E_1)^{-T} \Phi(E_2)}_{\text{holonomy}}$$

$Q^2 = 0$ may fail (disc bubble)

When $Q^2 = 0$, $H^i(\mathcal{H}_{1,2}) = HF((L_1, E_1), (L_2, E_2))$

Floer cohomology

(L, E) $\xrightarrow{\text{B-twist}}$ M $\xrightarrow{\text{rank } E} \int_L \Omega \cdot \mu$ $\mu \in H^0(M, \wedge^1 T^*M)$

$$\text{Tr}_{(L_1, E_1), (L_2, E_2)}^{(-1)^F} = \text{rank } E_1 \cdot \text{rank } E_2 \cdot \#(L_1 \cap L_2)$$

= : Riemann's bilinear identity

B-branes in $NL\sigma M$ M

$$N \subset M, \quad \mathcal{E} \xrightarrow{\text{hol}} N \quad \rightsquigarrow \quad \mathcal{E} \text{ coherent sheaf}$$

plx hol

$\mathcal{E}_1, \mathcal{E}_2$ vector bundles on M : 0-mode approx for $\mathcal{H}_{\mathcal{E}_1, \mathcal{E}_2}, Q$

$$\mathcal{H}_{\mathcal{E}_1, \mathcal{E}_2}^i = \Omega^{0,i}(M, \mathcal{E}_1^* \otimes \mathcal{E}_2)$$

$$Q = \bar{\partial}_{\mathcal{E}_1^* \otimes \mathcal{E}_2}$$

$$\begin{aligned} H^i(\mathcal{H}_{\mathcal{E}_1, \mathcal{E}_2}) &= H^{0,i}(M, \mathcal{E}_1^* \otimes \mathcal{E}_2) \\ &= \text{Ext}^i(\mathcal{E}_1, \mathcal{E}_2) \end{aligned}$$

$$\text{Tr}_{\mathcal{E}_1, \mathcal{E}_2} (-1)^F = \sum_p (-1)^p \text{Ext}^p(\mathcal{E}_1, \mathcal{E}_2) = \chi(\mathcal{E}_1, \mathcal{E}_2)$$

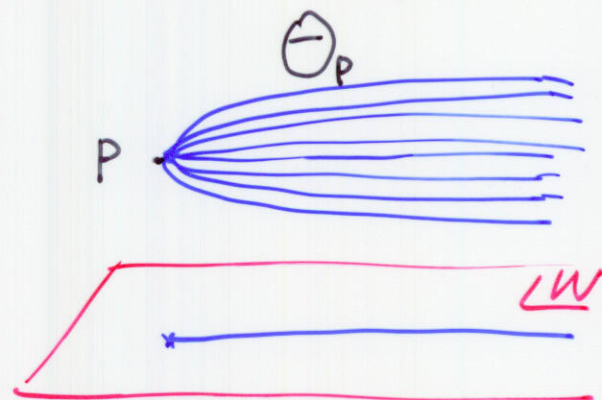
$$\begin{aligned} \mathcal{E} \quad \begin{array}{c} \text{A-twist} \\ \text{---} \end{array} &= \int_M \text{ch}(\mathcal{E}^*) \sqrt{\text{td}(X)} e^{B+i\omega} \mathcal{O} \\ \text{---} & \\ \mathcal{O} \in \text{Hob}(M) & \\ & + \dots \\ & \uparrow \\ & \mathcal{O}(e^{-\frac{1}{\hbar}}) \text{ correction} \end{aligned}$$

A-branes in LG (M, W)

Assume $\text{Crit}(W)$
all non-degenerate

branes: Lefschetz thimbles

$$P \in \text{Crit}(W) \quad \Theta_P = \bigcup \left(\text{gradient flow lines for } \text{Re} W \text{ starting } P \right)$$



