# On Monge-Ampère Equations with Homogenous Right-Hand Sides 

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#### Abstract

We study the regularity and behavior at the origin of solutions to the two-dimensional degenerate Monge-Ampère equation $\operatorname{det} D^{2} u=|x|^{\alpha}$ with $\alpha>-2$. We show that when $\alpha>0$, solutions admit only two possible behaviors near the origin, radial and nonradial, which in turn implies $C^{2, \delta}$-regularity. We also show that the radial behavior is unstable. For $\alpha<0$ we prove that solutions admit only the radial behavior near the origin. (c) 2008 Wiley Periodicals, Inc.


## 1 Introduction

We consider the degenerate two-dimensional Monge-Ampère equation

$$
\begin{equation*}
\operatorname{det} D^{2} u=|x|^{\alpha}, \quad x \in B_{1}, \tag{1.1}
\end{equation*}
$$

on the unit disc $B_{1}=\{|x| \leq 1\}$ of $\mathbb{R}^{2}$ and in the range of exponents $\alpha>-2$. Our goal is to investigate the behavior of solutions $u$ near the origin, where the equation becomes degenerate.

The study of (1.1) is motivated by the Weyl problem with nonnegative curvature, posed in 1916 by Weyl [13] himself: Given a Riemannian metric $g$ on the 2 -sphere $\$^{2}$ whose Gauss curvature is everywhere positive, does there exist a global $C^{2}$ isometric embedding $X:\left(\mathbb{S}^{2}, g\right) \rightarrow\left(\mathbb{R}^{3}, d s^{2}\right)$, where $d s^{2}$ is the standard flat metric on $\mathbb{R}^{3}$ ?
H. Lewy [10] solved the problem under the assumption that the metric $g$ is analytic. The solution to the Weyl problem, under the regularity assumption that $g$ has continuous fourth-order derivatives, was given in 1953 by L. Nirenberg [12].
P. Guan and Y.Y. Li [6] considered the question: If the Gauss curvature of the metric $g$ is nonnegative instead of strictly positive and $g$ is smooth, is it still possible to have a smooth isometric embedding? It was shown in [6] that for any $C^{4}$ Riemannian metric $g$ on $\$^{2}$ with nonnegative Gaussian curvature, there is always a $C^{1,1}$ global isometric embedding into $\left(\mathbb{R}^{3}, d s^{2}\right)$.

Examples show that for some analytic metrics with positive Gauss curvature on $\mathbb{S}^{2}$ except at one point, there exists only a $C^{2,1}$ but not a $C^{3}$ global isometric embedding into $\left(\mathbb{R}^{3}, d s^{2}\right)$. Note that the phenomenon is global, since C. S. Lin [11] has shown that for any smooth two-dimensional Riemannian metric with nonnegative Gauss curvature there exists a smooth local isometric embedding into $\left(\mathbb{R}^{3}, d s^{2}\right)$.

This leads to the following question, which was posed in [6]:
Under what conditions on a smooth metric $g$ on $\mathbb{\$}^{2}$ with nonnegative Gauss curvature, is there a $C^{2, \alpha}$ global isometric embedding into $\left(\mathbb{R}^{3}, d s^{2}\right)$, for some $\alpha>0$, or even a $C^{2,1}$ ?
The problem can be reduced to a partial differential equation of Monge-Ampère type that becomes degenerate at the points where the Gauss curvature vanishes. It is well-known that in general one may have solutions to degenerate Monge-Ampère equations that are at most $C^{1,1}$.

One may consider a smooth Riemannian metric $g$ on $\mathbb{S}^{2}$ with nonnegative Gauss curvature that has only one nondegenerate zero. In this case, if we represent the $C^{1,1}$ embedding as a graph, answering the above question amounts to studying the regularity at the origin of the degenerate Monge-Ampère equation

$$
\begin{equation*}
\operatorname{det} D^{2} u=f \quad \text { on } B_{1} \tag{1.2}
\end{equation*}
$$

in the case where the forcing term $f$ vanishes quadratically at $x=0$. More precisely, it suffices to assume that $f(x)=|x|^{2} g(x)$, where $g$ is a positive Lipschitz function. This leads to equation (1.1) when $\alpha=2$.

In addition to the results mentioned above, degenerate equations of the form (1.2) on $\mathbb{R}^{2}$ were previously considered by P . Guan in [5] in the case where $f \in$ $C^{\infty}\left(B_{1}\right)$ and

$$
\begin{equation*}
A^{-1}\left(x_{1}^{2 l}+B x_{2}^{2 m}\right) \leq f\left(x_{1}, x_{2}\right) \leq A\left(x_{1}^{2 l}+B x_{2}^{2 m}\right) \tag{1.3}
\end{equation*}
$$

for some constants $A>0, B \geq 0$, and positive integers $l \leq m$. The $C^{\infty}$ regularity of the solution $u$ of (1.2) was shown in [5], under the additional condition that $u_{x_{2} x_{2}} \geq C_{0}>0$. It was conjectured in [5] that the same result must be true under the weaker condition that $\Delta u \geq C_{0}>0$. This was recently shown by P. Guan and I. Sawyer in [8].

Equation (1.1) has also an interpretation in the language of optimal transportation with quadratic cost $c(x, y)=|x-y|^{2}$. In this setting the problem consists in transporting the density $|x|^{\alpha} d x$ from a domain $\Omega_{x}$ into the uniform density $d y$ in the domain $\Omega_{y}$ in such a way that we minimize the total "transport cost," namely,

$$
\int_{\Omega_{x}}|y(x)-x|^{2}|x|^{\alpha} d x
$$

Then, by a theorem of Y. Brenier [1], the optimal map $x \mapsto y(x)$ is given by the gradient of a solution of the Monge-Ampère equation (1.1). The behavior of these
solutions at the origin gives information on the geometry of the optimal map near the singularity of the measure $|x|^{\alpha} d x$.

We will next state the results of this paper. We assume that $u$ is a solution of equation (1.1). Then $u$ is $C^{\infty}$ smooth away from the origin. The following results describe the regularity of $u$ at the origin. We begin with the case when $\alpha>0$.
THEOREM 1.1 If $\alpha>0$, then $u \in C^{2, \delta}$ for a small $\delta$ depending on $\alpha$.
Theorem 1.1 is a consequence of Theorem 1.2, which shows that there are exactly two types of behavior near the origin.

Theorem 1.2 If $\alpha>0$ and

$$
\begin{equation*}
u(0)=0, \quad \nabla u(0)=0, \tag{1.4}
\end{equation*}
$$

then there exist positive constants $c(\alpha)$ and $C(\alpha)$ depending on $\alpha$ such that either $u$ has the radial behavior

$$
\begin{equation*}
c(\alpha)|x|^{2+\frac{\alpha}{2}} \leq u(x) \leq C(\alpha)|x|^{2+\frac{\alpha}{2}}, \tag{1.5}
\end{equation*}
$$

or, in an appropriate system of coordinates, the nonradial behavior

$$
\begin{equation*}
u(x)=\frac{a}{(\alpha+2)(\alpha+1)}\left|x_{1}\right|^{2+\alpha}+\frac{1}{2 a} x_{2}^{2}+O\left(\left(\left|x_{1}\right|^{2+\alpha}+x_{2}^{2}\right)^{1+\delta}\right) \tag{1.6}
\end{equation*}
$$

for some $a>0$.
The nonradial behavior (1.6) was first shown by P. Guan in [5] under the condition that $u_{x_{2} x_{2}} \geq C_{0}>0$ near the origin, and was recently generalized in [8] to assume only that $\Delta u \geq C_{0}>0$.

The next result states that the radial behavior is unstable.
Theorem 1.3 Suppose $\alpha>0$, and let $u_{0}$ be the radial solution to (1.1),

$$
u_{0}(x)=c_{\alpha}|x|^{2+\frac{\alpha}{2}}
$$

and consider the Dirichlet problem

$$
\operatorname{det} D^{2} u=|x|^{\alpha}, \quad u=u_{0}-\varepsilon \cos (2 \theta) \text { on } \partial B_{1} .
$$

Then $u-u(0)$ has the nonradial behavior (1.6) for small $\varepsilon$.
Subsequences of blowup solutions satisfying (1.5) converge to homogeneous solutions, as shown next.

Theorem 1.4 Under the assumptions of Theorem 1.2, if u satisfies (1.5), then for any sequence of $r_{k} \rightarrow 0$ the blowup solutions

$$
r_{k}^{-2-\frac{\alpha}{2}} u\left(r_{k} x\right)
$$

have a subsequence that converges uniformly on compact sets to a homogeneous solution of (1.1).

