$C^{2,\alpha}$ REGULARITY OF FLAT FREE BOUNDARIES FOR THE THIN ONE-PHASE PROBLEM.

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ABSTRACT. We prove $C^{2,\alpha}$ regularity of sufficiently flat free boundaries, for the *thin* one-phase problem in which the free boundary occurs on a lower dimensional subspace. This problem appears also as a model of a one-phase free boundary problem in the context of the fractional Laplacian $(-\Delta)^{1/2}$.

1. INTRODUCTION

Let g(x,s) be a continuous non-negative function in the ball $B_1 \subset \mathbb{R}^{n+1} = \mathbb{R}^n \times \mathbb{R}$, which vanishes on a subset of $\mathbb{R}^n \times \{0\}$ and it is even in the *s* variable. We consider the following free boundary problem

(1.1)
$$\begin{cases} \Delta g = 0, & \text{in } B_1^+(g) := B_1 \setminus \{(x,0) : g(x,0) = 0\}, \\ \frac{\partial g}{\partial U} = 1, & \text{on } F(g) := \partial_{\mathbb{R}^n} \{x \in \mathcal{B}_1 : g(x,0) > 0\} \cap \mathcal{B}_1, \end{cases}$$

where

(1.2)
$$\frac{\partial g}{\partial U}(x_0) := \lim_{t \to 0^+} \frac{g(x_0 + t\nu(x_0), 0)}{\sqrt{t}}, \quad x_0 \in F(g)$$

with $\nu(x_0)$ the normal to F(g) at x_0 pointing toward $\{x : g(x, 0) > 0\}$ and $\mathcal{B}_r \subset \mathbb{R}^n$ the *n*-dimensional ball of radius r (centered at 0).

If F(g) is C^2 then it can be shown (see Section 7) that any function g which is harmonic in $B_1^+(g)$ has an asymptotic expansion at a point $x_0 \in F(g)$,

$$g(x,s) = \alpha(x_0)U((x-x_0) \cdot \nu(x_0), s) + o(|x-x_0|^{1/2} + s^{1/2}).$$

Here U(t,s) is the real part of \sqrt{z} which in the polar coordinates

$$t = r\cos\theta, \quad s = r\sin\theta, \quad r \ge 0, \quad -\pi \le \theta \le \pi,$$

is given by

(1.3)
$$U(t,s) = r^{1/2} \cos \frac{\theta}{2}$$

Then, the limit in (1.2) represents the coefficient $\alpha(x_0)$ in the expansion above (which justifies our notation)

$$\frac{\partial g}{\partial U}(x_0) = \alpha(x_0)$$

and our free boundary condition requires that $\alpha \equiv 1$ on F(g).

Solutions to our free boundary problem (1.1) are critical points to the energy functional

$$E(g) := \int_{B_1} |\nabla g|^2 \, dx \, ds + \frac{\pi}{2} \mathcal{H}^n(\{g > 0\} \cap \mathcal{B}_1).$$

If the second term is replaced by $\mathcal{H}^{n+1}(\{g > 0\})$, we obtain the classical one-phase free boundary problem (see for example [AC].) In our case the free boundary occurs on the lower dimensional subspace $\mathbb{R}^n \times \{0\}$ and for this reason we refer to (1.1) as to the *thin one-phase* free boundary problem.

This free boundary problem was first considered by Caffarelli, Roquejoffre and Sire [CafRS] as a model of a one-phase Bernoulli type free boundary problem in the context of the fractional Laplacian. It is relevant in applications when turbulence or long-range interactions are present, for example in flame propagation and also in the propagation of surfaces of discontinuities. For further information on this model see [CafRS] and the references therein.

In this paper we are interested in the question of regularity for the free boundary F(g). Concerning this issue the authors of [CafRS] proved that in dimension n = 2, a Lipschitz free boundary is C^1 . In [DR], the first author and Roquejoffre showed that in any dimension if the free boundary F(g) is sufficiently flat then it is $C^{1,\alpha}$.

This paper is the first of a series of papers, which investigate the regularity of F(g) and in particular the question of whether Lipschitz free boundaries are smooth. This basic question was answered positively in the case of minimal surfaces by De Giorgi [DG] and by Caffarelli [C1] for the standard one-phase free boundary problem.

Our strategy to obtain the regularity of Lipschitz free boundaries is to use a Weiss-type monotonicity formula [W] combined with flatness results and ad hoc Schauder type estimates near the free boundary. To implement this method we need to obtain first $C^{2,\alpha}$ estimates for flat free boundaries, which we achieve in this paper. Unlike the case of minimal surfaces and of the standard one-phase problem, $C^{2,\alpha}$ estimates do not seem to follow easily from $C^{1,\alpha}$. It appears that $C^{2,\alpha}$ is the critical regularity needed to obtain C^{∞} smoothness of the free boundary, as well as the regularity needed to implement our blow-up analysis.

The following is the main result of this paper (see Section 2 for the precise definition of viscosity solution to (1.1)).

Theorem 1.1. There exists $\bar{\epsilon} > 0$ small depending only on n, such that if g is a viscosity solution to (1.1) satisfying

 $\{x \in \mathcal{B}_1 : x_n \le -\bar{\epsilon}\} \subset \{x \in \mathcal{B}_1 : g(x,0) = 0\} \subset \{x \in \mathcal{B}_1 : x_n \le \bar{\epsilon}\},\$

then F(g) is a $C^{2,\alpha}$ graph in $\mathcal{B}_{\frac{1}{2}}$ for every $\alpha \in (0,1)$ with $C^{2,\alpha}$ norm bounded by a constant depending on α and n.

The proof of Theorem 1.1 follows the lines of the flatness theorem in [DR], which is inspired by the regularity theory developed by the second author in [S]. In this case the proof is more technical since we need to approximate the free boundary quadratically. To do so, we introduce a family of approximate solutions $V_{S,a,b}$ which have the same role as quadratic polynomials in the regularity theory of elliptic equations. Such family will be used also in a subsequent paper to obtain boundary Schauder type estimates for solutions to our problem.

In the last section of this paper we also prove some useful general facts about viscosity solutions g to our free boundary problem (1.1), such as $C^{1/2}$ -optimal regularity, asymptotic expansion near regular points of the free boundary and compactness.

The paper is organized as follows. In Section 2 we recall notation, definitions and some basic results from [DR], including the linearized problem associated to (1.1). Section 3 is devoted to the construction of the quadratic approximate solutions $V_{S,a,b}$. In Section 4 we prove a Harnack type inequality for solutions to (1.1). In Section 5 we establish the improvement of flatness result via a compactness argument which makes crucial use of the Harnack inequality of Section 4. Our argument reduces the problem to studying the regularity of solutions to the linearized problem. This is pursued in Section 6. We conclude the paper with Section 7 where we provide some general facts about viscosity solutions to (1.1).

2. Definitions and basic lemmas

In this section we recall notation, definitions and some necessary results from [DR].

2.1. **Basic facts.** Throughout the paper, constants which depend only on the dimension n will be called universal. In general, small constants will be denoted by c, c_i and large constants by C, C_i and they may change from line to line in the body of the proofs.

A point $X \in \mathbb{R}^{n+1}$ will be denoted by $X = (x, s) \in \mathbb{R}^n \times \mathbb{R}$, and sometimes $x = (x', x_n)$ with $x' = (x_1, \ldots, x_{n-1})$.

A ball in \mathbb{R}^{n+1} with radius r and center X is denoted by $B_r(X)$ and for simplicity $B_r = B_r(0)$. Also \mathcal{B}_r denotes the *n*-dimensional ball $B_r \cap \{s = 0\}$.

Let $v \in C(B_1)$ be a non-negative function. We associate to v the following sets:

$$B_1^+(v) := B_1 \setminus \{(x,0) : v(x,0) = 0\} \subset \mathbb{R}^{n+1}$$
$$B_1^+(v) := B_1^+(v) \cap \mathcal{B}_1 \subset \mathbb{R}^n;$$
$$F(v) := \partial_{\mathbb{R}^n} \mathcal{B}_1^+(v) \cap \mathcal{B}_1 \subset \mathbb{R}^n.$$

Often subsets of \mathbb{R}^n are embedded in \mathbb{R}^{n+1} , as it will be clear from the context.

We consider the thin one-phase free boundary problem

(2.1)
$$\begin{cases} \Delta g = 0, & \text{in } B_1^+(g), \\ \frac{\partial g}{\partial U} = 1, & \text{on } F(g), \end{cases}$$

where

$$\frac{\partial g}{\partial U}(x_0) := \lim_{t \to 0^+} \frac{g(x_0 + t\nu(x_0), 0)}{\sqrt{t}}, \quad X_0 = (x_0, 0) \in F(g).$$

Here $\nu(x_0)$ denotes the unit normal to F(g), the free boundary of g, at x_0 pointing toward $\mathcal{B}_1^+(g)$.

We now recall the notion of viscosity solutions to (2.1), introduced in [DR].

Definition 2.1. Given g, v continuous, we say that v touches g by below (resp. above) at $X_0 \in B_1$ if $g(X_0) = v(X_0)$, and

 $g(X) \ge v(X)$ (resp. $g(X) \le v(X)$) in a neighborhood O of X_0 .

If this inequality is strict in $O \setminus \{X_0\}$, we say that v touches g strictly by below (resp. above).

Definition 2.2. We say that $v \in C(B_1)$ is a (strict) comparison subsolution to (2.1) if v is a non-negative function in B_1 which is even with respect to s = 0 and it satisfies

(i) v is C^2 and $\Delta v \ge 0$ in $B_1^+(v)$;

(ii)
$$F(v)$$
 is C^2 and if $x_0 \in F(v)$ we have

$$v(x_0 + t\nu(x_0), 0) = \alpha(x_0)\sqrt{t} + o(\sqrt{t}), \text{ as } t \to 0^+,$$

with

$$\alpha(x_0) \ge 1,$$

where $\nu(x_0)$ denotes the unit normal at x_0 to F(v) pointing toward $\mathcal{B}_1^+(v)$;

(iii) Either v is not harmonic in $B_1^+(v)$ or $\alpha(x_0) > 1$ at all $x_0 \in F(v)$.

Similarly one can define a (strict) comparison supersolution.

Definition 2.3. We say that g is a viscosity solution to (2.1) if g is a continuous non-negative function in B_1 which is even with respect to s = 0 and it satisfies

- (i) $\Delta g = 0$ in $B_1^+(g)$;
- (ii) Any (strict) comparison subsolution (resp. supersolution) cannot touch g by below (resp. by above) at a point $X_0 = (x_0, 0) \in F(g)$.

Remark 2.4. We remark that if g is a viscosity solution to (2.1) in B_{λ} , then

$$g_{\lambda}(X) = \lambda^{-1/2} g(\lambda X), \quad X \in B_1$$

is a viscosity solution to (2.1) in B_1 .

Finally, we state for completeness the boundary Harnack inequality which will be often used throughout the paper. This version follows from the boundary Harnack inequality proved in [CFMS].

Theorem 2.5 (Boundary Harnack Inequality). Let v be harmonic in $B_1^+(v)$ and let F(v) be a Lipschitz graph in the e_n -direction (pointing towards the positive phase) with $0 \in F(v)$. If w is harmonic in $B_1^+(w) = B_1^+(v)$, then

$$\frac{w}{v} \le C\frac{w}{v}(\frac{1}{2}e_n) \quad in \ B_{3/4},$$

with C depending only on n and on the Lipschitz constant of F(v).

2.2. The function \tilde{g} . Here and henceforth we denote by P the half-hyperplane

$$P := \{ X \in \mathbb{R}^{n+1} : x_n \le 0, s = 0 \}$$

and by

$$L := \{ X \in \mathbb{R}^{n+1} : x_n = 0, s = 0 \}.$$

Also, throughout the paper we call $U(X) := U(x_n, s)$, where U is the function defined in (1.3).

Let g be a continuous non-negative function in \overline{B}_{ρ} . As in [DR], we define the multivalued map \tilde{g} which associate to each $X \in \mathbb{R}^{n+1} \setminus P$ the set $\tilde{g}(X) \subset \mathbb{R}$ via the formula

(2.2)
$$U(X) = g(X - we_n), \quad \forall w \in \tilde{g}(X)$$

We write $\tilde{g}(X)$ to denote any of the values in this set.

This change of variables has the same role as the partial Hodograph transform for the standard one-phase problem. Our free boundary problem becomes a problem

with fixed boundary for \tilde{g} , and the limiting values of \tilde{g} on L give the free boundary of g as a graph in the e_n direction.

Recall that if g satisfies the ϵ -flatness assumption

(2.3)
$$U(X - \epsilon e_n) \le g(X) \le U(X + \epsilon e_n) \quad \text{in } B_{\rho}, \text{ for } \epsilon > 0$$

then $\tilde{g}(X) \neq \emptyset$ for $X \in B_{\rho-\epsilon} \setminus P$ and $|\tilde{g}(X)| \leq \epsilon$, hence we can associate to g a possibly multi-valued function \tilde{g} defined at least on $B_{\rho-\epsilon} \setminus P$ and taking values in $[-\epsilon, \epsilon]$ which satisfies

(2.4)
$$U(X) = g(X - \tilde{g}(X)e_n).$$

Moreover if g is strictly monotone in the e_n -direction in $B^+_{\rho}(g)$, then \tilde{g} is single-valued.

We recall the following lemmas from [DR].