Finite model for $\mathcal{H}_{\Theta_1, \Theta_2}, Q$

$$\mathcal{H}_{\Theta_1, \Theta_2} = \bigoplus \mathbb{C}[r]$$

$$\gamma: ([0, 1]; 0, 1) \rightarrow (M, \Theta_1, \Theta_2)$$

$$\dot{\gamma} = -\text{grad } \text{Im} W$$

$$Q[r] = \sum_{\substack{r' \\ \mu(r') = \mu(r) + 1}} [r'] \cdot \# \left(\left[\begin{array}{c} \text{hatched box} \\ \xrightarrow{\text{hol}} \\ \text{hatched box} \end{array} \right]_{\Theta_1, \Theta_2} \right) e^{-A(\text{Im} W)}$$

define Θ_P^\pm by $W_P \begin{cases} W(\Theta_P^+) \\ W(\Theta_P) \\ W(\Theta_P^-) \end{cases}$ small rotation

$$H(\mathcal{H}_{\Theta_1, \Theta_2}) = \begin{cases} \mathbb{C}^{\#(\Theta_1^-, \Theta_2^+)} \\ \mathbb{C} \end{cases} \leftarrow 0 \text{ if } \text{Im} W(\Theta_1) < \text{Im} W(\Theta_2) \text{ "directed"}$$

$$\Theta \int_{\text{Fun} M/W} e^{-iW} f \Omega = \int_{\Theta^-} e^{-iW} f \Omega$$

B-twist

B-branes in LG (M, W)

$$\sum_{\text{cptx}} \mathbb{C} \subset M \quad W|_{\mathbb{Z}} = \text{constant.}$$

Z_a, Z_b : $\mathcal{H}_{a,b}$ \mathbb{Z}_2 graded, $Q^2 = W(Z_a) - W(Z_b)$

If $W(Z_a) \neq W(Z_b)$, Q -fails to make a complex

If $W(Z_a) = W(Z_b)$, $Q^2 = 0$. 0-mode approx:

$$\mathcal{H}_{a,b} = \Omega^{0,*}(Z_a \cap Z_b, \wedge N_{Z_a \cap Z_b})$$

$$Q = \bar{\partial} + \partial W.$$

• $Z_a \cap Z_b = \emptyset \Rightarrow H^*(Q) = 0$

• $Z_a = Z_b = \{p\} \quad N_{\{p\}} = \mathbb{C}^n$

$$\mathcal{H}_{p,p} = \Omega^{0,*}(\{p\}, \wedge \mathbb{C}^n) = \wedge \mathbb{C}^n$$

$$0 \leftarrow \mathbb{C} \xleftarrow{\partial W(p)} \wedge^1 \mathbb{C}^n \xleftarrow{\partial W(p)} \wedge^2 \mathbb{C}^n \leftarrow \dots \leftarrow \wedge^n \mathbb{C}^n \leftarrow 0$$

$p \notin \text{Crit}(W)$: $\partial W(p) \neq 0 \Rightarrow H^*(Q) = 0$

$p \in \text{Crit}(W)$: $\partial W(p) = 0 \Rightarrow H^*(Q) = \wedge \mathbb{C}^n$

Alternative (more general) $M = \mathbb{C}^n$, $W = \text{Polynomial}$

..... Kontsevich

Objects : matrix factorizations of W

$$\mathcal{D} = \begin{pmatrix} 0 & f \\ g & 0 \end{pmatrix} \quad 2N \times 2N$$

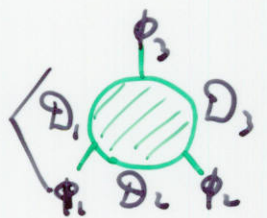
entries in $\mathbb{C}[X] = \mathbb{C}[X_1, \dots, X_n]$

$$\mathcal{D}^2 = W \cdot \mathbb{1}_{2N} \quad (\text{i.e. } fg = gf = W \mathbb{1}_N)$$


(or equivalently $M = \mathbb{C}[X]^N \oplus \mathbb{C}[X]^N$ \mathbb{Z}_2 -graded $\mathbb{C}[X]$ module
 \mathcal{D} : odd endomorphism of M , $\mathcal{D}^2 = W \mathbb{1}_M$)