In the case $-2<\alpha<0$ solutions have only the radial behavior. Actually, we prove a stronger result by showing that $u$ converges to the radial solution $u_{0}$ in the following sense:

Theorem 1.5 If $-2<\alpha<0$ and (1.4) holds, then

$$
\lim _{x \rightarrow 0} \frac{u(x)}{u_{0}(x)}=1 .
$$

Our results are based on the following argument: assume that a section of $u$, say $\{u<1\}$, is "much longer" in the $x_{1}$-direction compared to the $x_{2}$-direction. If $v$ is an affine rescaling of $u$ so that $\{v<1\}$ is comparable to a ball, then $v$ is an approximate solution of

$$
\operatorname{det} D^{2} v(x) \approx c\left|x_{1}\right|^{\alpha} .
$$

Hence, the geometry of small sections of solutions of this new equation provides information on the behavior of the small sections of $u$. For example, if the sections of $v$ are "much longer" in the $x_{1}$-direction (case $\alpha>0$ ), then the corresponding sections of $u$ degenerate more and more in this direction, producing the nonradial behavior (1.6). If the sections of $v$ are longer in the $x_{2}$-direction (case $\alpha<0$ ), then the sections of $u$ tend to become round and we end up with a radial behavior near the origin.

We close this introduction with the following remarks.
Remark 1.6. From the proofs one can see that the theorems above, with the exception of the instability result, are still valid for the equation with more general right-hand side

$$
\operatorname{det} D^{2} u=|x|^{\alpha} g(x)
$$

with $g \in C^{\delta}\left(B_{1}\right), g>0$.

## Remark 1.7.

(1) We will show in the proof of Theorem 1.1 that solutions of (1.1), with $\alpha>0$, which satisfy the radial behavior (1.5) at the origin are of class $C^{2, \alpha / 2}$.
(2) Theorems 1.1 and 1.2 and the results of Guan in [5] and Guan and Sawyer in [8] imply that solutions of (1.1), with $\alpha$ a positive integer, which satisfy the nonradial behavior (1.6) at the origin are $C^{\infty}$ smooth.

Remark 1.8. Equations of the form

$$
\begin{equation*}
\operatorname{det} D^{2} w=|\nabla w|^{\beta}, \quad \beta=-\alpha, \tag{1.7}
\end{equation*}
$$

for which the set $\{\nabla w=0\}$ is compactly included in the domain of definition, can be reduced to (1.1) by defining $u$ to be the Legendre transform of $w$. Hence Theorem 1.5 establishes the sharp regularity of solutions $w$ of equation (1.7) when $0<\beta<2$.

The paper is organized as follows. In Section 2 we introduce tools and notation to be used later in the paper. In Section 3 we prove Theorem 1.2. In Section 4 we establish the radial behavior of solutions when $-2<\alpha<0$, showing Theorem 1.5. In Section 5 we investigate homogeneous solutions and give the proof of Theorem 1.4. In Section 6 we prove Theorem 1.3. Finally, in Section 7 we show that Theorem 1.2 implies Theorem 1.1.

## 2 Preliminaries

In this section we investigate the geometry of the sections of $u$, namely the sets

$$
S_{t, x_{0}}^{u}:=\left\{u(x)<u\left(x_{0}\right)+\nabla u\left(x_{0}\right) \cdot\left(x-x_{0}\right)+t\right\} .
$$

We omit the indices $u$ and $x_{0}$ whenever there is no possibility of confusion. We recall some facts about such sections.

John's lemma (cf. [9, theorem 1.8.2]) states that any bounded convex set $\Omega \subset$ $\mathbb{R}^{n}$ is balanced with respect to its center of mass. That is, if $\Omega$ has center of mass at the origin, there exists an ellipsoid $E$ (with center of mass 0 ) such that

$$
E \subset \Omega \subset k(n) E
$$

for a constant $k(n)$ depending only on the dimension $n$.
Sections $S_{t, x_{0}}^{u}$ of solutions to Monge-Ampère equations with doubling measure $\mu$ on the right-hand side also satisfy a balanced property with respect to $x_{0}$. We recall the following definition:

Definition 2.1 (Doubling Measure) The measure $\mu$ is doubling with respect to ellipsoids in $\Omega$ if there exists a constant $c>0$ such that for any point $x_{0} \in \Omega$ and any ellipsoid $x_{0}+E \subset \Omega$

$$
\begin{equation*}
\mu\left(x_{0}+E\right) \geq c \mu\left(\left(x_{0}+2 E\right) \cap \Omega\right) \tag{2.1}
\end{equation*}
$$

The following theorem, due to L. Caffarelli [2], holds.
THEOREM 2.2 (Caffarelli) Let $u: \Omega \rightarrow \mathbb{R}$ be a (Alexandrov) solution of

$$
\operatorname{det} D^{2} u=\mu
$$

with $\mu$ a doubling measure. Then for each $S_{t, x_{0}} \subset \Omega$ there exists a unimodular matrix $A_{t}$ such that

$$
\begin{equation*}
k_{0}^{-1} A_{t} B_{r} \subset S_{t, x_{0}}-x_{0} \subset k_{0} A_{t} B_{r} \tag{2.2}
\end{equation*}
$$

with

$$
r=t\left(\mu\left(S_{t, x_{0}}\right)\right)^{-\frac{1}{n}}, \quad \operatorname{det} A_{t}=1
$$

for a constant $k_{0}(c, n)>0$.

The ellipsoid $E=A_{t} B_{r}$ remains invariant if we replace $A_{t}$ with $A_{t} O$ with $O$ orthogonal; thus we may assume that $A$ is triangular. If (2.2) is satisfied we write

$$
S_{t} \sim A_{t}
$$

and say that the eccentricity of $S_{t}$ is proportional to $\left|A_{t}\right|$.
The measure that appears in (1.1), namely

$$
\mu:=|x|^{\alpha} d x
$$

is clearly doubling with respect to ellipsoids for $\alpha>0$. We will see in Section 4 that this property is still true for $-1<\alpha<0$ but fails for $-2<\alpha \leq-1$.

Next we discuss the case when the right-hand side in the Monge-Ampère equation depends only on one variable, i.e.,

$$
\begin{equation*}
\operatorname{det} D^{2} u=h\left(x_{1}\right) \tag{2.3}
\end{equation*}
$$

We will show in Section 3 that such equations are satisfied by blowup limits of solutions to det $D^{2} u=|x|^{\alpha}$ at the origin, when $\alpha>0$. These equations remain invariant under affine transformations. Also, by taking derivatives along the $x_{2}-$ direction, one obtains the Pogorelov-type estimate

$$
u_{22} \leq C
$$

in the interior of the sections of $u$.
Assume that $u$ satisfies equation (2.3) in $B_{1} \subset \mathbb{R}^{n}$, in any dimension $n \geq 2$, and perform the following partial Legendre transformation:

$$
\begin{equation*}
y_{1}=x_{1}, \quad y_{i}=u_{i}(x), \quad i \geq 2, \quad u^{*}(y)=x^{\prime} \cdot \nabla_{x^{\prime}} u-u(x), \tag{2.4}
\end{equation*}
$$

with $x^{\prime}=\left(x_{2}, \ldots, x_{n}\right)$. The function $u^{*}$ is obtained by taking the Legendre transform of $u$ on each slice $x_{1}=$ const. We claim that $u^{*}$ (which is convex in $y^{\prime}$ and concave in $y_{1}$ ) satisfies

$$
\begin{equation*}
u_{11}^{*}+h\left(y_{1}\right) \operatorname{det} D_{y^{\prime}}^{2} u^{*}=0 . \tag{2.5}
\end{equation*}
$$

To see this, we first notice that by the change of variable

$$
v\left(x_{1}, x^{\prime}\right) \rightarrow u\left(x_{1}, x^{\prime}+x_{1} \xi^{\prime}\right)
$$

$v$ satisfies the same equation as $u$ and

$$
v^{*}(y)=u^{*}(y)-y_{1} \xi^{\prime} \cdot y^{\prime} .
$$

Thus we may assume that $D^{2} u$ is diagonal at $x$. Now it is easy to check that

$$
u_{1}^{*}=-u_{1}, \quad \nabla_{y^{\prime}} u^{*}=x^{\prime},
$$

and

$$
u_{11}^{*}=-u_{11}, \quad D_{y^{\prime}}^{2} u^{*}=\left[D_{x^{\prime}}^{2} u\right]^{-1} .
$$

Hence $u^{*}$ satisfies (2.5).
Remark 2.3. The following hold:
(1) The partial Legendre transform of $u^{*}$ is $u$, i.e., $\left(u^{*}\right)^{*}=u$.
(2) The inequality $|u-v| \leq \varepsilon$ implies that $\left|u^{*}-v^{*}\right| \leq \varepsilon$ on their common domain of definition.
(3) In dimension $n=2$, the partial Legendre transform of the function

$$
p\left(x_{1}, x_{2}\right)=a\left|x_{1}\right|^{2+\alpha}+b x_{1} x_{2}+d x_{2}^{2}
$$

is given by

$$
\begin{align*}
p^{*}\left(y_{1}, y_{2}\right) & =\left(a\left|x_{1}\right|^{2+\alpha}+b x_{1} x_{2}+d x_{1}^{2}\right)^{*} \\
& =-a y_{1}^{2+\alpha}+\frac{1}{4 d}\left(y_{2}-b y_{1}\right)^{2} . \tag{2.6}
\end{align*}
$$

Notice that $p$ is a solution of the equation $\operatorname{det} D^{2} u=c\left|x_{1}\right|^{\alpha}$ for an appropriate constant $c$, and $p^{*}$ is a solution of the equation $w_{11}+c\left|y_{1}\right|^{\alpha} w_{22}=0$.