Lemma 2.6. Let g, v be non-negative continuous functions in B_{λ} with v strictly increasing in the e_n -direction in $B_{\lambda}^+(v)$. Assume that g and v satisfy the flatness assumption (2.3) in B_{λ} for $\epsilon > 0$ small. If

$$v \leq g \quad in \ B_{\lambda},$$

then

$$\tilde{v} \leq \tilde{g} \quad on \ B_{\lambda-\epsilon} \setminus P_{\epsilon}$$

Viceversa, if

$$\tilde{v} \leq \tilde{g} \quad on \ B_{\sigma} \setminus P,$$

for some $0 < \sigma < \lambda - \epsilon$, then

 $v \leq g$ on $B_{\sigma-\epsilon}$.

Lemma 2.7. Let g, v be respectively a solution and a subsolution to (2.1) in B_2 , with v strictly increasing in the e_n -direction in $B_2^+(v)$. Assume that g and v satisfy the flatness assumption (2.3) in B_2 for $\epsilon > 0$ small. If,

(2.5)
$$\tilde{v} + \sigma \leq \tilde{g} \quad in \ (B_{3/2} \setminus B_{1/2}) \setminus P_{3/2}$$

for some $\sigma > 0$, then

(2.6)
$$\tilde{v} + \sigma \leq \tilde{g} \quad in \ B_{3/2} \setminus P_{\tau}$$

Finally, given a Lipschitz function ϕ defined on $B_{\lambda}(\bar{X})$, with values in [-1,1], then for all $\epsilon > 0$ small there exists a unique function φ_{ϵ} defined at least on $B_{\lambda-\epsilon}(\bar{X})$ such that

(2.7)
$$U(X) = \varphi_{\epsilon}(X - \epsilon \phi(X)e_n), \quad X \in B_{\lambda}(X),$$

that is

$$\tilde{\varphi_{\epsilon}} = \epsilon \phi$$

Moreover such function φ_{ϵ} is increasing in the e_n -direction.

If g satisfies the flatness assumption (2.3) in B_1 and ϕ is as above then (say $\lambda < 1/4, \bar{X} \in B_{1/2}$,)

(2.8)
$$\tilde{\varphi}_{\epsilon} \leq \tilde{g} \quad \text{in } B_{\lambda}(\bar{X}) \setminus P \Rightarrow \varphi_{\epsilon} \leq g \quad \text{in } B_{\lambda-\epsilon}(\bar{X})$$

The following Proposition will be used in the compactness argument for the proof of the improvement of flatness in Section 6. **Proposition 2.8.** Let ϕ be a smooth function in $B_{\lambda}(\bar{X}) \subset \mathbb{R}^{n+1} \setminus P$. Define (for $\epsilon > 0$ small) the function φ_{ϵ} as above by

(2.9)
$$U(X) = \varphi_{\epsilon}(X - \epsilon \phi(X)e_n).$$

Then,

(2.10)
$$\Delta \varphi_{\epsilon} = \epsilon \Delta (U_n \phi) + O(\epsilon^2), \quad in \ B_{\lambda/2}(\bar{X})$$

with the function in $O(\epsilon^2)$ depending on $\|\phi\|_{C^5}$ and λ .

Proof. For notional simplicity we drop the subindex ϵ in the definition of φ_{ϵ} . From formula (2.9) and Taylor's theorem, we have that

(2.11)
$$U(X) = \varphi(X) - \epsilon \varphi_n(X) \phi(X) + \epsilon^2 \Psi(X), \quad \text{in } B_{\lambda/2}(\bar{X})$$

with $\|\Psi\|_{C^3(B_{\lambda/2}(\bar{X}))} \leq C$ and C depending on $\|\phi\|_{C^5}$ and λ . Thus,

$$U_n(X) = \varphi_n(X) + O(\epsilon).$$

Combining this formula for $\varphi_n(X)$ and (2.11) we obtain

$$U(X) = \varphi(X) - \epsilon U_n(X)\phi(X) + O(\epsilon^2).$$

Hence, using that U is harmonic,

$$0 = \Delta U(X) = \Delta \varphi(X) - \Delta(\epsilon U_n \phi)(X) + O(\epsilon^2),$$

as desired.

We remark that in fact the function in $O(\epsilon^2)$ only depends on λ if we choose ϵ small enough depending on $\|\phi\|_{C^5}$.

2.3. The linearized problem. We recall here the linearized problem associated to (2.1). Here and later U_n denotes the x_n -derivative of the function U. Recall that

$$P := \{ X \in \mathbb{R}^{n+1} : x_n \le 0, s = 0 \}, \quad L := \{ X \in \mathbb{R}^{n+1} : x_n = 0, s = 0 \}.$$

Given $h \in C(B_1)$ and $X_0 = (x'_0, 0, 0) \in B_1 \cap L$, we call

$$|\nabla_r h|(X_0) := \lim_{(x_n, s) \to (0, 0)} \frac{h(x'_0, x_n, s) - h(x'_0, 0, 0)}{r}, \quad r^2 = x_n^2 + s^2.$$

Once the change of unknowns (2.2) has been done, the linearized problem associated to (2.1) is

(2.12)
$$\begin{cases} \Delta(U_n h) = 0, & \text{in } B_1 \setminus P, \\ |\nabla_r h| = 0, & \text{on } B_1 \cap L. \end{cases}$$

Definition 2.9. We say that h is a solution to (2.12) if $h \in C(B_1)$, h is even with respect to $\{s = 0\}$ and it satisfies

- (i) $\Delta(U_n h) = 0$ in $B_1 \setminus P$;
- (ii) *h* cannot be touched by below (resp. by above) at any $X_0 = (x'_0, 0, 0) \in B_1 \cap L$, by a continuous function ϕ which satisfy

$$\phi(X) = \phi(X_0) + a(X_0) \cdot (x' - x'_0) + b(X_0)r + O(|x' - x'_0|^2 + r^{3/2}),$$

with $b(X_0) > 0$ (resp. $b(x_0) < 0$).

In Section 6, we will prove a quadratic expansion for solutions to the linearized problem which yields the following corollary.

$$\Box$$

Corollary 2.10. Let h be a solution to (2.12) such that $|h| \leq 1$. Given any $\alpha \in (0, 1)$, there exists η_0 depending on α , such that h satisfies

$$|h(X) - (h(0) + \xi_0 \cdot x' + \frac{1}{2}(x')^T M_0 x' - \frac{a_0}{2}r^2 - b_0 r x_n)| \le \frac{1}{4}\eta_0^{2+\alpha} \quad in \ B_{\eta_0},$$

with $r^2 = x_n^2 + s^2$, for some $a_0, b_0 \in \mathbb{R}, \xi_0 \in \mathbb{R}^{n-1}, M_0 \in S^{(n-1)\times(n-1)}$ with

 $|\xi_0|, |a_0|, |b_0|, ||M_0|| \le C, \quad C \text{ universal}$

and

$$a_0 + b_0 - tr M_0 = 0.$$

3. A family of functions.

In this section we introduce a family of functions $V_{S,a,b}$ which approximate our solution quadratically. These functions will be often used as comparison subsolutions/supersolutions. We establish here some of their basic properties, including their behavior under the change of coordinates $V \to \tilde{V}$ (see Proposition 3.5).

We start by presenting some basic properties of the solution U defined in the introduction. Recall that

$$U(t,s) := \rho^{1/2} \cos \frac{\beta}{2},$$

where

$$t = \rho \cos \beta, \quad s = \rho \sin \beta, \quad \rho \ge 0, \quad -\pi \le \beta \le \pi.$$

We will use the following properties of the function U:

(i) $\Delta U = 0$, U > 0 in $\mathbb{R}^{n+1} \setminus P$.

(ii)
$$U_t = \frac{1}{2}\rho^{-1/2}\cos\frac{\beta}{2} = \frac{1}{2\rho}U$$
 and $U_t > 0$ in $\mathbb{R}^{n+1} \setminus P$.

Since U_t is positive harmonic in $\mathbb{R}^2 \setminus \{(t,0), t \leq 0\}$, homogenous of degree -1/2and vanishes continuously on $\{(t,0), t < 0\}$ one can see from boundary Harnack inequality (or by direct computation) that values of U_t at nearby points with the same second coordinate are comparable in diadic rings. Precisely we have

(3.1)
$$\frac{U_t(t_1,s)}{U_t(t_2,s)} \le C \quad \text{if} \quad |t_1 - t_2| \le \frac{1}{2} |(t_2,s)|$$

Next we introduce the family $V_{S,a,b}$. For any $a, b \in \mathbb{R}$ we define the following family of (two-dimensional) functions (given in polar coordinates (ρ, β))

(3.2)
$$v_{a,b}(t,s) := (1 + \frac{a}{4}\rho + \frac{b}{2}t)\rho^{1/2}\cos\frac{\beta}{2},$$

that is

$$v_{a,b}(t,s) = (1 + \frac{a}{4}\rho + \frac{b}{2}t)U(t,s) = U(t,s) + o(\rho^{1/2}).$$

Given a surface $S = \{x_n = h(x')\} \subset \mathbb{R}^n$, we call $\mathcal{P}_{S,X}$ the 2D plane passing through X = (x, s) and perpendicular to S, that is the plane containing X and generated by the s-direction and the normal direction from (x, 0) to S.

We define the family of functions

(3.3)
$$V_{\mathcal{S},a,b}(X) := v_{a,b}(t,s), \quad X = (x,s),$$

with $t = \rho \cos \beta$, $s = \rho \sin \beta$ respectively the first and second coordinate of X in the plane $\mathcal{P}_{\mathcal{S},X}$. In other words, t is the signed distance from x to \mathcal{S} (positive above \mathcal{S} in the x_n -direction.)

If

$$\mathcal{S} := \{ x_n = \frac{1}{2} (x')^T M x' + \xi' \cdot x' \},\$$

for some $M \in S^{(n-1)\times(n-1)}, \xi' \in \mathbb{R}^{n-1}$ we use the notation

(3.4)
$$V_{M,\xi',a,b}(X) := V_{\mathcal{S},a,b}(X).$$

This will be the case throughout most of the paper.

Definition 3.1. For $\delta > 0$ small, we define the following classes of functions

$$\mathcal{V}_{\delta} := \{ V_{M,\xi',a,b} : \|M\|, |\xi'|, |a|, |b| \le \delta \},\$$

and

$$\mathcal{V}^0_{\delta} := \{ V_{M,\xi',a,b} \in \mathcal{V}_{\delta} : a+b-trM = 0 \}.$$

Most of the times we will work with functions in the class \mathcal{V}_{δ} , since we deal with the flat case. Notice that if we rescale $V = V_{M,\xi',a,b}$ that is

$$V_{\lambda}(X) = \lambda^{-1/2} V(\lambda X), \quad X \in B_1,$$

then it easily follows from our definition that

$$V_{\lambda} = V_{\lambda M, \xi', \lambda a, \lambda b}.$$

In the next proposition we provide a condition for a function $V \in \mathcal{V}_{\delta}$ to be a subsolution/supersolution.

Proposition 3.2. Let $V = V_{M,\xi',a,b} \in \mathcal{V}_{\delta}$, with $\delta \leq \delta_0$ universal. There exists a universal constant $C_0 > 0$ such that if

$$(3.5) a+b-trM \ge C_0 \delta^2$$

then V is a comparison subsolution to (2.1) in B_2 .

Proof. Clearly from our formula for $v_{a,b}$ the function V satisfies the free boundary condition of Definition 2.2 with $\alpha(x_0) \equiv 1$. We need to check that $\Delta V(X) > 0$ at all $X \in B_2^+(V)$.

Since that V(X) depends only on (t, s) and

$$\Delta_x t = -\kappa(x)$$

where $\kappa(x)$ is the sum of the principal curvatures of the parallel surface to S (in \mathbb{R}^n) passing through x, we compute that

(3.6)
$$\Delta V(X) = \Delta_{(t,s)} v_{a,b} - (\partial_t v_{a,b}) \kappa(x).$$

From our formula for $v_{a,b}$, using polar coordinates we get that

(3.7)
$$\Delta_{(t,s)}v_{a,b} = \frac{1}{2}(a+b)\rho^{-1/2}\cos\frac{\beta}{2} = (a+b)U_t.$$

Also, since $\rho \leq 2$,

(3.8)
$$|\partial_t v_{a,b} - U_t| \le (|a| + |b|)\rho^{1/2} \cos\frac{\beta}{2} \le 8\delta U_t$$

Finally we use that $\kappa_i(x)$ the principal curvatures at x are given by,

(3.9)
$$\kappa_i(x) = \frac{\kappa_i(x^*)}{1 - t\kappa_i(x^*)}$$

where x^* is the projection of x onto S. Since $|\xi'|, ||M|| \leq \delta$ we obtain that

 $|\kappa_i(x^*)| \le C\delta, \quad |\kappa(x^*) - trM| \le C\delta^3$

for C universal, which in view of (3.9) give

$$(3.10) |\kappa(x) - trM| \le C\delta^2.$$

From (3.6) combined with (3.7), (3.8) and (3.10) we get that

(3.11)
$$|\Delta V(X) - (a+b-trM)U_t| \le \frac{1}{2}C_0\delta^2 U_t$$

for a C_0 universal. It follows that if

$$a+b-trM \ge C_0\delta^2$$

then $\Delta V(X) > 0$ as desired.

Next, we estimate V_n and ΔV outside a small cone with axis L.

Proposition 3.3. Let
$$V = V_{M,\xi',a,b} \in \mathcal{V}_{\delta}$$
 with $\delta \leq \delta_0$ universal, then
(3.12) $c \leq \frac{V_n}{U_n} \leq C$, in $B_2 \setminus (P \cup \{|(x_n,s)| \leq 10\delta |x'|\})$.