Morphisms : $\mathcal{H}_{\mathcal{D}_1, \mathcal{D}_2} = \text{Hom}^{\text{ev}}(M_1, M_2) \oplus \text{Hom}^{\text{od}}(M_1, M_2)$

$$Q\phi = \mathcal{D}_2 \phi - (-1)^{|\phi|} \phi \mathcal{D}_1$$



$$\langle \text{Diagram} \rangle = \text{res}_W (\text{Str}(\phi_3 \phi_2 \phi_1 (d\mathcal{D}_1)^n))$$



$$\langle \text{Diagram} \rangle = \text{res}_W (f \text{Str}(\phi (d\mathcal{D})^n))$$

Kapustin - Li

B-branes in LG orbifold W/\mathbb{Z}_d $x_i \rightarrow e^{\frac{2\pi i}{d}} x_i$

$$W(x_1, \dots, x_n) \text{ degree } d \quad (W(e^{\frac{2\pi i}{d}} x) = e^{2id} W(x))$$

require $\mathcal{D}(x_1, \dots, x_n)$ also "homogeneous" H-Walcher

$$e^{i\alpha R} \mathcal{D}(e^{\frac{2\pi i}{d}} x) e^{-i\alpha R} = e^{i\alpha} \mathcal{D}(x)$$

$\rightarrow F_V$ conserved (necessary for conformal invariance)

$$\Rightarrow e^{i\pi R} \mathcal{D}(e^{\frac{2\pi i}{d}} x) e^{-i\pi R} = -\mathcal{D}(x)$$

$$\rho_\varphi(\omega) \mathcal{D}(\omega X) \rho_\varphi(\omega)^{-1} = \mathcal{D}(X) \quad \omega \in \mathbb{Z}_d \text{-invariant}$$

$$\rho_\varphi(\omega) = e^{\pi i \varphi} (-1)^F e^{i\pi R} \quad \varphi \text{ chosen so that } \rho_\varphi(\omega)^d = 1$$

$\varphi \equiv \varphi + 2$ labels \mathbb{Z}_d representations

\mathbb{Z}_d action on $\text{Hom}(M_1, M_2)$:

$$\phi(X) \rightarrow \rho_{\varphi_2}(\omega) \phi(\omega X) \rho_{\varphi_1}(\omega)^{-1}$$

$$\mathcal{H}_{(\mathcal{D}_1, \varphi_1), (\mathcal{D}_2, \varphi_2)}^{\text{orb}} = \text{Hom}(M_1, M_2)^{\mathbb{Z}_d}$$

$$U(1)_R \text{ sym: } \phi(X) \rightarrow e^{i\alpha R_2} \phi(e^{\frac{2\pi i}{d}} X) e^{-i\alpha R_1}$$

Charge can be made $\subset \mathbb{Z}$ by dressing $e^{i\alpha(\varphi_2 - \varphi_1)}$

\mathbb{Z} grading $(\varphi_i \neq \varphi_i + 2 \text{ for } \mathbb{Z}\text{-graded B-branes})$ Walcher

matrix factorizations of W \leftrightarrow Eisenbud Maximal Cohen Macaulay modules over $R = \mathbb{C}[X]/W$

1980's Auslander, Reiten
Knorrer, Buchweitz, ...

Buchweitz (1990's)

$$\underline{MF}(W) \cong \underline{MCM}(R) \cong \underline{D}^b(R)$$

\uparrow projectives = 0 \uparrow modulo perfect cplx's

Orlov (2005)

$$W = G(\phi_1, \dots, \phi_N) \text{ degree } d$$

$$M = \{W=0\} \subset \mathbb{C}P^{N-1} \text{ hypersurface}$$

$$d=N : \underline{grMF}(W) \xrightarrow{\Psi} D^b(\text{Coh } M)$$

\uparrow
 \mathbb{Z} graded B-branes in LG orbifold
 $W = G/\mathbb{Z}_d$



$$d > N : \underline{grMF}(W) \simeq \langle \underbrace{\bullet, \dots, \bullet}_{(d-N) \text{ objects}}, D^b(\text{Coh } M) \rangle$$