From now on we will restrict our discussion to dimension $n=2$ and the special case where $h\left(x_{1}\right)=\left|x_{1}\right|^{\alpha}$.
Lemma 2.4 Assume that for some $\alpha>0$, $w$ solves the equation

$$
L w:=w_{11}+\left|y_{1}\right|^{\alpha} w_{22}=0 \quad \text { in } B_{1} \subset \mathbb{R}^{2}
$$

with $|w| \leq 1$. Then in $B_{1 / 2}, w$ satisfies

$$
\begin{aligned}
w(y)= & a_{0}+a_{1} \cdot y+a_{2} y_{1} y_{2} \\
& +a_{3}\left(\frac{1}{2} y_{2}^{2}-\frac{1}{(\alpha+2)(\alpha+1)}\left|y_{1}\right|^{2+\alpha}\right)+O\left(\left(y_{2}^{2}+\left|y_{1}\right|^{2+\alpha}\right)^{1+\delta}\right)
\end{aligned}
$$

with $\left|a_{i}\right|$ and $O(\cdot)$ bounded by a universal constant and $\delta=\delta(\alpha)>0$.
Proof: First we prove that $w_{2}$ is bounded in the interior. Since $L w_{2}=0$, the same argument applied inductively would imply that the derivatives of $w$ with respect to $y_{2}$ of any order are bounded in the interior.

To establish the bound on $w_{2}$, we show that

$$
\begin{equation*}
L\left(C w^{2}+\varphi^{2} w_{2}^{2}\right) \geq 0 \tag{2.7}
\end{equation*}
$$

for a smooth cutoff function $\varphi$, to be made precise later. Indeed, a direct computation shows that

$$
L\left(w^{2}\right)=2\left(w_{1}^{2}+\left|y_{1}\right|^{\alpha} w_{2}^{2}\right)
$$

and

$$
\begin{aligned}
L\left(\varphi^{2} w_{2}^{2}\right)= & L\left(\varphi^{2}\right) w_{2}^{2}+\varphi^{2} L\left(w_{2}^{2}\right)+2\left(\varphi^{2}\right)_{1}\left(w_{2}^{2}\right)_{1}+2\left|y_{1}\right|^{\alpha}\left(\varphi^{2}\right)_{2}\left(w_{2}^{2}\right)_{2} \\
= & L\left(\varphi^{2}\right) w_{2}^{2}+2 \varphi^{2}\left(w_{21}^{2}+\left|y_{1}\right|^{\alpha} w_{22}^{2}\right)+8\left(\varphi_{1} w_{2}\right)\left(\varphi w_{21}\right) \\
& +8\left|y_{1}\right|^{\alpha}\left(\varphi_{2} w_{2}\right)\left(\varphi w_{22}\right) ;
\end{aligned}
$$

hence

$$
\begin{aligned}
L\left(C w^{2}+\varphi^{2} w_{2}^{2}\right) \geq & 2 C\left|y_{1}\right|^{\alpha} w_{2}^{2}+2 \varphi^{2}\left(w_{21}^{2}+\left|y_{1}\right|^{\alpha} w_{22}^{2}\right) \\
& +L\left(\varphi^{2}\right) w_{2}^{2}+8\left(\varphi_{1} w_{2}\right)\left(\varphi w_{21}\right)+8\left|y_{1}\right|^{\alpha}\left(\varphi_{2} w_{2}\right)\left(\varphi w_{22}\right)
\end{aligned}
$$

By choosing the cutoff function $\varphi$ such that $\varphi_{1}=0$ for $\left|y_{1}\right| \leq \frac{1}{4}$, then

$$
L\left(\varphi^{2}\right) \geq-C_{1}\left|y_{1}\right|^{\alpha}, \quad\left|\varphi_{1} w_{2}\right| \leq C_{1}\left|y_{1}\right|^{\frac{\alpha}{2}}\left|w_{2}\right|,
$$

and we obtain (2.7) if $C$ is large. Therefore $w_{2}$ is bounded in the interior by the maximum principle.

The equation $w_{11}+\left|y_{1}\right|^{\alpha} w_{22}=0$ and the bound $\left|w_{22}\right| \leq C$ imply the bound

$$
\left|w_{11}\right| \leq C\left|y_{1}\right|^{\alpha} .
$$

Thus $w_{1}$ is bounded. The same estimates as above show that $w_{12}$ and $w_{122}$ are bounded as well. By Taylor's formula, namely,

$$
f(t)=f(0)+f^{\prime}(0) t+\int_{0}^{t}(t-s) f^{\prime \prime}(s) d s,
$$

and the equation $L w=0$, we conclude that

$$
\begin{aligned}
w\left(y_{1}, 0\right) & =w(0)+w_{1}(0) y_{1}-\frac{w_{22}(0)}{(\alpha+2)(\alpha+1)} y_{1}^{2+\alpha}+O\left(\left|y_{1}\right|^{3+\alpha}\right) \\
w\left(y_{1}, y_{2}\right) & =w\left(y_{1}, 0\right)+w_{2}\left(y_{1}, 0\right) y_{2}+\frac{w_{22}(0)}{2} y_{2}^{2}+O\left(\left|y_{2}\right|^{3}+\left|y_{1} y_{2}^{2}\right|\right)
\end{aligned}
$$

and

$$
w_{2}\left(y_{1}, 0\right)=w_{2}(0)+w_{12}(0) y_{1}+O\left(\left|y_{1}\right|^{2+\alpha}\right),
$$

from which the lemma follows.
Notation. By universal constants we understand positive constants that may also depend on the exponent $\alpha$. Also, when there is no possibility of confusion, we use the letters $c$ and $C$ for various universal constants that change from line to line.

## 3 Proof of Theorem 1.2

Throughout this section we assume that $\alpha>0$ and that $u$ satisfies

$$
u(0)=0, \quad \nabla u(0)=0 .
$$

We simply write $S_{t}$ for the section $S_{t, 0}^{u}$.

## Let

$$
\Gamma:=\left\{\left|x_{1}\right|^{2+\alpha}+x_{2}^{2}<1\right\}
$$

be the 1 -section of $\left|x_{1}\right|^{2+\alpha}+x_{2}^{2}$ at 0 . If a set $\Omega$ satisfies

$$
(1-\theta) \Gamma \subset \Omega \subset(1+\theta) \Gamma
$$

we write

$$
\Omega \in \Gamma \pm \theta .
$$

The following approximation lemma constitutes the basic step in the proof of Theorem 1.2.

Lemma 3.1 Assume that $u$ in the section $S_{1}$ satisfies

$$
\begin{equation*}
\operatorname{det} D^{2} u=c f(x), \quad\left|f(x)-\left|x_{1}\right|^{\alpha}\right| \leq \varepsilon \tag{3.1}
\end{equation*}
$$

and

$$
\begin{equation*}
S_{1} \in \Gamma \pm \theta \tag{3.2}
\end{equation*}
$$

with $\varepsilon \leq \varepsilon_{0}$ and $\varepsilon^{1 / 8} \leq \theta, \theta<1$ small. Then, for some small universal $t_{0}$, we have

$$
S_{t_{0}} \in A D_{t_{0}}\left(\Gamma \pm \theta t_{0}^{\delta}\right)
$$

where

$$
A:=\left(\begin{array}{cc}
a_{11} & 0 \\
a_{21} & a_{22}
\end{array}\right), \quad D_{t_{0}}:=\left(\begin{array}{cc}
t_{0}^{1 /(2+\alpha)} & 0 \\
0 & t_{0}^{1 / 2}
\end{array}\right)
$$

and

$$
|A-I| \leq C \theta, \quad C \text { universal. }
$$

Moreover, the constant $c$ in (3.1) satisfies

$$
\begin{equation*}
|c-2(1+\alpha)(2+\alpha)| \leq C \theta \tag{3.3}
\end{equation*}
$$

Proof: We consider the solution

$$
\begin{equation*}
v:=\frac{c^{1 / 2}}{[2(1+\alpha)(2+\alpha)]^{1 / 2}}\left(\left|x_{1}\right|^{2+\alpha}+x_{2}^{2}\right) \tag{3.4}
\end{equation*}
$$

of the equation

$$
\operatorname{det} D^{2} v=c\left|x_{1}\right|^{\alpha}
$$

and compute that

$$
\begin{equation*}
\operatorname{det} D^{2}\left(v+\sqrt{c \varepsilon}|x|^{2}\right)>c\left(\left|x_{1}\right|^{\alpha}+\varepsilon\right) \geq \operatorname{det} D^{2} u \tag{3.5}
\end{equation*}
$$

and

$$
\begin{equation*}
\operatorname{det} D^{2}\left(u+\sqrt{c \varepsilon}|x|^{2}\right)>c(f(x)+\varepsilon) \geq \operatorname{det} D^{2} v \tag{3.6}
\end{equation*}
$$

because $\left|f(x)-\left|x_{1}\right|^{\alpha}\right| \leq \varepsilon$ by assumption.
We first notice that the assumption (3.2) implies that the constant $c$ in equation (3.1) is bounded from above by a universal constant if $\varepsilon_{0}$ is small. This can be easily seen from equation (3.6), which, with the aid of the maximum principle, implies that $u+\sqrt{c \varepsilon}|x|^{2} \geq v$ on $\{u=1\}$ (notice that both $v$ and $w=u+\sqrt{c \varepsilon}|x|^{2}$ satisfy $v(0)=w(0)=0$ and $\nabla w(0)=\nabla v(0)=0)$. Since $\{u=1\} \in \Gamma \pm \theta$, this readily gives a bound on $c$ if we assume that $\theta$ is small.