If
$$V \in \mathcal{V}^0_{\delta}$$
 then

$$(3.13) \qquad |\Delta V(X)| \le C\delta^2 U_n(X) \quad in \ B_2 \setminus (P \cup \{|(x_n, s)| \le 10\delta |x'|\}).$$

Proof. From our formula

$$V_n(X) = \partial_t v_{a,b}(t,s) \frac{\partial t}{\partial x_n}$$

where t represents the signed distance from x to S. Since $\nabla_x t$ is the unit vector at x that has the direction of the normal from x to S, it makes an angle of order δ with respect to e_n . Hence since

$$\frac{\partial t}{\partial x_n} = \nabla_x t \cdot e_n$$

we get

(3.14)
$$1 \ge \frac{\partial t}{\partial x_n} \ge 1 - C\delta^2$$

and we obtain

$$\partial_t v_{a,b}(t,s) \ge V_n(X) \ge \frac{1}{2} \partial_t v_{a,b}(t,s).$$

From (3.8) we see that $\partial_t v_{a,b} \sim U_t$ and we obtain that

(3.15)
$$2\partial_t U(t,s) \ge V_n(X) \ge \frac{1}{4}\partial_t U(t,s)$$

Thus to obtain our claim we need to replace t with x_n in the inequality above.

Since in $B_{2|x|}$ the surface S is in a $4\delta|x|$ neighborhood of $x_n = 0$ we find that $|t - x_n| \le 4\delta|x|$. If X belongs to the domain in (3.12) then

$$|(x_n, s)| \ge 8\delta |x| \ge 2|t - x_n|$$

and we obtain from (3.1)

(3.16)
$$c \le \frac{U_t(t,s)}{U_t(x_n,s)} \le C$$

which together with (3.15) gives the desired conclusion (3.12).

Now (3.13) follows immediately. Indeed by formula (3.11) we have that

$$|\Delta V(X)| \le C\delta^2 U_t(t,s),$$

which combined with (3.16) gives the desired bound.

Remark 3.4. We remark that if $V \in \mathcal{V}^0_{\delta}$, then the rescaling $V_{\lambda}(X) = \lambda^{-1/2} V(\lambda X)$ with $\lambda \leq 1$, satisfies

(3.17)
$$c \le \frac{(V_{\lambda})_n}{U_n} \le C, \quad |\Delta V_{\lambda}(X)| \le C\delta^2 U_n(X),$$

in the dilation ball of factor $1/\lambda$

$$B_{2/\lambda} \setminus (P \cup \{ |(x_n, s)| \le 10\delta |x'| \}).$$

Indeed

$$\Delta V_{\lambda}(X) = \lambda^{3/2} \Delta V(\lambda X), \quad U_n(X) = \lambda^{1/2} U_n(\lambda X), \quad (V_{\lambda})_n(X) = \lambda^{1/2} V_n(\lambda X).$$

Now we study the behavior of $V \in \mathcal{V}_{\delta}$ under the transformation $V \to \tilde{V}$. This will be quite useful in the rest of the paper.

Proposition 3.5. Let $V = V_{M,\xi',a,b} \in \mathcal{V}_{\delta}$, with $\delta \leq \delta_0$ universal. Then V is strictly monotone increasing in the e_n -direction in $B_2^+(V)$. Moreover, \tilde{V} satisfies the following estimate in $B_2 \setminus P$

$$|\tilde{V}(X) - \gamma_V(X)| \le C_1 \delta^2, \quad \gamma_V(X) = \frac{a}{2}r^2 + brx_n - \frac{1}{2}(x')^T M x' - \xi' \cdot x'$$

with $r = \sqrt{x_n^2 + s^2}$ and C_1 a universal constant.

Proof. First we show that $v_{a,b}$ satisfies

(3.18)
$$U(t + \gamma_{a,b} - C\delta^2, s) \le v_{a,b}(t,s) \le U(t + \gamma_{a,b} + C\delta^2, s)$$

where $\rho^2 = t^2 + s^2$ and $\gamma_{a,b}$ is the following expression depending on t and s:

$$\gamma_{a,b}(t,s) := \frac{a}{2}\rho^2 + b\rho t.$$

Indeed since (see properties of U listed at the beginning of this action)

$$|U_{tt}| \le C\rho^{-1}U_t$$

we have that if $|\mu| \leq \rho/2$ then

$$|U(t+\mu,s) - (U(t,s) + \mu U_t(t,s))| \le \mu^2 |U_{tt}(t',s)| \le C\mu^2 \rho^{-1} U_t(t,s),$$

where in the last inequality we used (3.1). Thus, since $U_t = U/(2\rho)$,

$$(1 + \frac{\mu}{2\rho} + C\frac{\mu^2}{\rho^2})U(t,s) \ge U(t+\mu,s) \ge (1 + \frac{\mu}{2\rho} - C\frac{\mu^2}{\rho^2})U(t,s).$$

Choosing

$$\mu = \tilde{\mu} \pm 4C \frac{\tilde{\mu}^2}{\rho}$$

we obtain that

$$U(t+\tilde{\mu}+4C\frac{\tilde{\mu}^2}{\rho},s) \ge (1+\frac{\tilde{\mu}}{2\rho})U(t,s) \ge U(t+\tilde{\mu}-4C\frac{\tilde{\mu}^2}{\rho},s),$$

provided that $|\tilde{\mu}/\rho| < c$, with c sufficiently small. Since

$$v_{a,b} = \left(1 + \frac{a}{4}\rho + \frac{b}{2}t\right)U$$

we can apply the inequality above with

$$\tilde{\mu} = \frac{a}{2}\rho^2 + bt\rho,$$

hence $|\tilde{\mu}|/\rho \leq C\delta$ and obtain the claim.

When t is the signed distance from x to the surface S we have

$$t = 0$$
 on $S := \{x_n = h(x') := \frac{1}{2}x'^T M x' + \xi' \cdot x'\}$

and by (3.14)

$$1 \ge \frac{\partial t}{\partial x_n} \ge 1 - C\delta^2 \quad \text{in } B_1$$

thus, by integrating this inequality on the segment $(x', h(x')), (x', x_n)$ we get

$$|t - (x_n - h(x'))| \le C\delta^2$$

Since in B_1 , the surface S and $x_n = 0$ are within distance δ from each other we have $|t - x_n| \leq C\delta$ and hence

$$|\gamma_{a,b}(t,s) - \gamma_{a,b}(x_n,s)| \le \|\nabla v_{a,b}\|_{L^{\infty}} |t - x_n| \le C\delta^2.$$

From the last two inequalities we have that

$$|(t + \gamma_{a,b}(t,s)) - (x_n + \gamma_V(X))| \le C\delta^2,$$

with

$$\gamma_V(X) = \gamma_{a,b}(x_n, s) - \frac{1}{2}x'^T M x' - \xi' \cdot x'.$$

Using this fact and (3.18) (and the monotonicity of U in the e_n direction) we obtain

$$U(X + (\gamma_V(X) - C\delta^2)e_n) \le V(X) \le U(X + (\gamma_V(X) + C\delta^2)e_n),$$

and the estimate for \tilde{V} is proved.

Finally, we remark that the monotonicity of V follows from (3.15).

Remark 3.6. Notice that from the last inequality in the proof above, we obtain that if $V \in \mathcal{V}_{\delta}$, then V satisfies the 4δ -flatness assumption in B_1 (see also (2.3)):

$$U(X - 4\delta e_n) \le V(X) \le U(X + 4\delta e_n).$$

This could be also checked easily directly from the definition of V.

We conclude this section with by comparing the functions V corresponding to two nearby surfaces.

Lemma 3.7. Let S_i , i = 1, 2 be surfaces with curvature bounded by 2. Let

$$V_i = V_{\mathcal{S}_i, a_i, b_i}, \quad |a_i|, |b_i| \le 2, \quad i = 1, 2.$$

Assume that,

$$S_i \cap B_{2\sigma} = \{x_n = h_i(x')\}, \quad \sigma \le c$$

with h_i Lipschitz graphs, $h_i(0) = 0, |\nabla h_i| \le 1$ and c universal. If
 $|a_1 - a_2|, |b_1 - b_2| \le \epsilon, ||h_1 - h_2||_{L^{\infty}} \le \epsilon \sigma^2,$

for some small $\epsilon \leq c$, then

$$V_1(X) \le V_2(X + C\epsilon\sigma^2 e_n) \quad in \ B_{\sigma}$$

Proof. After a rescaling of factor $1/\sigma$, we need to prove our lemma for $\sigma = 1$ and with the curvature of S_i , a_i , b_i and ϵ smaller than c universal.

First we prove that for $0 < \lambda \leq 1$,

$$v_{a_1,b_1}(t,s) \le v_{a_2,b_2}(t+C\epsilon\lambda^2,s), \quad \lambda \le \rho = |(t,s)| \le 2\lambda.$$

By (3.15), $\partial_t v_{a,b}$ is proportional to $\partial_t U$ in the disk of radius 2. Since on the segment with endpoints (t, s) and $(t + C\epsilon\lambda^2, s)$ all the values of $\partial_t U$ are comparable (see (3.1)) we obtain (using $2\rho U_t = U$)

$$\begin{aligned} v_{a_2,b_2}(t+C\epsilon\lambda^2,s) &\geq v_{a_2,b_2}(t,s) + C\epsilon\lambda^2 U_t(t,s) \\ &\geq U(1+\frac{a_2}{4}\rho + \frac{b_2}{2}t + C\epsilon\frac{\lambda^2}{\rho}) \\ &\geq U(1+\frac{a_1}{4}\rho + \frac{b_1}{2}t) \\ &\geq v_{a_1,b_1}(t,s), \end{aligned}$$

and our claim is proved.

Since v_{a_2,b_2} is increasing in the first coordinate, we obtain that

 $v_{a_1,b_1}(t,s) \le v_{a_2,b_2}(t+C\epsilon,s), \quad |(t,s)| \le 1.$

On the other hand, from the hypotheses on h_i we see that in B_1

$$t_1 + C\epsilon \le \bar{t}_2,$$

where \bar{t}_2 is the distance to $S_2 - C' \epsilon e_n$, for some C' large depending on the C above. Hence in B_1 we have

$$V_1(X) = v_{a_1,b_1}(t_1,s) \le v_{a_2,b_2}(t_1 + C\epsilon,s) \le v_{a_2,b_2}(\bar{t}_2,s) = V_2(X + C'\epsilon e_n).$$

4. HARNACK INEQUALITY

In this section we state and prove a Harnack type inequality for solutions to our free boundary problem (2.1). This will allow us to obtain some compactness of flat solutions after the transformation $g \to \tilde{g}$ (see Corollary 4.2) which is a crucial ingredient in Theorem 1.1.

Theorem 4.1 (Harnack inequality). There exist $\bar{\epsilon} > 0$ small and $\bar{C} > 0$ large universal, such that if g solves (2.1) and it satisfies

(4.1)
$$V(X + a_0 e_n) \le g(X) \le V(X + b_0 e_n) \quad in \ B_\rho(X^*) \subset B_1,$$

with $V = V_{M,\xi',a,b} \in \mathcal{V}^0_{\delta}$, and
 $\overline{z} = 2, \qquad b_0 - a_0$

$$\bar{C}\delta^2 \le \frac{b_0 - a_0}{\rho} \le \bar{\epsilon},$$

with $|a_0|, |b_0| \le 1$, then

(4.2)
$$V(X + a_1 e_n) \le g(X) \le V(X + b_1 e_n)$$
 in $B_{\bar{\eta}\rho}(X^*)$,

with

$$a_0 \leq a_1 \leq b_1 \leq b_0, \quad b_1 - a_1 = (1 - \bar{\eta})(b_0 - a_0),$$

for a small universal constant $\bar{\eta} \in (0, 1/2)$.

In the particular case when V = U, this statement was proved in [DR]. Our proof follows the same lines as the one in [DR] but it requires a more careful analysis since the function V is no longer a precise solution.

From this statement we get the desired corollary to be used in the proof of our main result. Precisely, assume g satisfies (4.1) in B_1 with $a_0 = -\epsilon$, $b_0 = \epsilon$ for some small $\epsilon \ll \bar{\epsilon}$, and δ such that $\bar{C}\delta^2 \leq \epsilon$. Notice that from Remark 3.6, the functions V and g are $(4\delta + \epsilon)$ -flat in B_1 .

Then at any point $X^* \in B_{1/2}$ we can apply Harnack inequality repeatedly for a sequence of radii $\rho_m = \frac{1}{2} \bar{\eta}^m$ and obtain

$$V(X + a_m e_n) \le g(X) \le V(X + b_m e_n) \quad \text{in } B_{\frac{1}{2}\bar{\eta}^m}(X^*),$$

with

(4.3)
$$b_m - a_m = (b_0 - a_0)(1 - \bar{\eta})^m = 2\epsilon (1 - \bar{\eta})^m,$$

for all m's, $m \ge 1$ such that

(4.4)
$$4\epsilon \frac{(1-\bar{\eta})^{m-1}}{\bar{\eta}^{m-1}} \le \bar{\epsilon}$$

This implies that for all such m's, the function \tilde{g} satisfies

(4.5)
$$\hat{V} + a_m \leq \tilde{g} \leq \tilde{V} + b_m, \quad \text{in } B_{\frac{1}{2}\bar{\eta}^m - 4\delta - \epsilon}(X^*) \setminus P_{\underline{q}}$$

with a_m, b_m as in (4.3). Define the following (possibly multivalued) function

(4.6)
$$\tilde{g}_{\epsilon,V}(X) := \frac{\tilde{g}(X) - V(X)}{\epsilon}, \quad X \in B_{1-4\delta-\epsilon} \setminus P,$$

and notice that

provided that

(4.8)
$$4\delta + \epsilon \le \epsilon^{1/2} \le \bar{\eta}^m/4$$

If $\epsilon \leq \bar{\epsilon} \bar{\eta}^{2m_0}$ for some nonnegative integer m_0 then our inequalities above (4.4), (4.8) and hence also (4.5) hold for all $m \leq m_0$. We thus obtain the following corollary.