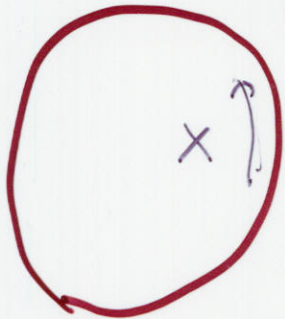
semi-orthogonal decomposition

$$d < N : D^b(\text{Coh } M) \simeq \langle \underbrace{\mathcal{O}_M(-N-d), \dots, \mathcal{O}_M}_{(N-d) \text{ objects}}, \underline{grMF}(W) \rangle$$

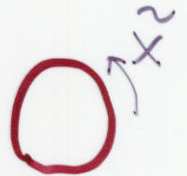
Mirror Symmetry for D-branes

T-duality

$$(S', R) \cong (\tilde{S}', \frac{\alpha'}{R})$$

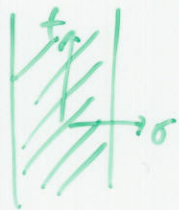


momentum \leftrightarrow winding #
winding # \leftrightarrow momentum



$$\Rightarrow \begin{aligned} \partial_t X &= \partial_\sigma \tilde{X} \\ \partial_\sigma X &= \partial_t \tilde{X} \end{aligned}$$

Open string:



Neumann b.c.

$$\partial_\sigma X = 0 \text{ on } \partial\Sigma$$

Dirichlet b.c.

$$\partial_t \tilde{X} = 0 \text{ on } \partial\Sigma$$

D1-brane wrapped on S' \leftrightarrow D0-brane at a point of \tilde{S}'

$U(1)$ holonomy $\oint_{S'} A \leftrightarrow$ position in \tilde{S}'

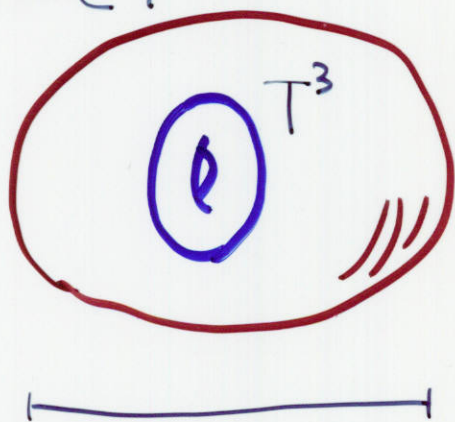
$$\tilde{S}' = \{ \text{D0-brane in T-dual} \} = \{ \text{D1-brane in } S' \} = H^1(S', U(1))$$

Space of $U(1)$ holonomy

$$\tilde{S}' = \text{dual space } H^1(S', U(1))$$

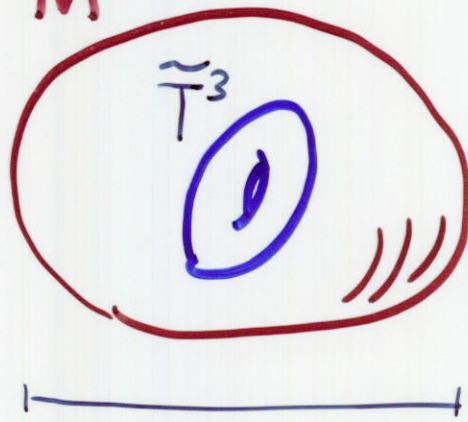
Strominger-Yau-Zaslow

$$M = CY^3$$



mirror
↔

$$\tilde{M}$$



$$\left\{ \begin{array}{l} \text{D3 brane wrapped on} \\ \text{Special Lagrangian } T^3 \subset M \end{array} \right\} = \tilde{M}$$

$$M = \left\{ \begin{array}{l} \text{D3 brane wrapped on} \\ \text{SLAG } \tilde{T}^3 \subset \tilde{M} \end{array} \right\}$$