We will next show that

$$
\begin{equation*}
\{v<1\} \in \Gamma \pm 2 \theta \tag{3.7}
\end{equation*}
$$

which implies the bound (3.3). Indeed, if

$$
\{v<1\} \subset(1-2 \theta) \Gamma
$$

then $v>u+\tilde{c} \theta|x|^{2}$ on $\{u=1\}$ for a universal $\tilde{c}$; thus

$$
v>u+\sqrt{c \varepsilon}|x|^{2} \quad \text { on }\{u=1\}
$$

since, by the assumptions of the lemma, $\sqrt{\varepsilon}<\varepsilon^{1 / 8} \leq \theta$ and $\varepsilon \leq \varepsilon_{0}$ with $\varepsilon_{0}$ sufficiently small. We conclude from the maximum principle (see (3.5)) that $v>$ $u+\sqrt{c \varepsilon}|x|^{2}$ in $S_{1}$. This is a contradiction, since $u(0)=v(0)=0$. If

$$
(1+2 \theta) S_{1} \subset\{v<1\}
$$

then similarly we obtain $v+\sqrt{c \varepsilon}|x|^{2}<u$ in $S_{1}$, a contradiction.
Let $w$ be the solution of the problem

$$
\operatorname{det} D^{2} w=c x_{1}^{2} \text { in } S_{1}, \quad w=u \text { on } \partial S_{1}
$$

By the maximum principle

$$
w+\sqrt{c \varepsilon}\left(|x|^{2}-2\right) \leq u \leq w-\sqrt{c \varepsilon}\left(|x|^{2}-2\right) ;
$$

thus

$$
|w-u| \leq C \sqrt{\varepsilon}
$$

Also from (3.7) we obtain

$$
|w-v| \leq C \theta
$$

Hence, by Remark 2.3, the corresponding partial Legendre transforms defined in Section 2 satisfy in $B_{1 / 2}$

$$
\begin{gather*}
\left|w^{*}-v^{*}\right| \leq C \theta  \tag{3.8}\\
\left|w^{*}-u^{*}\right| \leq C \sqrt{\varepsilon}, \quad u^{*}(0)=0, \quad \nabla u^{*}(0)=0 \tag{3.9}
\end{gather*}
$$

and $w^{*}$ and $v^{*}$ solve the same linear equation

$$
w_{11}^{*}+c\left|y_{1}\right|^{\alpha} w_{22}^{*}=0
$$

Using Lemma 2.4 for the difference $w^{*}-v^{*}$ together with (2.6), (3.4), (3.3), and (3.8) yields to

$$
\begin{align*}
w^{*}= & -\left|y_{1}\right|^{2+\alpha}+\frac{1}{4} y_{2}^{2}+a+b_{1} y_{1}+b_{2} y_{2}  \tag{3.10}\\
& +\theta\left(c y_{1} y_{2}+d_{1}\left|y_{1}\right|^{2+\alpha}+d_{2} y_{2}^{2}+O\left(\left(\left|y_{1}\right|^{2+\alpha}+y_{2}^{2}\right)^{1+\delta}\right)\right)
\end{align*}
$$

with the coefficients $a, b_{i}, c$, and $d_{i}$ bounded by a universal constant.
From (3.9) we find that

$$
w^{*}\left(0, y_{2}\right) \geq-C \sqrt{\varepsilon} \quad \text { and } \quad w^{*}\left(y_{1}, 0\right) \leq C \sqrt{\varepsilon}
$$

since, from the convexity in $y_{2}$ and concavity in $y_{1}$ of $u^{*}$,

$$
u^{*}\left(0, y_{2}\right) \geq 0 \quad \text { and } \quad u^{*}\left(y_{1}, 0\right) \leq 0
$$

This and (3.10) imply the bounds

$$
|a| \leq C \varepsilon^{\frac{1}{2}}, \quad\left|b_{1}\right| \leq C \varepsilon^{\frac{1}{4}}, \quad\left|b_{2}\right| \leq C \varepsilon^{\frac{1}{4}} .
$$

Thus, if $\left|y_{1}\right|^{2+\alpha}+y_{2}^{2} \leq 10 t_{0}$, then

$$
w^{*}=-\left(1-d_{1} \theta\right)\left|y_{1}\right|^{2+\alpha}+\left(\frac{1}{4}+d_{2} \theta\right) y_{2}^{2}+c \theta y_{1} y_{2}+O\left(\varepsilon^{\frac{1}{4}}+\theta t_{0}^{1+\delta}\right)
$$

Hence, by performing the partial Legendre transform on $w^{*}$ (using that $\left(w^{*}\right)^{*}=w$ and (2.6)), we obtain

$$
\begin{equation*}
\left|w-\left[e_{1}\left|x_{1}\right|^{2+\alpha}+e_{2}\left(x_{2}+e_{3} x_{1}\right)^{2}\right]\right| \leq C\left(\varepsilon^{\frac{1}{4}}+\theta t_{0}^{1+\delta}\right) \tag{3.11}
\end{equation*}
$$

for

$$
\left|x_{1}\right|^{2+\alpha}+4 e_{2}^{2}\left(x_{2}+e_{3} x_{1}\right)^{2} \leq 10 t_{0}
$$

with $\left|e_{1}-1\right|,\left|e_{2}-1\right|$, and $\left|e_{3}\right|$ bounded by $C \theta$.
We next observe that if $p(x)=e_{1}\left|x_{1}\right|^{2+\alpha}+e_{2}\left(x_{2}+e_{3} x_{1}\right)^{2}$, then the function

$$
\tilde{p}(y):=\frac{1}{t_{0}} p(F y)
$$

with $F$ given by

$$
F^{-1}:=\left(\begin{array}{cc}
t_{0}^{-1 /(2+\alpha)} & 0 \\
0 & t_{0}^{-1 / 2}
\end{array}\right)\left(\begin{array}{cc}
e_{1}^{1 /(2+\alpha)} & 0 \\
e_{3} e_{2}^{1 / 2} & e_{2}^{1 / 2}
\end{array}\right)=D_{t_{0}}^{-1} A^{-1}
$$

satisfies

$$
\tilde{p}(y)=\left|y_{1}\right|^{2+\alpha}+y_{2}^{2}
$$

Hence, defining

$$
\tilde{w}(y)=\frac{1}{t_{0}} w(F y)
$$

we conclude from (3.11) that

$$
\left|\tilde{w}(y)-\left(\left|y_{1}\right|^{2+\alpha}+y_{2}^{2}\right)\right| \leq C\left(\varepsilon^{\frac{1}{4}} t_{0}^{-1}+\theta t_{0}^{\delta}\right) \quad \text { for }\left|y_{1}\right|^{2+\alpha}+y_{2}^{2} \leq 2
$$

Since $|\tilde{w}-\tilde{u}| \leq C \varepsilon^{1 / 2} t_{0}^{-1}$ (because $|w-u| \leq C \varepsilon^{1 / 2}$ ), we find for $\varepsilon<\min \left(\theta^{8}, \varepsilon_{0}\right)$, with $\varepsilon_{0}$ small, that

$$
\{\tilde{u}<1\} \in \Gamma \pm \gamma
$$

with

$$
\gamma=C\left(\varepsilon^{\frac{1}{4}} t_{0}^{-1}+\theta t_{0}^{\delta}\right) \leq \theta t_{0}^{\delta^{\prime}}
$$

The proof is now complete since $S_{t_{0}}=F\{\tilde{u}<1\}=A D_{t_{0}}\{\tilde{u}<1\}$.
The proof given above also proves the following lemma.
Lemma 3.2 Assume that u satisfies

$$
\operatorname{det} D^{2} u=c f(x) \text { on } S_{1} \quad \text { and } \quad B_{1 / k_{0}} \subset S_{1} \subset B_{k_{0}}
$$

Then, given $\theta_{0}$, there exist $\varepsilon_{1}\left(\theta_{0}, k_{0}\right)$ and $t_{1}\left(\theta_{0}, k_{0}\right)$ small such that if

$$
\left|f(x)-\left|x_{1}\right|^{\alpha}\right| \leq \varepsilon_{1}, \quad \text { then } \quad S_{t_{1}} \in A_{0} D_{t_{1}}\left(\Gamma \pm \theta_{0}\right)
$$

with

$$
A_{0}:=\left(\begin{array}{cc}
a_{0,11} & 0  \tag{3.12}\\
a_{0,21} & a_{0,22}
\end{array}\right)
$$

and

$$
c\left(k_{0}\right) \leq a_{0, i i} \leq C\left(k_{0}\right), \quad\left|a_{0,12}\right| \leq C\left(k_{0}\right),
$$

for some universal constants $c\left(k_{0}\right)$ and $C\left(k_{0}\right)$.