Corollary 4.2. Let g solve (2.1) and satisfy for $\epsilon \leq \overline{\epsilon}$ V

$$V(X - \epsilon e_n) \le g(X) \le V(X + \epsilon e_n)$$
 in B_1 ,

with

$$V = V_{M,\xi',a,b} \in \mathcal{V}^0_{\delta}, \quad \bar{C}\delta^2 \le \epsilon,$$

for $\bar{\epsilon}, \bar{C} > 0$ universal constants. If

$$\epsilon \leq \bar{\epsilon} \ \bar{\eta}^{2m_0},$$

for some nonnegative integer m_0 (with $\bar{\eta} > 0$ small universal), then the function $\tilde{g}_{\epsilon,V}$ defined in (4.6) satisfies

 $a_{\epsilon}(X) \leq \tilde{g}_{\epsilon,V}(X) \leq b_{\epsilon}(X), \quad in \ B_{1/2} \setminus P$

with

$$b_{\epsilon} - a_{\epsilon} \le 2(1 - \bar{\eta})^{m_0},$$

and $a_{\epsilon}, b_{\epsilon}$ having a modulus of continuity bounded by the Hölder function αt^{β} for α, β depending only on $\bar{\eta}$.

The proof of Harnack inequality will follow from the Proposition below.

Proposition 4.3. There exist $\bar{\epsilon}, \bar{\delta} > 0$ and $\bar{C} > 0$ universal, such that if g solves (2.1) and it satisfies

(4.9)
$$V(X - \epsilon e_n) \le g(X) \le V(X + \epsilon e_n) \quad in \ B_1, \quad for \ 0 < \epsilon \le \overline{\epsilon}$$

with

$$V = V_{M,\xi',a,b} \in \mathcal{V}^0_{\delta}, \quad \delta \le \bar{\delta}, \quad \bar{C}\delta^2 \le \epsilon,$$

then either

$$g(X) \le V(X + (1 - \eta)\epsilon e_n)$$
 in B_{η} ,

or

$$g \ge V(X - (1 - \eta)\epsilon e_n)$$
 in B_{η} ,

for a small universal constant $\eta \in (0, 1)$.

First we show that if $g \ge V$ and they separate of order ϵ at one point, then they separate also of order ϵ away from a neighborhood of $L = \{x_n = 0, s = 0\}$. This follows from the boundary Harnack inequality. Below are the details.

Lemma 4.4. If g solves (2.1) and it satisfies

(4.10)
$$g(X) \ge V(X - \epsilon e_n) \quad in \ B_1,$$

(4.11)
$$g(\bar{X}) \ge V(\bar{X}) \quad at \ some \ \bar{X} \in B_{\frac{1}{8}}(\frac{1}{4}e_n),$$

with $V = V_{M,\xi',a,b} \in \mathcal{V}^0_{\delta}, \ \bar{C}\delta^2 \leq \epsilon \ for \ \bar{C} > 0 \ universal, \ then$

(4.12)
$$g(X) \ge V(X - (1 - \tau)\epsilon e_n) \quad in \mathcal{C},$$

with

$$\mathcal{C} := \{ (x', x_n, s) : \frac{d}{2} \le |(x_n, s)| \le \frac{1}{2}, |x'| \le \frac{1}{2} \}, \quad d = \frac{1}{8\sqrt{n-1}}$$

and τ a small universal constant $\tau \in (0, 1)$.

Proof. We have

$$V(X - (1 - \tau)\epsilon e_n) = V(X - \epsilon e_n) + \tau \epsilon V_n(X + \lambda \epsilon e_n)$$

for some λ with $|\lambda| < 1$. Hence by (3.1),(3.12) for ϵ small enough and $X \in \mathcal{C}$

$$V(X - (1 - \tau)\epsilon e_n) \le V(X - \epsilon e_n) + C\tau\epsilon U_n(X + \lambda\epsilon e_n)$$
$$\le V(X - \epsilon e_n) + C_1\tau\epsilon U_n(X).$$

Thus, if $h(X) := g(X) - V(X - \epsilon e_n)$ we need to show that

(4.13)

and then choose $\tau = c_1/C_1$.

To obtain (4.13), notice that by a similar computation as the one above in view of (4.11) and (3.1),(3.12) we get that for ϵ small enough

 $h \ge c_1 \epsilon U_n, \quad \text{in } \mathcal{C},$

(4.14)
$$h(\bar{X}) \ge V(\bar{X}) - V(\bar{X} - \epsilon e_n) \ge cU_n(\bar{X})\epsilon \ge c_2\epsilon.$$

Also, by (4.10) we have

$$h \ge 0$$
 in B_1 .

Finally, by (3.13)

$$|\Delta h| \le C\delta^2 U_n \le C_2 \delta^2, \quad \text{in } \widetilde{\mathcal{C}} \setminus P.$$

where $\widetilde{\mathcal{C}} \supset \mathcal{C}$ is the d/4-neighborhood of \mathcal{C} .

Thus in view of (4.14) and Harnack inequality we have that (for \overline{C} large enough)

(4.15)
$$h \ge c_2 \epsilon - C\delta^2 \ge c_3 \epsilon, \quad \text{in } B_{1/8}(\frac{1}{4}e_n).$$

Denote by

$$D := \widetilde{\mathcal{C}} \setminus (B_{1/8}(\frac{1}{4}e_n) \cup P)$$

and let q_1, q_2 satisfy in D

$$(4.16) \qquad \qquad \Delta q_1 = 0, \quad \Delta q_2 = -1$$

with boundary conditions respectively

$$q_1 = 0$$
 on $\partial \widetilde{\mathcal{C}} \cup P$, $q_1 = 1$ on $\partial B_{1/8}(\frac{1}{4}e_n)$

and

$$q_2 = 0 \quad \text{on } \partial D.$$

By boundary Harnack inequality, q_1 is comparable to the distance function s in a neighborhood of $P \cap C \subset \subset \widetilde{C}$. Since q_2 is Lipschitz continuous in a neighborhood of $P \cap C$, we then obtain

(4.17)
$$q_1 \ge c_4 q_2 \quad \text{in } \mathcal{C} \setminus B_{1/8}(\frac{1}{4}e_n),$$

with $c_4 > 0$ universal. By the maximum principle,

$$h \ge q := c_3 \epsilon \, q_1 - C_2 \delta^2 \, q_2 \quad \text{in } D,$$

since $h \ge q$ on ∂D and $\Delta h \le \Delta q$ in D. Hence, by (4.17) we get that (for \overline{C} large enough)

$$h \ge \epsilon \frac{c_3}{2} q_1 \ge c_5 \epsilon U_n, \quad \text{in } \mathcal{C} \setminus B_{1/8}(\frac{1}{4}e_n),$$

where in the last inequality we used that (by boundary Harnack inequality) q_1 and U_n are comparable. This inequality together with (4.15) gives the desired claim (4.13).

We are now ready to present the proof of Proposition 4.3.

Proof of Proposition 4.3. Assume that

(4.18)
$$g(\bar{X}) - V(\bar{X}) \ge 0, \quad \bar{X} = \frac{1}{2}e_n.$$

Then in view of assumption (4.9) from Lemma 4.4, after the change of variables $g \to \tilde{g}$ we get that

(4.19)
$$\tilde{g}(X) \ge \tilde{V}(X) + \tau \epsilon - \epsilon \quad \text{in } \mathcal{C}' \setminus P$$

with

$$\mathcal{C}' := \{ (x', x_n, s) : d \le |(x_n, s)| \le \frac{1}{4}, |x'| \le \frac{1}{2} \}, \quad d = \frac{1}{8\sqrt{n-1}}.$$

Denote by

$$W(X) := V_{M + \frac{c}{n-1}\epsilon I, \xi', a, b+2c\epsilon}(X) \in \mathcal{V}_{\delta + \epsilon},$$

with c small to be made precise later. Then in view of Proposition 3.5 we have

(4.20)
$$-2C_1(\delta+\epsilon)^2 \le (\tilde{V}-\tilde{W}) + c\epsilon(2rx_n - \frac{1}{2(n-1)}|x'|^2) \le 2C_1(\delta+\epsilon)^2.$$

First we choose c small depending on τ such that

$$\tilde{V} \ge \tilde{W} - \frac{\tau}{2}\epsilon,$$

where we used that $\bar{C}\delta^2 \leq \epsilon \leq \bar{\epsilon}$ with $\bar{C} \geq C(\tau)$ and $\bar{\epsilon}$ small enough. Then, if \bar{C} is sufficiently large depending on c,

(4.21)
$$\tilde{V} \ge \tilde{W} + \tau^* \epsilon$$
, on $\{|(x_n, s)| \le d, |x'| = \frac{1}{2}\} \setminus P$,

for some $\tau^* > 0$ small, say $\tau^* < \tau/2$. These combined with (4.19) give

(4.22)
$$\widetilde{g}(X) \ge \widetilde{W}(X) + \tau^* \epsilon - \epsilon \quad \text{in } (\mathcal{C}' \cup \{|(x_n, s)| \le d, |x'| = \frac{1}{2}\}) \setminus P.$$

Moreover, if \overline{C} is large enough we get that W satisfies (3.5) and hence W is a subsolution. Thus from Lemma 2.7 and the inequality above we conclude that

(4.23)
$$\widetilde{g}(X) \ge \widetilde{W}(X) + \tau^* \epsilon - \epsilon \quad \text{in } \{ |(x_n, s)| \le d, |x'| \le \frac{1}{2} \} \setminus P.$$

Finally, from (4.20) we see that there is a small neighborhood around the origin $B_{\eta} \subset \{ |(x_n, s)| \leq d, |x'| \leq \frac{1}{2} \}$ (η small universal depending on the constants above, $\eta < \tau^*/2$) such that

$$\tilde{W} \ge \tilde{V} - \frac{\tau^*}{2}\epsilon, \quad \text{in } B_{2\eta} \setminus P.$$

Hence, from (4.23) we conclude that

$$\tilde{g} \ge \tilde{V} + \eta \epsilon - \epsilon$$
 in $B_{2\eta} \setminus P$,

for some small universal constant η , and the lemma is proved after the change of variable $\tilde{g} \to g$.

We conclude this section with the proof of Theorem 4.1.

Proof of Theorem 4.1. After a translation of the origin we may assume that we satisfy our flatness hypothesis (4.1) in $B_{\rho}(X^*) \subset B_2$ with

$$(x^*)' = 0, \quad a_0 + b_0 = 0, \quad V \in \mathcal{V}^0_{2\delta}$$

We dilate the picture by a factor of $2/\rho$ and work with the rescalings

$$g_{\rho}(X) = (\frac{\rho}{2})^{-1/2} g(\frac{\rho}{2}X), \quad V_{\rho}(X) = (\frac{\rho}{2})^{-1/2} V(\frac{\rho}{2}X).$$

which are defined in a ball of radius 2 included in $B_{4/\rho}$. Notice that, if $V \in \mathcal{V}_{2\delta}^0$ then $V_{\rho} \in \mathcal{V}_{2\delta}^0$.

After dropping the subindex ρ for simplicity of notation, we may assume that the flatness condition (4.1) holds in some ball $B_2(X^*) \subset \mathbb{R}^{n+1}$, with $V \in \mathcal{V}^0_{2\delta}$,

$$a_0 = -\epsilon, \quad b_0 = \epsilon, \quad (x^*)' = 0$$

and

$$\bar{C}\delta^2 \le \epsilon \le \bar{\epsilon}.$$

We need to prove the conclusion (4.2) in a ball $B_{2\bar{\eta}}(X^*)$.

We distinguish three cases depending on whether X^* is close to L, close to P, or far from P.

In Case 2 and Case 3 we will use the following properties from Remark 3.4.

(4.24)
$$c \leq \frac{V_n}{U_n} \leq C$$
, $|\triangle V| \leq C\delta^2 U_n$ in $B_2(X^*) \setminus (P \cup \{|(x_n, s)| \leq 20\delta |x'|\})$

Below η is the universal constant from Proposition 4.3.

Case 1. $|X^*| < \eta/4$.

In this case, since $B_1 \subset B_2(X^*)$ we follow under the assumptions of Proposition 4.3. Hence we can conclude that for any $\bar{\eta} \leq \eta/4$ in $B_{2\bar{\eta}}(X^*) \subset B_{\eta}$ either

$$g(X) \le V(X + (1 - \eta)\epsilon e_n),$$

or

$$g(X) \ge V(X - (1 - \eta)\epsilon e_n)$$

and our conclusion is satisfied for all $\bar{\eta} \leq \eta/4$.