Mirror Symmetry

$$= T\text{-duality} + \underline{\text{Quantum Correction}}$$

$\sim T^3$ degeneration

Gross

W.D. Puan

D. Morrison

;

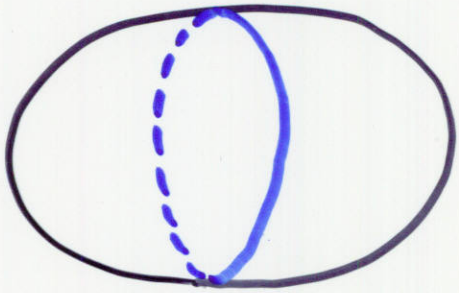
} topological construction
of the mirror

$$X = \mathbb{C}P^1$$

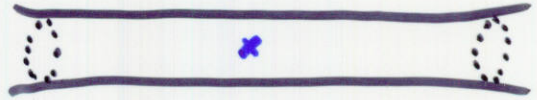


$$W = e^{-Y} + e^{-t+Y}$$

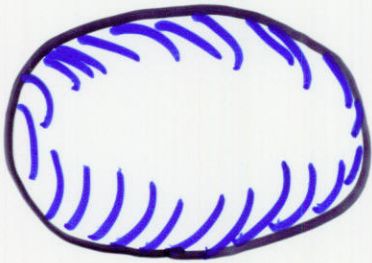
A-branes in $\mathbb{C}P^1$



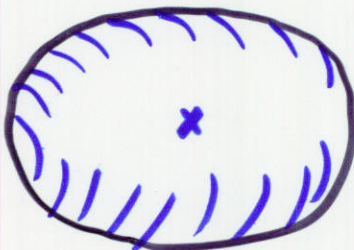
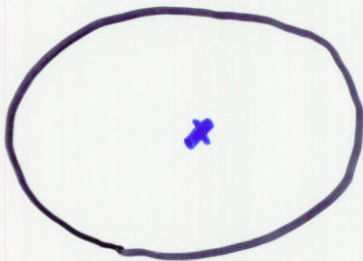
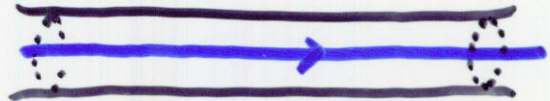
B-branes in LG



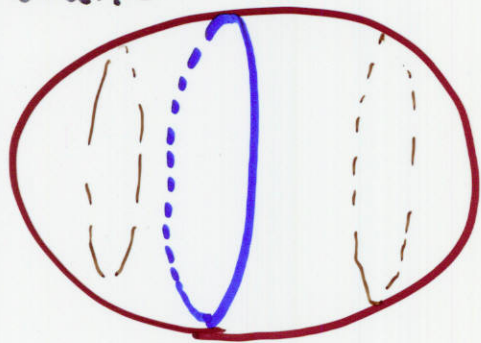
B-branes in $\mathbb{C}P^1$



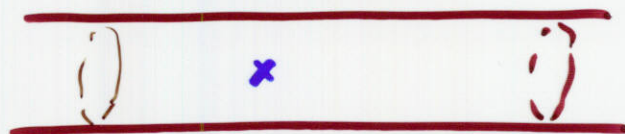
A-branes in LG



D1-brane



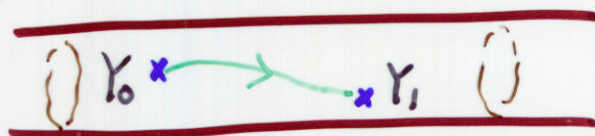
D0-brane



$$\text{Area}(\text{D1}) = \text{Re } Y_x$$

$$\int_{\text{D1}} A = \text{Im } Y_x$$

For two D0-branes in LG



$$Q_B^2 = W(Y_0) - W(Y_1) = e^{-Y_0} + e^{-t+Y_0} - e^{-Y_1} - e^{-t+Y_1}$$

$Q_B^2 = 0$ iff $Y_1 = Y_0$ or $Y_1 = t - Y_0$

$$\mathcal{H}_{Y_0, Y_1}(\text{surv}) = \begin{cases} \wedge \mathbb{C} = \mathbb{C} \otimes \mathbb{C} & \text{if } Y_0 = Y_1 \in \text{Crit}(W) \\ 0 & \text{otherwise} \end{cases}$$