The next proposition shows that if the section $S_{t}$ has large eccentricity for some $t$, then $u$ enjoys the nonradial behavior (1.6) at the origin.

Proposition 3.3 Assume that u solves the equation

$$
\operatorname{det} D^{2} u=|x|^{\alpha} \quad \text { on } S_{1}
$$

and that $S_{1}$ has large eccentricity, i.e.,

$$
F B_{1 / k_{0}} \subset S_{1} \subset F B_{k_{0}}, \quad F:=c\left(\begin{array}{cc}
b & 0 \\
0 & 1 / b
\end{array}\right)
$$

with $b \geq C_{0}$. Then there exists az-system of coordinates such that

$$
\begin{equation*}
u(z)=\frac{a}{(\alpha+2)(\alpha+1)}\left|z_{1}\right|^{2+\alpha}+\frac{1}{2 a} z_{2}^{2}+O\left(\left(\left|z_{1}\right|^{2+\alpha}+z_{2}^{2}\right)^{1+\delta}\right) \tag{3.13}
\end{equation*}
$$

for some $a>0$.
Theorem 1.2 readily follows from Proposition 3.3. Indeed, since $|x|^{\alpha} d x$ is a doubling measure, we know from Theorem 2.2 that there exists $k_{0}$ depending on $\alpha$ such that each section $S_{t} \subset B_{1}$ satisfies

$$
F_{t} B_{1 / k_{0}} \subset S_{t} \subset F_{t} B_{k_{0}}
$$

for some symmetric matrix $F_{t}$. If for any $t>0$ the quotient between the largest and the smallest eigenvalue of $F_{t}$ is greater than $C_{0}^{2}$, then we satisfy the hypothesis of the proposition above for a rescaling of $u$, and therefore obtain the nonradial behavior. In the case when the quotient is bounded from above by $C_{0}^{2}$ for all small $t>0$, then we clearly obtain the radial behavior.

Proof of Proposition 3.3: The proof will be based on an inductive argument, where at each step we use Lemma 3.1.

Define

$$
v_{1}(x):=u(F x),
$$

and compute $v_{1}$ such that it satisfies the equation

$$
\operatorname{det} D^{2} v_{1}(x)=(\operatorname{det} F)^{2}|F x|^{\alpha}=c^{4+\alpha} b^{\alpha}\left|\left(x_{1}, b^{-2} x_{2}\right)\right|^{\alpha} .
$$

Also,

$$
\left\{v_{1}<1\right\}=F^{-1} S_{1} .
$$

If $b$ is large, then $v_{1}$ satisfies the hypothesis of Lemma 3.2. Hence, for some fixed $\theta_{0}$ we obtain

$$
S_{t_{1}}=F\left\{v_{1}<t_{1}\right\} \in F A_{0} D_{t_{1}}\left(\Gamma \pm \theta_{0}\right)
$$

with $A_{0}$ satisfying (3.12).
We assume by induction that for $t=t_{1} t_{0}^{k}$ we have

$$
S_{t} \in F A_{k} D_{t}\left(\Gamma \pm \theta_{0} t_{0}^{(k-1) \delta}\right)
$$

with

$$
A_{k}:=\left(\begin{array}{cc}
a_{k, 11} & 0 \\
a_{k, 21} & a_{k, 22}
\end{array}\right)
$$

and

$$
\begin{equation*}
\frac{c}{2} \leq a_{k, i i} \leq 2 C, \quad\left|a_{k, 21}\right| \leq 2 C . \tag{3.14}
\end{equation*}
$$

We will show that

$$
S_{t_{0} t} \in F A_{k+1} D_{t_{0} t}\left(\Gamma \pm \theta_{0} t_{0}^{k \delta}\right)
$$

where

$$
A_{k+1}=A_{k} E_{k} \quad \text { and } \quad E_{k}:=\left(\begin{array}{cc}
e_{k, 11} & 0 \\
e_{k, 21} & e_{k, 22}
\end{array}\right)
$$

with

$$
\begin{equation*}
\left|e_{k, i i}-1\right| \leq C \theta_{0} t_{0}^{(k-1) \delta}, \quad\left|e_{k, 21}\right| t^{-\frac{\alpha}{2(2+\alpha)}} \leq C \theta_{0} t_{0}^{(k-1) \delta} \tag{3.15}
\end{equation*}
$$

Notice that condition (3.15) implies the bound

$$
\begin{equation*}
\left|A_{k+1}-A_{k}\right| \leq C \theta_{0} t_{0}^{(k-1) \delta} \tag{3.16}
\end{equation*}
$$

To prove this inductive step, we observe that the function

$$
v_{t}(x):=t^{-1} u\left(F A_{k} D_{t} x\right)
$$

satisfies in $\left\{v_{t}<1\right\}$ the equation

$$
\operatorname{det} D^{2} v_{t}=c_{t}|\tilde{x}|^{\alpha}
$$

with

$$
|\tilde{x}|^{\alpha}=\left|\left(a_{k, 11} t^{\frac{1}{2+\alpha}} x_{1}, b^{-2}\left(a_{k, 21} t^{\frac{1}{2+\alpha}} x_{1}+a_{k, 22} t^{\frac{1}{2}} x_{2}\right)\right)\right|^{\alpha}=c_{t}^{\prime} f_{t}(x)
$$

and

$$
\left|f_{t}-\left|x_{1}\right|^{\alpha}\right| \leq b^{-2} t^{\frac{\alpha}{2(2+\alpha)}}
$$

Also,

$$
\left\{v_{t}<1\right\} \in \Gamma \pm \theta_{0} t_{0}^{(k-1) \delta}
$$

since $S_{t} \in F A_{k} D_{t}\left(\Gamma \pm \theta_{0} t_{0}^{(k-1) \delta}\right)$ by the inductive assumption. Hence, if $\delta=$ $\delta(\alpha)$ is chosen small, then $v_{t}$ satisfies the assumptions of Lemma 3.1, leading to

$$
\left\{v_{t}<t_{0}\right\} \in \tilde{A} D_{t_{0}}\left(\Gamma \pm \theta_{0} t_{0}^{k \delta}\right)
$$

with

$$
\begin{equation*}
|\tilde{A}-I| \leq C \theta_{0} t_{0}^{(k-1) \delta} \tag{3.17}
\end{equation*}
$$

Thus

$$
S_{t_{0} t} \in F A_{k} D_{t} \tilde{A} D_{t_{0}}\left(\Gamma \pm \theta_{0} t_{0}^{k \delta}\right)
$$

Defining $E_{k}$ such that $D_{t} \tilde{A}=E_{k} D_{t}$, we see from (3.17) that $E_{k}$ satisfies (3.15). We conclude the proof of the induction step by first choosing $\theta_{0}$ small so that (3.12) and (3.16) imply that (3.14) is always satisfied.