Case 2. $|X^*| \ge \eta/4$, and $B_{\frac{\eta}{32}}(X^*) \cap P = \emptyset$. In this case, if $\bar{\epsilon}$ is small enough then it follows from (4.24) that the function

$$h(X) := g(X) - V(X - \epsilon e_n) \ge 0$$

satisfies

$$|\Delta h| \le C\delta^2 U_n \quad \text{in } B := B_{\frac{\eta}{64}}(X^*).$$

Notice also that by Harnack inequality

(4.25)
$$\frac{U_n(X)}{U_n(Y)} \le C \quad \text{for } X, Y \in B$$

with C universal. Assume that

$$g(X^*) \ge V(X^*).$$

Then, in view of (4.24) and (4.25)

$$h(X^*) = g(X^*) - V(X^* - \epsilon e_n) \ge c\epsilon U_n(X^*).$$

Hence by Harnack inequality, (4.25) and the condition $\bar{C}\delta^2 \leq \epsilon$

$$h \ge c\epsilon U_n(X^*) - C\delta^2 ||U_n||_{L^{\infty}(B)} \ge c'\epsilon U_n(X^*) \text{ in } B_{\frac{\eta}{128}}(X^*).$$

Thus, using (4.24) we have that for τ small enough

$$h \ge c' \epsilon \sup_{B} V_n \ge V(X - (1 - \tau)\epsilon e_n) - V(X - \epsilon e_n) \quad \text{in } B_{\frac{\eta}{128}}(X^*),$$

from which our desired conclusion follows with any $\bar{\eta}$ such that $2\bar{\eta} \leq \min\{\eta/128, \tau\}$.

Case 3. $|X^*| \ge \eta/4$ and $B_{\frac{\eta}{32}(X^*)} \cap P \neq \emptyset$.

In this case we argue similarly as in the previous case but we need to make use of the boundary Harnack inequality.

Assume that $X^* \in \{s > 0\}$ and call $X_0^* = (x^*, 0)$ the projection of X^* onto $\{s = 0\}$. If $\bar{\epsilon}$ is small enough then it follows from (4.24) that the function

$$h(X) := g(X) - V(X - \epsilon e_n) \ge 0,$$

satisfies

$$|\Delta h| \le C\delta^2 U_n$$
 in $B := B_{\frac{\eta}{8}}(X_0^*) \cap \{s > 0\},$

for a universal constant C. Denote by $Y^* = X_0^* + \frac{\eta}{16}e_n$ and assume that

 $g(Y^*) \ge V(Y^*).$

As in the previous case, by Harnack inequality

(4.26)
$$h \ge c\epsilon U_n(Y^*) \quad \text{in } B_{\frac{\eta}{22}}(Y^*).$$

Now we argue similarly as in Lemma 4.4.

Denote by

$$D := (B_{\eta/8}(X_0^*) \setminus B_{\eta/32}(Y^*)) \cap \{s > 0\}.$$

Let q_1, q_2 satisfy in D

$$\Delta q_1 = 0, \quad \Delta q_2 = -1$$

with boundary conditions respectively,

$$q_1 = 1$$
 on $\partial B_{\eta/32}(Y^*)$, $q_1 = 0$ on $\partial (B_{\eta/8}(X_0^*) \cap \{s > 0\})$

and

$$q_2 = 0 \quad \text{on } \partial D$$

By the maximum principle, in view of (4.26) we obtain that

$$h \ge c\epsilon U_n(Y^*)q_1 - C\delta^2 q_2$$
 in D_1

Moreover,

$$q_1 \ge cq_2$$
 in $D \cap B_{\eta/16}(X_0^*)$.

Hence using that $\bar{C}\delta^2 \leq \epsilon$ we get

$$h(X) \ge c' \epsilon U_n(Y^*) q_1(X) \ge c \epsilon U_n(X) \quad \text{in } B_{\eta/16}(X_0^*) \cap \{s > 0\}$$

where in the last inequality we used that $U_n(Y^*)q_1$ is comparable to U_n in view of boundary Harnack inequality.

Now we use (3.1) and (4.24) to conclude

$$h(X) = h(x, x_{n+1}) \ge c\epsilon \sup_{B_{\frac{\eta}{8}}(X_0^*)} U_n(y, x_{n+1}) \ge c\epsilon \sup_{B_{\frac{\eta}{8}}(X_0^*)} V_n(y, x_{n+1})$$

$$\ge V(X - (1 - \tau)\epsilon e_n) - V(X - \epsilon e_n) \quad \text{in } B_{\frac{\eta}{16}}(X_0^*) \supset B_{\frac{\eta}{32}}(X^*).$$

Then our desired statement holds for $\bar{\eta} \leq \min\{\tau/2, \eta/64\}$.

5. Improvement of flatness.

In this section we prove our main Theorem 1.1. We start with the following quadratic improvement of flatness proposition. We show that if a solution g stays in a $\lambda^{2+\alpha}$ neighborhood of a function $V \in \mathcal{V}_1^0$ in a ball B_{λ} then in $B_{\eta\lambda}$, g is in a $(\lambda\eta)^{2+\alpha}$ neighborhood of another function V in the same class.

Proposition 5.1. Given $\alpha \in (0,1)$, there exist $\lambda_0, \eta_0 \in (0,1)$ and C > 0 large depending on α and n, such that if g solves (2.1), $0 \in F(g)$ and g satisfies

(5.1)
$$V(X - \lambda^{2+\alpha} e_n) \le g(X) \le V(X + \lambda^{2+\alpha} e_n), \quad in \ B_\lambda \ with \ 0 < \lambda \le \lambda_0$$

for $V = V_{M,0,a,b} \in \mathcal{V}_1^0$, then in a possibly different system of coordinates denoted by $\overline{E} = \{\overline{e}_1, \ldots, \overline{e}_n, \overline{e}_{n+1}\},\$

(5.2)
$$\bar{V}(X - (\eta_0 \lambda)^{2+\alpha} \bar{e}_n) \le g(X) \le \bar{V}(X + (\eta_0 \lambda)^{2+\alpha} \bar{e}_n), \quad in \ B_{\eta_0 \lambda}$$

for some $\bar{V} = V_{\bar{S},\bar{a},\bar{b}}$ (defined in (3.3)) with \bar{S} given in the \bar{E} coordinates by

$$\bar{\mathcal{S}} = \{ \bar{x}_n = \frac{1}{2} (\bar{x}')^T \bar{M} \bar{x}' \}$$

and

$$|\bar{M} - M||, |\bar{a} - a|, |\bar{b} - b| \le C\lambda^{\alpha}, \quad \bar{a} + \bar{b} - tr\bar{M} = 0$$

Moreover, for any $\sigma \in (0, 1]$, the surfaces \bar{S} and S separate in B_{σ} at most $C(\lambda^{\alpha}\sigma^2 + \lambda^{1+\alpha}\sigma)$.

Proof. Let η_0, C be the constants in Corollary 2.10.

The proof is by compactness. Assume that no such λ_0 exists, then we can find a sequence of λ_k 's, tending to 0, g_k and V_k satisfying (5.1) for which (5.2) fails. We rescale g_k and V_k . For simplicity of notation we drop the dependence on k and denote

$$g_{\lambda}(X) = \lambda^{-1/2} g(\lambda X), \quad V_{\lambda}(X) = \lambda^{-1/2} V(\lambda X), \quad X \in B_1.$$

Notice that

$$V_{\lambda} = V_{\lambda M, 0, \lambda a, \lambda b} \in \mathcal{V}_{\lambda}^{0},$$

and

$$V_{\lambda}(X - \lambda^{1+\alpha}e_n) \le g_{\lambda}(X) \le V_{\lambda}(X + \lambda^{1+\alpha}e_n)$$
 in B_1

Let

$$\epsilon = \lambda^{1+\alpha}, \quad \delta = \lambda$$

and define

(5.3)
$$w_{\lambda} := \frac{\tilde{g}_{\lambda} - \gamma_{V_{\lambda}}}{\epsilon}$$

Thus by Proposition 3.5

$$w_{\lambda} = \frac{\tilde{g}_{\lambda} - \tilde{V}_{\lambda}}{\epsilon} + \frac{\tilde{V}_{\lambda} - \gamma_{V_{\lambda}}}{\epsilon} = \widetilde{(g_{\lambda})}_{\epsilon, \tilde{V}_{\lambda}} + O(\frac{\delta^2}{\epsilon})$$

and hence by Corollary 4.2 we get that w_{λ} converges uniformly to a Holder continuous function w_0 as k tends to ∞ (and $\lambda \to 0$), with $w_0(0) = 0$ and $|w_0| \leq 1$.

We claim that w_0 is a viscosity solution of the linearized problem

(5.4)
$$\begin{cases} \Delta(U_n w_0) = 0, & \text{in } B_{1/2} \setminus P, \\ |\nabla_r w_0| = 0, & \text{on } B_{1/2} \cap L. \end{cases}$$

We start by showing that $U_n w_0$ is harmonic in $B_{1/2} \setminus P$.

Let $\tilde{\varphi}$ be a smooth function which touches w_0 strictly from below at $X_0 \in B_{1/2} \setminus P$. We need to show that

(5.5)
$$\Delta(U_n\tilde{\varphi})(X_0) \le 0.$$

Since w_{λ} converges uniformly to w_0 in $B_{1/2} \setminus P$ we conclude that there exist a sequence of constants $c_{\lambda} \to 0$ and a sequence of points $X_{\lambda} \in B_{1/2} \setminus P$, $X_{\lambda} \to X_0$ such that $\tilde{\psi}_{\lambda} := \epsilon(\tilde{\varphi} + c_{\lambda}) + \gamma_{\tilde{V}_{\lambda}}$ touches \tilde{g}_{λ} by below at X_{λ} for a sequence of λ 's tending to 0.

Define the function ψ_{λ} by the following identity

(5.6)
$$\psi_{\lambda}(X - \bar{\psi}_{\lambda}(X)e_n) = U(X).$$

Then according to (2.8) ψ_{λ} touches g_{λ} from below at $Y_{\lambda} = X_{\lambda} - \tilde{\psi}_{\lambda}(X_{\lambda})e_n \in B_1^+(g_{\lambda})$. Thus, since g_{λ} satisfies (2.1) in B_1 it follows that

(5.7)
$$\Delta \psi_{\lambda}(Y_{\lambda}) \le 0.$$

In a neighborhood of X_0 , $\gamma_{V_{\lambda}}/\lambda$ has bounded C^k norms (depending on $|X_0|$) hence $\tilde{\psi}_{\lambda}/\lambda$ has also bounded C^k norms. By Proposition 2.8

$$\begin{aligned} \Delta \psi_{\lambda} &= \lambda \Delta (U_n(\psi_{\lambda}/\lambda)) + O(\lambda^2) \\ &= \Delta (U_n \tilde{\psi}_{\lambda}) + O(\lambda^2) \\ &= \Delta (U_n(\epsilon \tilde{\varphi} + \gamma_{\tilde{V}_{\lambda}}))(X_{\lambda}) + O(\lambda^2) \\ &= \epsilon \Delta (U_n \tilde{\varphi}) + O(\lambda^2) \end{aligned}$$

where we have used that

$$\Delta(U_n \gamma_{V_\lambda}) = 0$$

This can be checked either explicitly or by using Theorem 6.1. In conclusion

$$\epsilon \Delta(U_n \tilde{\varphi})(Y_\lambda) + O(\lambda^2) \le 0.$$

We divide by $\epsilon = \lambda^{1+\alpha}$ and let $\lambda \to 0$. Using that $Y_{\lambda} \to X_0$ we obtain

$$\Delta(U_n\tilde{\varphi})(X_0) \le 0,$$

as desired.

Next we need to show that

$$|\nabla_r w_0|(X_0) = 0, \quad X_0 = (x'_0, 0, 0) \in B_{1/2} \cap L,$$

in the viscosity sense of Definition 2.9.

We argue by contradiction. Assume for simplicity (after a translation) that there exists a function ϕ which touches w_0 by below at 0 with $\phi(0) = 0$ and such that

$$\phi(X) = \xi' \cdot x' + \beta r + O(|x'|^2 + r^{3/2}),$$

with

$$\beta > 0.$$

Then we can find constants σ, \tilde{r} small and A large such that the polynomial

$$q(X) = \xi' \cdot x' - \frac{A}{2}|x'|^2 + 2A(n-1)x_nr$$

touches ϕ by below at 0 in a tubular neighborhood $N_{\bar{r}} = \{|x'| \leq \tilde{r}, r \leq \tilde{r}\}$ of 0, with

$$\phi - q \ge \sigma > 0$$
, on $N_{\tilde{r}} \setminus N_{\tilde{r}/2}$.

This implies that

(5.8) $w_0 - q \ge \sigma > 0, \quad \text{on } N_{\tilde{r}} \setminus N_{\tilde{r}/2},$

and

(5.9)
$$w_0(0) - q(0) = 0.$$

In particular, by continuity near the origin we can find a point X^* such that

(5.10)
$$w_0(X^*) - q(X^*) \le \frac{o}{8}, \quad X^* \in N_{\tilde{r}} \setminus P \text{ close to } 0.$$

Now, let us define

$$W_{\lambda} := V_{\lambda M + A\epsilon I, -\epsilon\xi', \lambda a, \lambda b + 2\epsilon A(n-1)} \in \mathcal{V}_{2\delta}.$$

Then in view of Proposition 3.5 we have

$$\widetilde{W}_{\lambda} = \epsilon q + \widetilde{\gamma}_{V_{\lambda}} + O(\delta^2)$$

and moreover, W_{λ} is a subsolution to our problem since $\epsilon \gg \delta^2$.