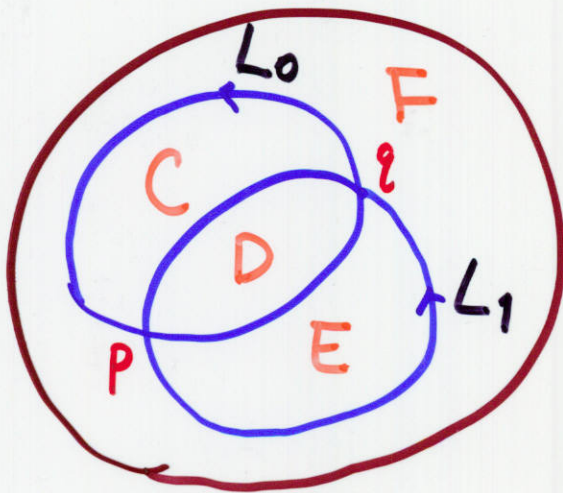
$$\text{Crit}(W) = \{ e^{-Y} = \pm e^{-Y/2} \} \rightarrow$$

$$\text{Area}(\text{D1}) = \frac{r}{2}$$

$$\int_{\text{D1}} A = 0 \text{ or } \pi$$

Computation in $NL\sigma M$

FOOO



$$C \cup D =: D_0$$

$$D \cup E =: D_1$$

$$Q_A[q] = e^{-A(D)}[p] - e^{-A(F)}[q]$$

$$Q_A[p] = e^{-A(C)}[q] - e^{-A(E)}[q]$$

$$Q_A^2[q] = \left[e^{-A(D_0)} - e^{-A(D_1)} - e^{-r+A(D_1)} + e^{-r+A(D_0)} \right][q]$$

$$\underline{Q_A^2 = 0 \text{ iff } A(D_1) = A(D_0) \text{ or } A(D_1) = r - A(D_0)}$$

If $A(D_1) = A(D_0)$ (i.e. $L_1 \sim L_0$)

$$Q_A[q] = e^{-A(D)} (1 - e^{-r+2A(D)})[p]$$

$$Q_A[p] = 0$$

$$\therefore HF(L_0, L_1) = \begin{cases} \sigma[q] \otimes C[p] & A(D_0) = A(D_1) = \frac{r}{2} \\ 0 & \text{otherwise.} \end{cases}$$

The prediction of mirror symmetry
was indeed correct.

The analysis is extended to general toric Fano.

Choi-Hyun Cho \leftarrow YG Oh

Also, compare

$$\rightarrow Q_A^2 = e^{-A(D_0)} + e^{-r+A(D_0)} - e^{-A(D_1)} - e^{-r+A(D_1)}$$

$$\rightarrow Q_B^2 = W|_{Y_0} - W|_{Y_1} = e^{-Y_0} + e^{-t+Y_0} - e^{-Y_1} - e^{-t+Y_1}$$

— identical! (as the should be)

$Q_B^2 = W|_{z_0} - W|_{z_1}$ has enough info to determine
 W (up to a shift)

\Rightarrow Computation of Q_A^2 in NLGM can reproduce
the LG superpotential of the mirror LG.