Define

$$
A^{*}:=\lim _{k \rightarrow \infty} A_{k} .
$$

We will next prove that

$$
\begin{equation*}
S_{t} \in F A^{*} D_{t}\left(\Gamma \pm C^{\prime} t^{\delta}\right) \tag{3.18}
\end{equation*}
$$

As before, let $t=t_{1} t_{0}^{k}$. Notice that

$$
A^{*}=A_{k} E_{k}^{*}, \quad E_{k}^{*}:=\Pi_{i=k}^{\infty} E_{i},
$$

and it is straightforward to check from (3.15) that

$$
\begin{equation*}
\left|e_{k, i i}^{*}-1\right| \leq C_{1} t^{\delta}, \quad\left|e_{k, 12}^{*}\right| t^{-\frac{\alpha}{2(2+\alpha)}} \leq C_{1} t^{\delta} . \tag{3.19}
\end{equation*}
$$

We have

$$
A_{k} D_{t}=A^{*}\left(E_{k}^{*}\right)^{-1} D_{t}=A^{*} D_{t} \tilde{E}
$$

with

$$
\left|\tilde{e}_{k, i i}-1\right| \leq C_{2} t^{\delta}, \quad\left|\tilde{e}_{k, 12}\right| \leq C_{2} t^{\delta} .
$$

Now (3.18) follows since

$$
\tilde{E}\left(\Gamma \pm C_{2} t^{\delta}\right) \subset \Gamma \pm C^{\prime} t^{\delta} .
$$

Finally, from (3.18) we see that

$$
u\left(F A^{*} x\right)=\left|x_{1}\right|^{2+\alpha}+x_{2}^{2}+O\left(\left(\left|x_{1}\right|^{2+\alpha}+x_{2}^{2}\right)^{1+\delta}\right)
$$

which implies that in a $z$-system of coordinates

$$
u(z)=\beta_{1}\left|z_{1}\right|^{2+\alpha}+\beta_{2} z_{2}^{2}+O\left(\left(\left|z_{1}\right|^{2+\alpha}+z_{2}^{2}\right)^{1+\delta}\right)
$$

The rescaled functions

$$
r^{-1} u\left(r^{\frac{1}{2+\alpha}} z_{1}, r^{\frac{1}{2}} z_{2}\right)
$$

converge, as $r \rightarrow 0$, to

$$
\tilde{u}(z):=\beta_{1}\left|z_{1}\right|^{2+\alpha}+\beta_{2} z_{2}^{2} .
$$

Moreover, this function solves the limiting equation

$$
\operatorname{det} D^{2} \tilde{u}=\left|z_{1}\right|^{\alpha} .
$$

Hence

$$
2(2+\alpha)(1+\alpha) \beta_{1} \beta_{2}=1,
$$

which implies (3.13).

## 4 Negative Powers

In this section we consider the equation

$$
\begin{equation*}
\operatorname{det} D^{2} u=|x|^{\alpha} \quad \text { in } \Omega \subset \mathbb{R}^{2} \tag{4.1}
\end{equation*}
$$

in the negative range of exponents $-2<\alpha<0$. We will assume, throughout the section, that $0 \in \Omega$ and

$$
u(0)=0, \quad \nabla u(0)=0
$$

Our goal is to prove the following proposition, which shows that solutions of equation (4.1) admit radial behavior only near the origin. This is in contrast with the case $0<\alpha<\infty$, where both the radial behavior and the nonradial behavior (3.13) occur (see Proposition 3.3).

Proposition 4.1 There exist positive constants c and $C$ (depending on $u$ ) such that

$$
c|x|^{2+\frac{\alpha}{2}} \leq u(x) \leq C|x|^{2+\frac{\alpha}{2}}
$$

near the origin.
We distinguish two cases depending on whether the measure $|x|^{\alpha} d x$ is doubling with respect to all ellipsoids (see the discussion in Section 2).

Case 1: $-1<\alpha<0$. In this case the measure

$$
\mu:=|x|^{\alpha} d x
$$

is doubling with respect to ellipsoids. Indeed, it suffices to show that there exists $c>0$ such that for any ellipsoid $E$, we have

$$
\begin{equation*}
\mu\left(x_{0}+E\right) \geq c \mu\left(x_{0}+2 E\right) \tag{4.2}
\end{equation*}
$$

Since $-1<\alpha$, the density

$$
\left(x_{1}^{2}+x_{2}^{2}\right)^{\frac{\alpha}{2}}
$$

is doubling on each line $x_{2}=$ const with the doubling constant independent of $x_{2}$. This implies that the density $\mu=|x|^{\alpha} d x$ is doubling with respect to any line in the plane. From this and the fact that $x_{0}+2 E$ can be covered with translates of $x_{0}+\frac{E}{2}$ over a finite number of directions, we obtain (4.2).

From Theorem 2.2, there exists a matrix $A_{t}$ such that $S_{t} \sim A_{t}$, i.e.,

$$
\begin{equation*}
k_{0}^{-1} A_{t} B_{r} \subset S_{t} \subset k_{0} A_{t} B_{r} \tag{4.3}
\end{equation*}
$$

with

$$
r=t\left(\mu\left(S_{t}\right)\right)^{-\frac{1}{2}}, \quad \operatorname{det} A_{t}=1
$$

In this case Proposition 4.1 follows from the lemma below.
LEMMA 4.2 There exist universal constants $C>0$ large and $\delta>0$ such that if $S_{t} \sim A_{t}$ with $\left|A_{t}\right|>C$, then

$$
\begin{equation*}
S_{\delta t} \sim A_{\delta t} \quad \text { with }\left|A_{\delta t}\right| \leq \frac{\left|A_{t}\right|}{2} \tag{4.4}
\end{equation*}
$$

In particular, $\left|A_{t}\right| \leq C\left|A_{t_{0}}\right|$ if $t \leq t_{0}$.

Proof: We will use a compactness argument. Assume, by contradiction, that the conclusion of the lemma is not true. Then we can find a sequence of solutions $u_{k}$ of (4.1) with sections $S_{t_{k}}^{u_{k}}$ at 0 such that $S_{t_{k}}^{u_{k}} \sim A_{t_{k}}^{u_{k}}$ with $\left|A_{t_{k}}^{u_{k}}\right| \rightarrow \infty$ and (4.4) does not hold for any $\delta>0$.

Without loss of generality, we may assume that

$$
A_{t_{k}}^{u_{k}}:=\left(\begin{array}{cc}
a_{k} & 0  \tag{4.5}\\
0 & a_{k}^{-1}
\end{array}\right), \quad a_{k} \rightarrow \infty .
$$

We renormalize the functions $u_{k}$ as

$$
\begin{equation*}
v_{k}(x):=\frac{1}{t_{k}} u_{k}\left(r_{k} A_{t_{k}}^{u_{k}} x\right) \tag{4.6}
\end{equation*}
$$

so that

$$
\operatorname{det} D^{2} v_{k}=c_{k}\left|A_{t_{k}}^{u_{k}} x\right|^{\alpha}=c_{k}^{\prime}\left|x_{1}^{2}+a_{k}^{-4} x_{2}^{2}\right|^{\frac{\alpha}{2}}
$$

and

$$
k_{0}^{-1} B_{1} \subset S_{1}^{v_{k}} \subset k_{0} B_{1}
$$

Since $S_{t_{k}}^{u_{k}} \sim A_{t_{k}}^{u_{k}}$, i.e., $r_{k}=t_{k}\left(\mu\left(S_{t_{k}}^{u_{k}}\right)\right)^{-1 / 2}$, the Monge-Ampère measure $\operatorname{det} D^{2} v_{k} d x$ satisfies

$$
\operatorname{det} D^{2} v_{k}\left(S_{1}^{v_{k}}\right)=r_{k}^{2} t_{k}^{-2} \mu\left(S_{t_{k}}^{u_{k}}\right)=1
$$

Hence, as $k \rightarrow \infty$ we can find a subsequence of the $v_{k}$ 's that converge uniformly to a function $v$ that satisfies

$$
\begin{equation*}
\operatorname{det} D^{2} v=c\left|x_{1}\right|^{\alpha} d x \tag{4.7}
\end{equation*}
$$

and

$$
k_{0}^{-1} B_{1} \subset S_{1}^{v} \subset k_{0} B_{1}, \quad \operatorname{det} D^{2} v_{k}\left(S_{1}^{v}\right)=1
$$

Obviously the constant $c$ in (4.7) is bounded from above and below by universal constants. Since the right-hand side of (4.7) does not depend on $x_{2}$ and $v$ is constant on $\partial S_{1}^{v}$, Pogorelov's interior estimate holds and we obtain the bound

$$
v_{22}<C_{1} \quad \text { in }\left(2 k_{0}\right)^{-1} B_{1}
$$

This implies that the section $S_{\delta}^{v}$ contains a segment of size $\delta^{1 / 2}$ in the $x_{2}$-direction, namely,

$$
\begin{equation*}
\left\{x_{1}=0,\left|x_{2}\right| \leq\left(\delta / C_{1}\right)^{\frac{1}{2}}\right\} \subset S_{\delta}^{v} \tag{4.8}
\end{equation*}
$$

From Theorem 2.2 there exists

$$
A_{\delta}=\left(\begin{array}{cc}
a & 0  \tag{4.9}\\
b & a^{-1}
\end{array}\right), \quad 0<a<C(\delta), \quad|b| \leq C(\delta)
$$

with

$$
\begin{equation*}
k_{0}^{-1} A_{\delta} B_{r} \subset S_{\delta}^{v} \subset k_{0} A_{\delta} B_{r} \tag{4.10}
\end{equation*}
$$

and

$$
\begin{equation*}
r=\delta\left[\operatorname{det} D^{2} v\left(S_{\delta}^{v}\right)\right]^{-\frac{1}{2}} \tag{4.11}
\end{equation*}
$$