Thus, from the uniform convergence of w_{λ} to w_0 and (5.8) we get that (for all λ small)

(5.11)
$$\frac{\tilde{g}_{\lambda} - \widetilde{W}_{\lambda}}{\epsilon} = w_{\lambda} - q + O(\frac{\delta^2}{\epsilon}) \ge \frac{\sigma}{2} \quad \text{in } (N_{\tilde{r}} \setminus N_{\tilde{r}/2}) \setminus P.$$

Similarly, from the uniform convergence of w_{λ} to w_0 and (5.10) we get that for k large

(5.12)
$$\frac{(\tilde{g}_{\lambda} - \tilde{W}_{\lambda})(X^*)}{\epsilon} \le \frac{\sigma}{4}, \quad \text{at } X^* \in N_{\tilde{r}} \setminus P.$$

On the other hand, it follows from Lemma 2.7 and (5.11) that

$$\frac{\tilde{g}_{\lambda} - \tilde{W}_{\lambda}}{\epsilon} \ge \frac{\sigma}{2} \quad \text{in } N_{\tilde{r}} \setminus P,$$

which contradicts (5.12).

In conclusion w_0 solves the linearized problem. Hence, by Corollary 2.10 since $w_0(0) = 0, w_0$ satisfies

$$(5.13) \quad -\frac{1}{4}\eta_0^{2+\alpha} \le w_0(X) - (\xi_0 \cdot x' + \frac{1}{2}x'^T M_0 x' - \frac{a_0}{2}r^2 - b_0 x_n r) \le \frac{1}{4}\eta_0^{2+\alpha} \quad \text{in } B_{4\eta_0},$$

for some $\eta_0 \in (0, 1)$ universal and with

$$a_0 + b_0 - tr M_0 = 0, \qquad |\xi_0|, ||M_0||, |a_0|, |b_0| \le C.$$

From the uniform convergence of w_{λ} to w_0 , we get that for all k large enough

(5.14)
$$-\frac{1}{2}\eta_0^{2+\alpha} \le w_\lambda(X) - \frac{T_\lambda - \gamma_{\tilde{V}_\lambda}}{\epsilon} \le \frac{1}{2}\eta_0^{2+\alpha} \quad \text{in } B_{4\eta_0} \setminus P,$$

with

$$I_{\lambda} := V_{\lambda M - \epsilon M_0, -\epsilon \xi_0, \lambda a - \epsilon a_0, \lambda b - \epsilon b_0}.$$

In conclusion, from the definition (5.3) of w_{λ} , we get

(5.15)
$$\tilde{T}_{\lambda} - \frac{\epsilon}{2} \eta_0^{2+\alpha} \le \tilde{g}_{\lambda} \le \tilde{T}_{\lambda} + \frac{\epsilon}{2} \eta_0^{2+\alpha},$$

or

$$T_{\lambda}(X - \frac{\epsilon}{2}\eta_0^{2+\alpha}e_n) \le g_{\lambda}(X) \le T_{\lambda}(X + \frac{\epsilon}{2}\eta_0^{2+\alpha}e_n) \quad \text{in } B_{2\eta_0}.$$

We rescale g_{λ} back from the ball B_1 to B_{λ} and obtain

(5.16)
$$T(X - \frac{\epsilon\lambda}{2}\eta_0^{2+\alpha}e_n) \le g(X) \le T(X + \frac{\epsilon\lambda}{2}\eta_0^{2+\alpha}e_n) \quad \text{in } B_{2\lambda\eta_0},$$
with

for

$$\mathcal{S}_T := \{ x_n = \frac{1}{2} (x')^T M_T x' + \xi_T \cdot x' \},$$
$$M_T := M - \frac{\epsilon}{\lambda} M_0, \quad \xi_T := -\epsilon \xi_0, \quad a_T := a - \frac{\epsilon}{\lambda} a_0, \quad b_T := b - \frac{\epsilon}{\lambda} b_0$$

 $T = V_{\mathcal{S}_T, a_T, b_T},$

Next we show that in a different system of coordinates, called \overline{E} , the function T can be approximated by $V_{M_T,0,a_T,b_T}$.

Assume for simplicity that ξ_T points in the e_1 direction. Then we choose an orthogonal system of coordinates $\bar{E} := \{\bar{e}_1, \bar{e}_2, \dots, \bar{e}_{n+1}\}$ with

$$\bar{e}_i = e_i, \quad \text{if } i \neq 1, n$$

and \bar{e}_n normal to S_T at 0.

Notice that the E system of coordinates is obtained from the standard one after an orthogonal transformation of norm bounded by $C|\xi_T|$ which is smaller than $C\epsilon$.

A point in this system on coordinates is denoted by \overline{X} . We let,

$$\bar{\mathcal{S}} := \{ \bar{x}_n = \frac{1}{2} (\bar{x}')^T M_T \bar{x}' \},\$$

and we write \bar{S} as a graph in the e_n direction, that is

$$\bar{\mathcal{S}} := \{x_n = h(x')\}.$$

We claim that in a ball of radius σ the distance (in the e_n direction) between S_T and \bar{S} in B_{σ} is less that $C\epsilon\sigma^2$, for any $0 < \sigma \leq 1$.

Indeed, since $\bar{x} = Ox$ with O orthogonal and $||O - I|| \le C\epsilon$, we obtain by implicit differentiation

$$||D_{x'}^2 h - M_T||_{L^{\infty}(B_1)} \le C\epsilon, \quad \nabla_{x'} h(0) = \xi_T.$$

Thus in $B_{2\eta_0\lambda}$ we have that the surfaces

$$S_T \pm \frac{\epsilon}{2} \lambda \eta_0^{2+\alpha} e_n$$

lie between

$$\bar{S} \pm \epsilon \lambda \eta_0^{2+\alpha} \bar{e}_n$$

since $C\epsilon(\eta_0\lambda)^2 \ll \frac{\epsilon}{2}\lambda\eta_0^{2+\alpha}$. In view of this inclusion, using that $v_{a_{\underline{T}},b_{\underline{T}}}(t,s)$ is monotone in t, we obtain from (5.16) the desired conclusion (5.2) with $\overline{M} = M_T, \overline{a} = a_T, \overline{b} = b_T$.

Since the distance between \mathcal{S}_T and S in B_σ is less than $C(\frac{\epsilon}{\lambda}\sigma^2 + \epsilon\sigma)$ the proof is finished.

We can now prove our main Theorem 1.1. In fact we show that under our flatness assumption, a solution g can be approximated in a $C^{2,\alpha}$ fashion by a function $V \in \mathcal{V}_C^0$.

Theorem 5.2. There exists $\bar{\epsilon} > 0$ small universal such that if g solves (2.1) in B_1 with

(5.17)
$$\{x \in \mathcal{B}_1 : x_n \le -\bar{\epsilon}\} \subset \{x \in \mathcal{B}_1 : g(x,0) = 0\} \subset \{x \in \mathcal{B}_1 : x_n \le \bar{\epsilon}\},$$

then in an appropriate system of coordinates denoted by \bar{e}_i

$$V(\bar{X} - C\lambda^{2+\alpha}\bar{e}_n) \le g(\bar{X}) \le V(\bar{X} + C\lambda^{2+\alpha}\bar{e}_n) \quad in \ B_\lambda, \ for \ all \ 0 < \lambda < 1/C,$$

for some $V = V_{M_0,0,a_0,b_0} \in \mathcal{V}_C^0$, with C depending on n and α . In particular, $F(g) \cap \mathcal{B}_{1/2}$ is a $C^{2,\alpha}$ graph in the e_n direction for any $\alpha \in (0,1)$.

Proof. It suffices to prove the theorem for any fixed $\alpha \in (0,1)$ for some $\bar{\epsilon}(\alpha)$, $C(\alpha)$ depending on α . The dependence of $\bar{\epsilon}$ on α can be easily removed by fixing $\bar{\epsilon} := \bar{\epsilon}(\bar{\alpha})$, say with $\bar{\alpha} = 1/2$. Then by the conclusion (5.2) for $\bar{\alpha}$, appropriate rescalings of q satisfy the flatness assumption (5.17) also for $\bar{\epsilon}(\alpha)$ for any $\alpha \in (0, 1)$.

By Lemma 7.9 the rescaling

$$g_{\mu}(X) = \mu^{-\frac{1}{2}}g(\mu X)$$

satisfies

$$U(X - \tau^{2+\alpha} e_n) \le g_\mu(X) \le U(X + \tau^{2+\alpha} e_n) \quad \text{in } B_1,$$

provided that $\bar{\epsilon}$, μ are chosen small depending on $\tau \leq \lambda_0$, with λ_0 the universal constant in Proposition 5.1 and τ small universal to be made precise later. Thus g_{μ} satisfies in B_{τ} the hypotheses of Proposition 5.1 with M = 0, a = 0, b = 0. Then we can apply Proposition 5.1 repeatedly for all $\tau_k := \tau \eta_0^k$ since by choosing τ small enough we can guarantee that $\sum_k C \tau_k^{\alpha} \leq 1$ and hence the corresponding M_k , a_k , b_k have always norm less than 1. Thus we obtain

(5.18)
$$V_{\mathcal{S}_k, a_k, b_k}(X - \tau_k^{2+\alpha} e_n^k) \le g_\mu(X) \le V_{\mathcal{S}_k, a_k, b_k}(X + \tau_k^{2+\alpha} e_n^k) \quad \text{in } B_{\tau_k}.$$

Using that S_k and S_{k+1} separate (in the e_n -direction) in B_σ at most $C(\tau_k^\alpha \sigma^2 + \tau_k^{1+\alpha}\sigma)$ we conclude that as $k \to \infty$, the paraboloids S_k converge uniformly in B_1 to a limit parabolid S_* . Moreover, S_* also separates from S_k in B_σ by at most $C(\tau_k^\alpha \sigma^2 + \tau_k^{1+\alpha}\sigma)$ in the e_n^* direction where e_n^* is the normal to S_* at the origin. Finally, as $k \to \infty$, $a_k \to a_*$, $b_k \to b_*$, with

$$|a_k - a_*|, |b_k - b_*| \le C\tau_k^{\alpha}.$$

Now notice that in $B_{2\tau_k}$, the paraboloids S_k and S_* separate at most $C\tau_k^{2+\alpha}$, thus we can apply Lemma 3.7 and use the inequality (5.18) to obtain

$$V_{\mathcal{S}_*, a_*, b_*}(X - C\tau_k^{2+\alpha} e_n^*) \le g_{\mu}(X) \le V_{\mathcal{S}_*, a_*, b_*}(X + C\tau_k^{2+\alpha} e_n^*), \quad \text{in } B_{\tau_k}.$$

Rescaling back we obtain the desired claim.

6. The regularity of the linearized problem

We recall that the linearized problem associated to (2.1) is

(6.1)
$$\begin{cases} \Delta(U_n h) = 0, & \text{in } B_1 \setminus P, \\ |\nabla_r h| = 0, & \text{on } B_1 \cap L, \end{cases}$$

where

$$|\nabla_r h|(X_0) := \lim_{(x_n, s) \to (0, 0)} \frac{h(x'_0, x_n, s) - h(x'_0, 0, 0)}{r}, \quad r^2 = x_n^2 + s^2.$$

In this section we obtain a second order expansion near the origin for a solution h to (6.1).

Theorem 6.1. Let h be a solution to (6.1) such that $|h| \leq 1$. Then h satisfies

(6.2)
$$|h(X) - (h(0) + \xi_0 \cdot x' + \frac{1}{2}(x')^T M_0 x' - \frac{a_0}{2}r^2 - b_0 r x_n)| \le C|X|^3,$$

for some a_0, b_0, ξ_0, M_0 with $|\xi_0|, |a_0|, |b_0|, ||M_0|| \leq C, C$ universal and

$$a_0 + b_0 - tr M_0 = 0$$

Proof. This proof is a refinement of Theorem 8.1 in [DR] where the authors obtained a first order expansion for h, in particular

(6.3)
$$|h(X) - h(X_0)| = O(|X - X_0|), \quad X_0 \in L.$$

Also in [DR] it is shown that h and its derivatives of all orders in the x' direction are Holder continuous with norm controlled by a universal constant in $B_{1/2}$ (see Corollary 8.7.)

We wish to prove that

(6.4)
$$|h(x', x_n, x) - h(x', 0, 0) + \frac{a(x')}{2}r^2 + b(x')x_nr| \le Cr^3, \quad (x', 0, 0) \in B_{1/2} \cap L,$$

with C universal and $h(\cdot, 0, 0), a, b$ smooth functions of x'.

The function h solves

$$\Delta(U_n h) = 0 \quad \text{in } B_1 \setminus P$$

and since U_n is independent on x' we can rewrite this equation as

(6.5)
$$\Delta_{x_n,s}(U_nh) = -U_n\Delta_{x'}h.$$

Moreover, since $\Delta_{x'}h$ solves the same linear problem as h then any estimate for h also holds for $\Delta_{x'}h$.