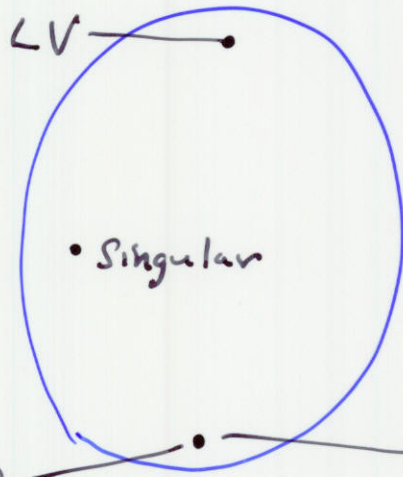
... This is one way to take into account
the quantum correction at degenerating
torus fibres.

MF mirror to real quintic

Kähler moduli
for quintic $\subset \mathbb{C}P^4$

$$\text{Fermat: } \sum_{i=1}^5 X_i^5 = 0$$

$$\frac{W = \sum_{i=1}^5 X_i^5}{\mathbb{Z}_5}$$



Complex structure moduli
space for mirror quintic

$$\sum_{i=1}^5 \tilde{X}_i^5 - 54 \tilde{X}_1 \cdots \tilde{X}_5 = 0$$

\mathbb{Z}_5

$$\frac{\tilde{W}_0 = \sum_{i=1}^5 \tilde{X}_i^5}{\mathbb{Z}_5^4}$$

RM \subset M real quintic $(\overset{\text{homeo}}{\approx} \mathbb{R}P^3)$

in $\tilde{\text{LGO}}$

$$\mathcal{D}_0 = \sum_{i=1}^5 (\tilde{X}_i^2 \eta_i + \tilde{X}_i^3 \bar{\eta}_i)$$

$$\eta_i = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}, \bar{\eta}_i = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

$$\text{on } \underbrace{(\mathbb{C}[\tilde{X}])^2 \hat{\otimes} \cdots \hat{\otimes} (\mathbb{C}[\tilde{X}])^2}_s$$

$$H^i(\mathcal{H}_{\mathcal{D}_0, \varphi_0, \mathcal{D}_0, \varphi_0}^{\text{orb}}) = \begin{cases} \mathbb{C} & 0 \\ \mathbb{C} & 1 \\ \mathbb{C} & 2 \\ \mathbb{C} & 3 \end{cases}$$

$$H^1: \Phi = (\eta_1 - \tilde{X}_1 \bar{\eta}_1) \cdots (\eta_5 - \tilde{X}_5 \bar{\eta}_5)$$

$$H^2: \tilde{X}_1 \cdots \tilde{X}_5 \cdot \mathbb{1} = \Phi^\dagger$$

They obey $\Phi^2 = -\Phi^T = -\tilde{X}_1 \dots \tilde{X}_5$

Deformation

$$W_0 \rightarrow W_\psi = W_0 - 54 \tilde{X}_1 \dots \tilde{X}_5$$

$$\mathcal{D}_0 \rightarrow \mathcal{D}_\psi = \mathcal{D}_0 + \psi \Phi$$

require $\mathcal{D}_\psi^2 = W_\psi$

$$\Rightarrow \boxed{\psi^2 = 54}$$

$\forall \psi \neq 0$, \exists two solutions $\psi = \pm \sqrt{54}$

$\longleftrightarrow \pi_1(\text{RM}) = \mathbb{Z}_2 \Rightarrow$ Two flat $U(1)$ bundles on RM

$$H^i(\mathcal{H}_{\mathcal{D}_\psi, \mathcal{D}_\psi}) = \begin{cases} \mathbb{C} \\ \mathbb{C} \\ \mathbb{C} \\ \mathbb{C} \end{cases} \text{ at } \psi=0, \quad \begin{cases} \mathbb{C} \\ 0 \\ 0 \\ \mathbb{C} \end{cases} \text{ at } \psi \neq 0$$

$\longleftrightarrow HF^i(\text{RM}, \text{RM}) = \begin{cases} \mathbb{C} \\ 0 \\ 0 \\ \mathbb{C} \end{cases}$ (near Large volume)

analytic continuation $\xrightarrow{\quad}$ $\begin{cases} \mathbb{C} \\ \mathbb{C} \\ \mathbb{C} \\ \mathbb{C} \end{cases}$ at $e^t = 0$