From (4.8) and (4.10) we have

$$
\frac{r}{a} \geq c_{1} \delta^{\frac{1}{2}}
$$

while from (4.9), (4.10), and (4.11) we get

$$
\delta^{2}=r^{2} \operatorname{det} D^{2} v\left(S_{\delta}^{v}\right) \geq c_{2} r^{2} \frac{r}{a}(a r)^{1+\alpha} .
$$

From the last two inequalities we obtain

$$
\begin{equation*}
a \leq C_{2} \delta^{\frac{-\alpha}{4(2+\alpha)}} \leq \frac{1}{4} \quad \text { for } \delta \text { small. } \tag{4.12}
\end{equation*}
$$

Since the $v_{k}$ 's converge uniformly to $v$, their $\delta$-sections also converge uniformly; thus

$$
S_{\delta}^{v_{k}} \sim A_{\delta} \quad \text { for } k \text { large }
$$

and hence

$$
S_{\delta t_{k}}^{u_{k}} \sim A_{t_{k}}^{u_{k}} A_{\delta}
$$

From (4.5), (4.9), and (4.12) we conclude

$$
\left|A_{t_{k}}^{u_{k}} A_{\delta}\right| \leq \frac{\left|A_{t_{k}}^{u_{k}}\right|}{3} \quad \text { for } k \text { large }
$$

which implies that the function $u_{k}$ satisfies (4.4), a contradiction.
Case 2: $-2<\alpha \leq-1$. In this case the measure $\mu$ is not doubling with respect to any convex set but it is still doubling with respect to convex sets that have the origin as the center of mass.

We proceed as in the first case but replace the sections $S_{t}$ with the sections $T_{t}$ that have 0 as the center of mass. The existence of these sections follows from the following lemma, due to L. Caffarelli [3, lemma 2].

Lemma 4.3 (Centered Sections) Let $u: \mathbb{R}^{n} \rightarrow \mathbb{R} \cup\{\infty\}$ be a globally defined convex function (we set $u=\infty$ outside $\Omega$ ). Also, assume $u$ is bounded in a neighborhood of 0 and the graph of $u$ does not contain an entire line. Then for each $t>0$, there exists a " $t$-section" $T_{t}$ centered at 0 ; that is, there exists $p_{t}$ such that the convex set

$$
T_{t}:=\left\{u(x)<u(0)+p_{t} \cdot x+t\right\}
$$

is bounded and has 0 as a center of mass.
Using the lemma above, one can obtain Theorem 2.2 (similarly as in [2]), with $S_{t}$ is replaced by $T_{t}$ : for every $T_{t} \subset \Omega$ as above, there exists a unitary matrix $A_{t}$ such that

$$
\begin{equation*}
k_{0}^{-1} A_{t} B_{r} \subset T_{t} \subset k_{0} A_{t} B_{r} \tag{4.13}
\end{equation*}
$$

with $r=t\left(\mu\left(T_{t}\right)\right)^{-1 / 2}$. If (4.13) is satisfied, we write $T_{t} \sim A_{t}$.

Next we will show the analogue of Lemma 4.2 for this case.
Lemma 4.4 There exist universal constants $C>0$ large and $\delta>0$ such that if $T_{t} \sim A_{t}$ with $\left|A_{t}\right|>C$, then $T_{\delta t} \subset T_{t}$ and

$$
\begin{equation*}
T_{\delta t} \sim A_{\delta t} \quad \text { with }\left|A_{\delta t}\right| \leq \frac{\left|A_{t}\right|}{2} \tag{4.14}
\end{equation*}
$$

Proof: We argue similarly as in the proof of Lemma 4.2. We assume by contradiction that the conclusion does not hold for a sequence of functions $u_{k}$. Proceeding as in the proof of Lemma 4.2, we work with the renormalizations $v_{k}$ of $u_{k}$ defined by (4.6), which satisfy

$$
\operatorname{det} D^{2} v_{k}=c_{k}^{\prime}\left|x_{1}^{2}+a_{k}^{-4} x_{2}^{2}\right|^{\alpha / 2}=: \mu_{k}
$$

and

$$
k_{0}^{-1} B_{1} \subset T_{1}^{v_{k}} \subset k_{0} B_{1}, \quad \operatorname{det} D^{2} v_{k}\left(T_{1}^{v_{k}}\right)=1 .
$$

As $k \rightarrow \infty$, we can find a subsequence of the $v_{k}$ 's that converge uniformly to a function $v$. Since $a_{k} \rightarrow \infty$ and $-2<\alpha \leq-1$, the corresponding measures $\mu_{k}$, when restricted to a line $x_{2}=$ const, converge weakly to the measure $c\left|x_{2}\right|^{1+\alpha} \delta_{\left\{x_{1}=0\right\}}$. This implies that the measures $\mu_{k}$ converge weakly to $c\left|x_{2}\right|^{1+\alpha} d \mathcal{H}_{\left\{x_{1}=0\right\}}^{1}$, where $d \mathcal{H}^{1}$ is the one-dimensional Hausdorff measure. The limit function $v$ therefore satisfies

$$
\begin{gather*}
\operatorname{det} D^{2} v=c\left|x_{2}\right|^{1+\alpha} d \mathcal{H}_{\left\{x_{1}=0\right\}}^{1}  \tag{4.15}\\
k_{0}^{-1} B_{1} \subset T_{1}^{v} \subset k_{0} B_{1}, \quad \operatorname{det} D^{2} v_{k}\left(T_{1}^{v}\right)=1 .
\end{gather*}
$$

Clearly $c$ is bounded from above and below by universal constants.
We notice that the measure $d \mathcal{H}_{\left\{x_{1}=0\right\}}^{1}$ is doubling with respect to any convex set with the center of mass on the line $\left\{x_{1}=0\right\}$. Using the same methods as in the case of the classical Monge-Ampère equation, one can show that the graph of $v$ contains no line segments when restricted to $\left\{x_{1}=0\right\}$ (see Lemma 4.5 below). From this and the fact that $v$ is the convex envelope of its restriction on $\partial T_{1}^{v}$ and $\left\{x_{1}=0\right\}$ (see (4.15)), we conclude that there exist two supporting planes with slopes $\beta e_{2} \pm \gamma e_{1}$ to the graph of $v$ at 0 . Moreover, it follows from the compactness of the equation (4.15) that $\gamma$ can be chosen to be universal, and the sections $T_{\delta}^{v}$ satisfy

$$
T_{\delta}^{v} \subset\left(2 k_{0}\right)^{-1} B_{1}
$$

when $\delta \leq \delta_{0}$, a universal constant. We have

$$
\begin{equation*}
T_{\delta}^{v} \subset\left\{\left|x_{1}\right| \leq c(\gamma) \delta\right\} . \tag{4.16}
\end{equation*}
$$

Let $A_{\delta}$ be of the form (4.9) with

$$
k_{0}^{-1} A_{\delta} B_{r} \subset T_{\delta}^{v} \subset k_{0} A_{\delta} B_{r}
$$

and

$$
\begin{equation*}
r=\delta\left[\operatorname{det} D^{2} v\left(T_{\delta}^{v}\right)\right]^{-\frac{1}{2}} \sim \delta\left(\frac{r}{a}\right)^{1+\frac{\alpha}{2}} \tag{4.17}
\end{equation*}
$$

On the other hand, (4.16) implies

$$
a r \leq c \delta,
$$

which together with (4.17) yields

$$
a \leq c \delta^{\frac{2+\alpha}{\sigma+\alpha}} \leq \frac{1}{4}
$$

for $\delta$ small enough. Now the contradiction follows as in Lemma 4.2.
Lemma 4.5 If $v$ satisfies (4.15), then

$$
v_{0}(t):=v(0, t)
$$

is strictly convex.
Proof: Assume that the conclusion does not hold. Then, after subtracting a linear function, we can assume that

$$
v \geq 0 \quad \text { in } T_{1}^{v}
$$

and

$$
v_{0}(t)=0 \text { for } t \leq 0, \quad v_{0}(t)>0 \text { for } t>0 .
$$

Let

$$
l_{\varepsilon}:=\varepsilon t+a_{\varepsilon}
$$

be such that

$$
\begin{equation*}
0 \in\left\{v_{0}<l_{\varepsilon}\right\}=\left(b_{\varepsilon}, c_{\varepsilon}\right) \rightarrow 0, \quad \frac{c_{\varepsilon}}{\left|b_{\varepsilon}\right|} \rightarrow 0 \quad \text { as } \varepsilon \rightarrow 0 \tag{4.18}
\end{equation*}
$$

We consider the linear function $p_{\varepsilon}$ in $\mathbb{R}^{2}$ such that $\left\{u<p_{\varepsilon}\right\}$ has center of mass on $\left\{x_{1}=0\right\}$ and $p_{\varepsilon}=l_{\varepsilon}$ on $\left\{x_{1}=0\right\}$. We claim that for $\varepsilon$ small, $\left\{u<p_{\varepsilon}\right\}$ is compactly included in $T_{1}^{v}$. Otherwise, the graph of $v$ would contain a segment passing through 0 ; hence $v=0$ is an open set that intersects the line $\left\{x_{1}=0\right\}$ and we contradict (4.15). Since $d \mathcal{H}_{\left\{x_{1}=0\right\}}^{1}$ is doubling with respect to the center of mass of $\left\{u<p_{\varepsilon}\right\}$, we conclude that this set is also balanced around 0 , which contradicts (4.18).