For each fixed x', we investigate the 2-dimensional problem

$$\Delta(U_t h) = U_t f, \quad \text{in } B_{1/2} \setminus \{t \le 0, s = 0\} \subset \mathbb{R}^2$$

with $h, f \in C^{0,\beta}$. Without loss of generality, for a fixed x' we may assume h(x', 0, 0) = 0. Thus in view of (6.3), the function

$$H := U_t h$$

is continuous in $B_{1/2} \subset \mathbb{R}^2$ and satisfies

$$\Delta H = U_t f \quad \text{in } B_{1/2} \setminus \{ t \le 0, s = 0 \}, \quad H = 0 \quad \text{on } B_{1/2} \cap \{ t \le 0, s = 0 \}.$$

Now, we consider the holomorphic transformation $z \to \frac{1}{2}z^2$

$$\Phi: (\zeta, y) \to (t, s) = (\frac{1}{2}(\zeta^2 - y^2), \zeta y)$$

which maps $B_1 \cap \{\zeta > 0\}$ into $B_{1/2} \setminus \{t \le 0, s = 0\}$ and call

$$\tilde{h}(\zeta, y) = h(t, s), \quad \tilde{f}(\zeta, y) = f(t, s), \quad \tilde{H}(\zeta, y) = H(t, s)$$

with $(\tilde{r}, \tilde{\theta})$, the polar coordinates in the (ζ, y) plane. Then, easy computations show that

(6.6)
$$\Delta \tilde{H} = \zeta \tilde{f} \quad \text{in } B_1 \cap \{\zeta > 0\}, \quad \tilde{H}(\zeta, y) = \frac{\zeta}{\tilde{r}^2} \tilde{h},$$

and

$$\tilde{H} = 0 \quad \text{on } \{\zeta = 0\}.$$

Since the right-hand side is in $C^{0,\beta}$ and \tilde{h}, \tilde{f} have the same regularity, we conclude from repeatedly applying Lemma 6.2 below that $\tilde{h}, \tilde{f} \in C^{\infty}$ with

$$||f||_{C^{k,\beta}(B^+_{1/2})}, ||h||_{C^{k,\beta}(B^+_{1/2})} \le C(k,\beta).$$

Notice that we can reflect \tilde{H} oddly and \tilde{h}, \tilde{f} evenly across $\{\zeta = 0\}$ and the resulting functions will still solve (6.6) in B_1 . Moreover from our assumptions, \tilde{f} and \tilde{h} are even with respect to y. Thus, we conclude that the Taylor polynomials for \tilde{f}, \tilde{h}

around the origin, are polynomials in ζ^2, y^2 . Now we use the Taylor expansion for \tilde{H} around 0, which is odd with respect to ζ and even in y, that is

$$\tilde{H}(\zeta, y) = \zeta (d_0 + d_1 \zeta^2 + d_2 y^2 + O(\tilde{r}^4))$$

with

$$6d_1 + 2d_2 = f(0) = f(0).$$

Thus,

$$\tilde{h}(\zeta, y) = \tilde{\zeta}^2 (d_0 + d_1 \zeta^2 + d_2 y^2 + O(\tilde{r}^4)).$$

In terms of the (t, s) coordinates this means that

$$h(t,s) = 2r(d_0 + 2d_1r(\cos\frac{\theta}{2})^2 + 2d_2r(\sin\frac{\theta}{2})^2) + O(r^3) =$$

= $2r(d_0 + (d_1 + d_2)r + (d_1 - d_2)t) + O(r^3)$
= $2d_0r - \frac{a}{2}r^2 - btr + O(r^3).$

In conclusion,

$$|h(X) - h(x', 0, 0) - 2d_0(x')r + \frac{a(x')}{2}r^2 + b(x')x_nr| \le Cr^3$$
 in $B_{1/2}$

with C universal,

$$||a||_{L^{\infty}(\{|x'|\leq 1/2\})}, ||b||_{L^{\infty}(\{|x'|\leq 1/2\})} \leq C,$$

and

$$a+b=\Delta_{x'}h(x',0,0).$$

Since h solves (6.1) we must also have $d_0(x') = 0$ and hence

$$|h(X) - h(x', 0, 0) + \frac{a(x')}{2}r^2 + b(x')x_nr| \le Cr^3$$
, in $B_{1/2}$.

Notice that a, b are smooth functions of x' with all order derivatives bounded by appropriate universal constants. Indeed due to the linearity of the problem it is easy to see that $D_{x'}^{\beta}a, D_{x'}^{\beta}b$ are the corresponding a and b for $D_{x'}^{\beta}h$. Writing the Taylor expansions at 0 for h(x', 0, 0) up to order 2 and a, b up to order 1 with $a_0 = a(0), b_0 = b(0)$ we get

$$\left| h(X) - \left(h(0) + \xi_0 \cdot x' + \frac{1}{2} (x')^T M_0 x' - \frac{a_0}{2} r^2 - b_0 x_n r \right) \right| \le C |X|^3.$$

In our proof above we used the following easy lemma.

Lemma 6.2. Let $\tilde{H} = \tilde{H}(\zeta, y)$ be a function defined on $\overline{B_1^+} \subset \mathbb{R}^2$, which vanishes continuously on $\{\zeta = 0\}$. If $\tilde{H} \in C^{k,\alpha}(\overline{B_1^+})$, $k \in \mathbb{N}, \alpha \in (0,1]$, then $\zeta^{-1}\tilde{H} \in C^{k-1,\alpha}(\overline{B_1^+})$, and

$$\|\zeta^{-1}H\|_{C^{k-1,\alpha}} \le \|H\|_{C^{k,\alpha}}.$$

Proof. Since $\tilde{H}(0, y) = 0$ we see that

$$\zeta^{-1}\tilde{H}(\zeta,y) = \int_0^1 \tilde{H}_\zeta(t\zeta,y) \, dt$$

and the lemma follows easily by taking derivatives in the equality above.

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7. Basic properties of a solution g.

We collect here some useful general facts about solutions g to our free boundary problem (2.1), such as $C^{1/2}$ -optimal regularity, asymptotic expansion near regular points of the free boundary and compactness.

First we recall some notation. Let $v \in C(B_1)$ be a non-negative function. We denote by

$$B_1^+(v) := B_1 \setminus \{(x,0) : v(x,0) = 0\} \subset \mathbb{R}^{n+1}$$

and by

$$F(v) := \partial_{\mathbb{R}^n} (B_1^+(v) \cap \mathcal{B}_1) \cap \mathcal{B}_1 \subset \mathbb{R}^n.$$

Also, we denote by P the half-hyperplane

$$P := \{ X \in \mathbb{R}^{n+1} : x_n \le 0, s = 0 \}.$$

Given a C^2 surface S in \mathbb{R}^{n-1} , we often work with functions of the form $V = V_{S,a,b}$ (see Definition 3.3). We remark that we can still apply the boundary Harnack inequality with V in a neighborhood of S since in this set V is comparable with a harmonic function H with F(H) = S.

Indeed, after a dilation we may assume that $V = V_{S,a,b} \in \mathcal{V}_{\delta}$, that is the curvatures of S in B_2 and |a|, |b| are bounded by δ small, universal. Let

$$V_1 := V_{\mathcal{S}, a-2n\delta, b}, \quad V_2 := V_{\mathcal{S}, a+2n\delta, b}$$

and notice that V_1 is a supersolution and V_2 is a subsolution in B_1 (see Proposition 3.2). Also

$$1/2V_2 \le V \le 2V_1 \le 2V_2,$$

hence there exists H between $1/2V_2$ and $2V_1$, with $1/4V \le H \le 4V$, H harmonic in $\{H > 0\}$ and F(H) = S.

We obtain the following version of the boundary Harnack inequality.

Lemma 7.1. Let $V := V_{\mathcal{S},a,b} \in \mathcal{V}_{\delta_0}$, for some small δ_0 universal and with $0 \in \mathcal{S}$. Let $w \in C(\overline{B}_1)$ be a non-negative function which is harmonic in $B_1^+(w)$. If $B_1^+(V) \subset B_1^+(w)$ then

$$w \ge cw(\frac{1}{2}e_n)V, \quad in \ B_{1/2}.$$

If $B_1^+(V) \subset B_1^+(w)$ then

$$w \le C \|w\|_{L^{\infty}(B_1)} V$$
, in $B_{1/2}$.

Proof. Let \bar{w} be the harmonic function in $B_{3/4}^+(V)$ with boundary value w on $\partial B_{3/4}$ and $\bar{w} = 0$ where $\{V = 0\}$.

If $B_1^+(V) \subset B_1^+(w)$ then in view of the observation above we can apply the boundary Harnack inequality with V and conclude that

$$w \ge \bar{w} \ge c \,\bar{w}(\frac{1}{2}e_n)V,$$
 in $B_{1/2}$.

On the other hand,

$$\bar{w}(\frac{1}{2}e_n) \ge c \inf_{B_{1/4}(\frac{3}{4}e_n) \cap \partial B_{3/4}} \bar{w}.$$

Using that w and \bar{w} coincide on $\partial B_{3/4}$ together with Harnack inequality we obtain our desired estimate. If $B_1^+(V) \subset B_1^+(w)$ then

$$w \le \bar{w} \le C\bar{w}(\frac{1}{2}e_n)V$$
, in $B_{1/2}$.

On the other hand,

$$\bar{w}(\frac{1}{2}e_n) \le \|\bar{w}\|_{L^{\infty}(B_{3/4})} = \|w\|_{L^{\infty}(B_{3/4})}$$

which yields our conclusion.

An immediate consequence is the following useful lemma.

Lemma 7.2. Let $V = V_{Sa,b} \in \mathcal{V}_{\delta_0}$ be a subsolution in B_1 , for some small universal δ_0 and with $0 \in S$. If w is harmonic in $B_1^+(w)$ and $B_1^+(V) \subset B_1^+(w)$ and

$$w \geq V - \epsilon$$
 in B_1 ,

then

$$w \ge (1 - C\epsilon)V$$
 in $B_{1/2}$.

Proof. Let $q: B_{3/4} \to \mathbb{R}$ be the harmonic function in $B_{3/4} \cap B_1^+(V)$ which has boundary values q = 1 on $\partial B_{3/4}$ and q = 0 on the set where V = 0. From our hypotheses on w and the maximum principle we obtain

$$w \ge V - \epsilon q$$
 in $B_{3/4}$.

On the other hand by Lemma 7.1, since $B_{3/4}^+(V) = B_{3/4}^+(q)$ we have $q \leq CV$ in $B_{1/2}$, which together with the inequality above implies the desired result.

Remark 7.3. From the proof of Lemma 7.2 we see that if the hypotheses on w hold only outside of the ball $B_{1/8}$, i.e

 $w \ge V - \epsilon$ on $B_1 \setminus B_{1/8}$, w harmonic in $B_1^+(V) \setminus B_{1/8}$

then the conclusion holds in the shell $B_{3/4} \setminus B_{1/4}$.

Next we prove optimal $C^{1/2}$ regularity for viscosity solutions.

Lemma 7.4 ($C^{1/2}$ -Optimal regularity). Assume g solves (2.1) in B_1 and $0 \in F(g)$. Then

$$g(x,0) \le C |d(x)|^{1/2} \quad in \ \mathcal{B}_{1/2}$$

where d(x) represents the distance from x to F(g). Also

$$||g||_{C^{1/2}(B_{1/2})} \le C(1+g(\frac{1}{2}e_{n+1})).$$

Proof. The first assertion follows in a standard way from the free boundary condition. By scaling, we need to show that if g is defined in B_2 , $0 \in F(g)$ and $\mathcal{B}_1(e_n) \subset B_2^+(g)$ then $u(e_n) \leq C$ for some large C universal.

By a rescaled version of Lemma 7.1 and Harnack inequality we have that in a neighborhood of 0,

$$g \ge cg(e_n)V_{\mathcal{S},2n,0}, \quad \mathcal{S} = \partial \mathcal{B}_1(e_n)$$

with $V_{S,2n,0}$ a subsolution near 0. The free boundary condition gives $1 \ge cg(e_n)$ which provides a bound for $g(e_n)$.

For the second inequality we write

 $g = g_0 + g_1$ in $D := B_{3/4} \cap \{s > 0\},\$

with g_0, g_1 harmonic in D and satisfying the following boundary conditions

$$g_0 = g$$
 on $\{s = 0\} \cap \partial D$, $g_0 = 0$ on $\{s > 0\} \cap \partial D$,

$$g_1 = 0$$
 on $\{s = 0\} \cap \partial D$, $g_1 = g$ on $\{s > 0\} \cap \partial D$.

From our estimate for g on $\{s = 0\}$, we obtain

$$||g_0||_{C^{1/2}(B_{1/2}\cap D)} \le C ||g||_{C^{1/2}(\mathcal{B}_{3/4})} \le C,$$

which together with the bound

$$||g_1||_{C^{1/2}(B_{1/2}\cap D)} \le Cg_1(\frac{1}{2}e_{n+1}) \le Cg(\frac{1}{2}e_{n+1}),$$

gives the desired conclusion.

Next we prove that if F(g) admits a tangent ball at 0 either from the positive or from the zero phase, then g has an asymptotic expansion of order $o(|X|^{1/2})$. This expansion also justifies our definition of viscosity solution to the free boundary problem (2.1). We remark however that this expansion holds also for an arbitrary harmonic function w which does not necessarily satisfy the free boundary condition.

Lemma 7.5 (Expansion at regular points from one side). Let $w \in C^{1/2}(B_1)$ be 1/2-Holder continuous, $w \ge 0$, with w harmonic in $B_1^+(w)$. If

$$0 \in F(w), \quad \mathcal{B}_{1/2}(1/2e_n) \subset B_1^+(w),$$

then

 $w = \alpha U + o(|X|^{1/2}), \quad for some \ \alpha > 0.$

The same conclusion holds for some $\alpha \geq 0$ if

$$\mathcal{B}_{1/2}(-1/2e_n) \subset \{w=0\}.$$

Proof. We define

$$\alpha := \inf_{\nu \notin P} \liminf_{t \to 0^+} \frac{w}{U}(t\nu).$$

First we notice that $\alpha > 0$. Indeed, by a rescaled version of Lemma 7.1

$$w \ge c w(\frac{1}{2}e_n) V_{\mathcal{S},0,0}, \quad \mathcal{S} = \partial \mathcal{B}_{\frac{1}{2}}(\frac{1}{2}e_n)$$

near the origin, for some c > 0. This implies that $\alpha \ge c w(\frac{1}{2}e_n) > 0$.

Assume by contradiction that the conclusion of the lemma does not hold with this choice of α . Then there exist $\delta_1 > 0$ and a sequence of points $y_k \to 0$ such that

(7.1)
$$|w(y_k) - \alpha U(y_k)| \ge \delta_1 |y_k|^{1/2}$$

Since w is 1/2-Holder continuous on B_1 , the rescalings

$$w_k(x) := |y_k|^{-\frac{1}{2}} w(|y_k|x),$$

are uniformly 1/2 Holder continuous and after passing to a subsequence we can assume that w_k converge uniformly on compact sets to a limiting function $w_* \in C(\mathbb{R}^n)$. We obtain

$$w_* \ge \alpha U, \quad \Delta w_* = 0 \quad \text{in } \mathbb{R}^n \setminus P_*$$

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and in view of (7.1) there exists a point y_* , $|y_*| = 1$ such that

$$w_*(y_*) \ge \alpha U(y_*) + \delta_1.$$

Using boundary Harnack inequality we find

(7.2)
$$w_* \ge \alpha (1+\delta_2) U \quad \text{in } B_1,$$

for some $\delta_2 > 0$ small. Now we let

$$V = V_{\frac{\delta_2}{2n}I, 0, \delta_2, 0}$$

and we notice that V is subharmonic in B_1 (by Proposition 3.2) and satisfies

$$V(X) = v_{\delta_2,0}(t,s) = (1 + \frac{\delta_2}{4}\rho)U(t,s)$$

$$\leq (1 + \frac{\delta_2}{2})U(t,s) \leq (1 + \frac{\delta_2}{2})U(x_n,s).$$

Thus (7.2) gives,

(7.3)
$$w_* \ge \alpha (1 + \frac{\delta_2}{4}) V \quad \text{in } B_1.$$

From the existence of a tangent ball at the origin included in $\{w > 0\}$ we see that for all large k, w_k is harmonic in the set where $\{V > 0\}$. Thus we conclude from (7.3) that in B_1

$$w_k \ge \alpha (1 + \frac{\delta_2}{4})V - \epsilon_k$$
, for some $\epsilon_k \to 0$.

By Lemma 7.2 we find that for all large k,

$$w_k \ge (1 - C\frac{\epsilon_k}{\alpha})\alpha(1 + \frac{\delta_2}{4})V \ge \alpha(1 + \frac{\delta_2}{8})V$$
 in $B_{1/2}$.

This implies that for any $\nu \notin P$

$$\liminf_{t \to 0^+} \frac{w}{U}(t\nu) = \liminf_{t \to 0^+} \frac{w_k}{U}(t\nu) \ge \alpha(1 + \frac{\delta_2}{8}),$$

which contradicts the minimality of α .

Remark 7.6. If we assume that F(w) admits a uniform tangent ball from its 0 side at all points in $\mathcal{B}_{1/2}$ then the hypothesis $w \in C^{1/2}(B_{1/4})$ is satisfied and therefore w has an expansion at all points in $F(w) \cap \mathcal{B}_{1/4}$. Indeed, by Lemma 7.1 we know that

$$w \le C \|w\|_{L^{\infty}} V_{\partial \mathcal{B}_r(x_0), 0, 0}$$

with $\partial \mathcal{B}_r(x_0)$ a tangent sphere to F(w) from the 0 side, and this implies

$$w(x) \le C \|w\|_{L^{\infty}} dist(x, \{w=0\})^{1/2}, \quad \forall x \in B_{1/4},$$

which gives $w \in C^{1/2}(B_{1/4})$.

In general, the term $o(|X|^{1/2})$ in the expansion for w can be improved in o(U) in the non-tangential direction to F(w). For example assume that $0 \in F(w) \in C^2$ and e_n is the normal to F(w) at 0 which points towards the positive phase. Then the non-tangential limit

$$\lim_{x \in \mathcal{C}, x \to 0} \ \frac{w}{U} = \alpha$$

where $\mathcal{C} \subset \mathbb{R}^n \setminus P$ is a cone whose closure does not contain $L = \{x_n = 0, s = 0\}$.

Indeed, by Lemma 7.5 and Remark 7.6 we have that $w = \alpha U + o(|X|^{1/2})$. Now the limit above follows by applying boundary Harnack inequality for U and w in the sets $\mathcal{C}_1 \cap (B_r \setminus B_{r/2})$ for all r small, where

$$\mathcal{C}_1 := \{ |x'| > \mu | (x_n, s) | \}$$

is such that $\overline{\mathcal{C}} \subset \mathcal{C}_1 \cup \{0\}$.

Remark 7.7. In the definition of viscosity solutions for our free boundary problem (see Definition 2.3) we can restrict the test functions only to the class of subsolutions and supersolutions of the form $cV_{S,a,b}$.

Precisely we say that g is a solution to (2.1) if

1) $\triangle g = 0$ in $B_1^+(g)$;

2) for any point $X_0 \in F(g)$ there exists no $V_{\mathcal{S},a,b}$ such that in a neighborhood of $X_0, V_{\mathcal{S},a,b}$ is a subsolution and

$$g \ge \alpha V_{S,a,b}, \quad \text{for some } \alpha > 1$$

with S touching strictly F(g) at X_0 from the positive side.

Analogously there is no supersolution $V_{\mathcal{S},a,b}$ such that

$$g \leq \alpha V_{S,a,b}$$
, for some $\alpha < 1$

and S touches strictly F(g) at X_0 from the 0 side.

In order to prove this statement we need to show that if we can touch g by below at a point $X_0 \in F(g)$ with a comparison subsolution v as in Definition 2.2, then we can touch also with a subsolution $\alpha V_{S,a,b}$ as above. A similar statement holds for supersolutions.

Assume for simplicity that $X_0 = 0$, e_n is normal to F(v) at 0 and $g \ge v$ in \overline{B}_1 . Let \overline{v} be the harmonic replacement for v in $B_1^+(v)$. In view of of Remark 7.6

$$\bar{v} = \alpha U + o(|X|^{1/2}), \text{ for some } \alpha \ge 1.$$

We claim that $\alpha > 1$. Indeed, $\bar{v} - v \ge 0$ is superharmonic in $B_1^+(v)$ and vanishes continuously on $\{v = 0\} \cap B_1$. If $\bar{v} - v \equiv 0$, then our claim follows from the definition of a comparison subsolution. Otherwise, by the boundary Harnack inequality

$$\bar{v} - v \ge \sigma \bar{v}$$

in a neighborhood of the origin, for some $\sigma > 0$. Thus $\bar{v} \ge v/(1-\sigma)$ near the origin and again the claim follows from the expansion of v at the origin.

The rescalings $v_k = r_k^{-1/2} v(r_k x)$ converge uniformly on compact sets to αU , with $\alpha > 1$. As in the proof of Lemma 7.5 we obtain that there exists δ small such that for all large k,

$$v_k \ge V := (1+\delta)V_{\frac{\delta}{2\pi}I,0,\delta,0}$$

and F(V) touches strictly $F(v_k)$ at the origin from the positive side. Rescaling back we obtain the desired conclusion.

Next we prove a compactness result for viscosity solutions to (2.1) whose free boundaries converge in the Hausdorff distance.

Proposition 7.8 (Compactness). Assume g_k solve (2.1) and converge uniformly to g_* in B_1 , and $\{g_k = 0\}$ converges in the Hausdorff distance to $\{g_* = 0\}$. Then g_* solves (2.1) as well.

Proof. Clearly g_* is harmonic in $B_1^+(g_*)$. In view of Remark 7.7 we need to check say that if $0 \in F(g_*)$ there exists no subsolution $V_{M,0,a,b}$ such that in a neighborhood of 0,

$$g_* \ge \alpha V_{M,0,a,b}, \quad \text{for some } \alpha > 1$$

and F(V) touches strictly $F(g_*)$ at 0 from the positive side. A similar statement can be checked also for supersolutions.

Assume by contradiction that such a $V = V_{M,0,a,b}$ exists. Then after a dilation we may assume that $V \in \mathcal{V}_{\delta}$ for some small δ and V is a subsolution in B_1 .

For any $\epsilon > 0$ there exists $\sigma > 0$ such that for all $|t| \leq \sigma$ and all large k's

$$W_t(X) := \alpha V(X + te_n) \le g_k - \epsilon \quad \text{in } B_1,$$

and

$$F(W_{-\sigma}) \subset B_1^+(g_k), \qquad F(W_t) \setminus B_{1/8} \subset B_1^+(g_k).$$

By Lemma 7.2 and Remark 7.3 we obtain that

$$g_k \ge (1 - C\epsilon)W_{-\sigma} \quad \text{in } B_{1/2}$$

and

$$g_k \ge (1 - C\epsilon)W_t$$
 in $B_{3/4} \setminus B_{1/4}$, for all $|t| \le \sigma$.

By choosing ϵ small (depending on α) we see that the functions W_t are strict subsolutions to our free boundary problem, and hence the inequality above can be extended in the interior (see Lemma 2.7) i.e.,

$$g_k \ge (1 - C\epsilon)W_t$$
 in $B_{1/2}$.

Writing this for $t = \sigma$ we see that $\{g_k = 0\}$ stays outside a neighborhood of the origin and we contradict the convergence in the Hausdorff distance to $\{g_* = 0\}$.

We conclude this section by showing that our flatness assumption on the free boundary F(g), implies closeness of g and U.

Lemma 7.9. Assume g solves (2.1). Given any $\delta > 0$ there exist $\bar{\epsilon} > 0$ and $\mu > 0$ depending on δ such that if

(7.4)
$$\{x \in \mathcal{B}_1 : x_n \le -\bar{\epsilon}\} \subset \{x \in \mathcal{B}_1 : g(x,0) = 0\} \subset \{x \in \mathcal{B}_1 : x_n \le \bar{\epsilon}\},\$$

then

$$U(X - \mu \delta e_n) \le g(X) \le U(X + \mu \delta e_n) \quad in \ B_\mu$$

Proof. The proof is by compactness. Assume by contradiction that a sequence of functions g_k satisfies the hypotheses with $\bar{\epsilon}_k \to 0$ but the conclusion does not hold.

Notice that by Harnack inequality $g_k(e_{n+1}/2)$ is bounded be a multiple of $g_k(e_n/2)$ which in view of Lemma 7.4 is bounded by a universal constant. Hence by the second claim in Lemma 7.4 the g_k 's have uniformly bounded $C^{1/2}$ norms on compact subsets of B_1 . After passing to a subsequence we can assume that g_k converges uniformly on compact sets of B_1 to a function g_* with

(7.5)
$$\Delta g_* = 0 \quad \text{in } B_1 \setminus P, \qquad g_* = 0 \quad \text{on } P \cap B_1.$$

By Remark 7.6, g_* is $C^{1/2}$. Moreover, the derivatives of g_* in the x' direction satisfy again (7.5) and we obtain

$$\|D_{x'}^{\beta}g_*\|_{C^{1/2}(B_{1/2})} \le C(\beta).$$

 \square

Now we can separate the variables and write

$$\Delta_{x_n,s}g_* = -\Delta_{x'}g_*$$

and we can argue as in Theorem 6.1 to obtain

$$|g_*(X) - \alpha U(X)| \le C|X|^{3/2}$$

with C universal.

We now want to apply Proposition 7.8 to conclude that g_* solves (2.1) and hence $\alpha = 1$. To do so, we must guarantee that $g_* > 0$ in $B_1 \setminus P$. Otherwise $g_* \equiv 0$ and hence $||g_k||_{L^{\infty}(B_{1/2})} \to 0$. Let $\mathcal{B}_k := \mathcal{B}_{1/8}(x_k)$ be a ball tangent to $F(g_k)$ from the zero side at some point $y_k \in B_{1/8}$. Then, since $||g_k||_{L^{\infty}(B_{1/2})} \to 0$, we have by Lemma 7.1

$$g_k \leq \sigma_k V_{\partial B_k, 0, 0}, \quad \text{with } \sigma_k \to 0.$$

This contradicts the free boundary condition for g_k at y_k .

In conclusion, g_* solves (2.1) and

$$|g_*(X) - U(X)| \le C|X|^{3/2}$$

with C universal.

Rescaling we find

$$g_{k,\mu}(X) - U(X)| \le C\mu$$
 in B_2 , with $g_{k,\mu}(X) := \mu^{-1/2}g_k(\mu X).$

Thus

$$g_{k,\mu}(X) \ge U(X) - C\mu \ge U(X - \bar{\epsilon}_k \mu^{-1} e_n) - C\mu.$$

Now we use that $F(g_{k,\mu}) \subset \{|x_n| \leq \bar{\epsilon}_k \mu^{-1}\}$ and obtain by Lemma 7.2 that in B_1

$$g_{k,\mu} \ge (1 - C\mu)U(X - \bar{\epsilon}_k\mu^{-1}e_n) \ge U(X - (\bar{\epsilon}_k\mu^{-1} + C\mu)e_n)$$

where the last inequality follows once more from boundary Harnack inequality. A similar inequality bounds $g_{k,\mu}$ by above. We choose μ small depending on δ and obtain that $g_{k,\mu}$ satisfies the conclusion of the theorem

$$U(X - \delta e_n) \le \mu^{-1/2} g_k(\mu X) \le U(X + \delta e_n),$$

and we reach a contradiction.

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