We are now in a position to exhibit the final steps of the proof of Proposition 4.1 in the case $-2<\alpha \leq-1$.

Proof of Proposition 4.1: We choose $t_{0}$ small such that $T_{t_{0}} \subset \Omega$. The existence of $t_{0}$ follows from the fact that the graph of $u$ cannot contain any line segments.

From Lemma 4.4 we conclude that there exists a large constant $K>0$ depending on the eccentricity of $T_{t_{0}}$ such that

$$
T_{t} \sim A_{t} \quad \text { with }\left|A_{t}\right| \leq K \text { for all } t \leq \delta t_{0} .
$$

Claim. There exists $\gamma$ depending on $K$ such that $S_{\gamma t} \subset T_{t}$.
To show this, first observe that by rescaling we can assume that $t=1$. We use the compactness of the problem for fixed $K$. If there exist a sequence $\gamma_{k} \rightarrow 0$ and functions $u_{k}$ for which the conclusion does not hold, then the graph of the limiting function $u_{\infty}$ (of a subsequence of $\left\{u_{k}\right\}$ ) contains a line segment. This is a contradiction since $u_{\infty}$ solves the Monge-Ampère equation (4.1), which proves the claim.

If $t=1$, then from simple geometrical considerations and the claim above we obtain

$$
\gamma k_{0}^{-1} K^{-1} B_{1} \subset S_{\gamma} \subset k_{0} K B_{1}
$$

By rescaling, we find that $S_{t}$ has bounded eccentricity for $t$ small, and the proposition is proved.

## 5 Homogenous Solutions and Blowup Limits

We will consider in this section homogeneous solutions of the equation

$$
\operatorname{det} D^{2} w(x)=|x|^{\alpha} \quad \text { in } \mathbb{R}^{2}
$$

for $\alpha>-2$, namely, solutions of the form

$$
w(x)=r^{2+\frac{\alpha}{2}} g(\theta):=r^{\beta} g, \quad \beta=2+\frac{\alpha}{2} .
$$

In the polar system of coordinates

$$
D^{2} w(x)=r^{\beta-2}\left(\begin{array}{cc}
\beta(\beta-1) g & (\beta-1) g^{\prime} \\
(\beta-1) g^{\prime} & g^{\prime \prime}+\beta g
\end{array}\right) .
$$

Thus, the function $g$ satisfies the following ODE:

$$
\begin{equation*}
\beta g\left(g^{\prime \prime}+\beta g\right)-(\beta-1)\left(g^{\prime}\right)^{2}=\frac{1}{\beta-1} \tag{5.1}
\end{equation*}
$$

We consider $g$ as the new variable in a maximal interval $[a, b]$ where $g$ is increasing, and define $h$ on $[g(a), g(b)]$ as $g^{\prime}=\sqrt{2 h(g)}$. We have $g^{\prime \prime}=h^{\prime}(g)$; thus $h$ satisfies

$$
\beta t\left(h^{\prime}(t)+\beta t\right)-2(\beta-1) h(t)=\frac{1}{\beta-1} .
$$

Solving for $h$ we obtain

$$
\begin{equation*}
2 h_{c}(t)=c t^{2\left(1-\frac{1}{\beta}\right)}-\beta^{2} t^{2}-\frac{1}{(\beta-1)^{2}} \tag{5.2}
\end{equation*}
$$

for some $c$ positive.

The function $g$ on $[a, b]$ is the inverse of

$$
a+\int_{g(a)}^{\xi} \frac{1}{\sqrt{2 h_{c}(t)}} d t
$$

and the length of the interval $[a, b]$ is given by

$$
\begin{equation*}
b-a=\int_{\left\{h_{c}>0\right\}} \frac{1}{\sqrt{2 h_{c}(t)}} d t:=I_{c} . \tag{5.3}
\end{equation*}
$$

Solutions of (5.1) are periodic, of period $2(b-a)$; thus a global solution $g$ on the circle exists if and only if $I_{c}$ equals $\frac{\pi}{k}$ for some integer $k$.

Next we investigate the existence of such solutions. First we notice that for any quadratic polynomial $f(s)=-l^{2} s^{2}+d_{1} s+d_{2}$ of opening $-2 l^{2}$, we have

$$
\begin{equation*}
\int_{\{f>0\}} \frac{1}{\sqrt{f(s)}} d s=\frac{\pi}{l} . \tag{5.4}
\end{equation*}
$$

Therefore if $\phi(s)$ denotes any convex function that intersects the parabola $l^{2} s^{2}$ at two points, and we set $f(s)=-l^{2} s^{2}+d_{1} s+d_{2}$, with $d_{1} s+d_{2}$ denoting the line through the intersection points between $\phi(s)$ and $l^{2} s^{2}$, then

$$
\int_{\left\{\phi(s)-l^{2} s^{2}>0\right\}} \frac{1}{\sqrt{\phi(s)-l^{2} s^{2}}} d s \geq \int_{\{f>0\}} \frac{1}{\sqrt{f(s)}} d s=\frac{\pi}{l} .
$$

If $\phi(s)$ is concave, we obtain the opposite inequality.
Applying the above to $h_{c}(s)$, we find that depending on the convexity of the first term in (5.2), we obtain that the integral $I_{c}$ in (5.3) is less (or greater) than $\frac{\pi}{\beta}$ for $\beta<2$ (or $\beta>2$ ), i.e.,

$$
\begin{equation*}
I_{c}<\frac{\pi}{\beta} \text { if } \beta<2 \text { and } I_{c}>\frac{\pi}{\beta} \text { if } \beta>2 . \tag{5.5}
\end{equation*}
$$

On the other hand, by performing the change of variable $t=s^{\beta / 2}$ in the integral (5.3), we obtain the integral (5.4) with

$$
f(s):=c_{1} s-4 s^{2}-c_{2} s^{2-\beta}
$$

for some positive constants $c_{1}$ and $c_{2}$ depending on $c$. Hence, depending on the convexity of the last term of $f$, the integral $I_{c}$ is greater (or less) than $\frac{\pi}{2}$ for $\beta<2$ (or $\beta>2$ ), i.e.,

$$
\begin{equation*}
I_{c}>\frac{\pi}{2} \text { if } \beta<2 \text { and } I_{c}<\frac{\pi}{2} \text { if } \beta>2 . \tag{5.6}
\end{equation*}
$$

Let $-2<\alpha<0$, or equivalently $1<\beta<2$. It follows from (5.5) and (5.6) that $\frac{\pi}{2}<I_{c}<\frac{\pi}{\beta}$; hence $I_{c}=\frac{\pi}{k}$ for an integral $k$ only when $k=1$. This readily implies that the only homogeneous solution in this case is the radial one.

Assume next that $\alpha>0$. We will now show that in this case, depending on the value of $\beta$, more homogeneous solutions may exist. To this end, denote by $c_{0}=c_{0}(\alpha)$ the value of $c$ for which the two functions

$$
f_{1}(t)=c t^{2\left(1-\frac{1}{\beta}\right)} \quad \text { and } \quad f_{2}(t)=\beta^{2} t^{2}+\frac{1}{(\beta-1)^{2}}
$$

become tangent. When $c<c_{0}$, then the set where $h_{c}(t)>0$ is empty. As $c \rightarrow c_{0}^{+}$ the set $\left\{t: h_{c}(t)>0\right\}$ approaches the point $t_{0}$ at which the two functions $f_{1}(t)$ and $f_{2}(t)$ become tangent when $c=c_{0}$. Since $f_{1}^{\prime}\left(t_{0}\right)=f_{2}^{\prime}\left(t_{0}\right)$ when $c=c_{0}$, the point $t_{0}$ satisfies

$$
2 c\left(1-\frac{1}{\beta}\right) t_{0}^{1-2 / \beta}=2 \beta^{2} t_{0}
$$

which implies that

$$
c\left(1-\frac{1}{\beta}\right) t_{0}^{-2 / \beta}=\beta^{2} .
$$

As $c \rightarrow c_{0}^{+}, f_{1}(t)$ behaves as its Taylor quadratic polynomial, namely,

$$
f_{1}(t) \approx f\left(t_{0}\right)+f^{\prime}\left(t_{0}\right) t^{2}+\frac{f^{\prime \prime}\left(t_{0}\right)}{2} t^{2}
$$

and

$$
\frac{f^{\prime \prime}\left(t_{0}\right)}{2}=c\left(1-\frac{1}{\beta}\right)\left(1-\frac{2}{\beta}\right) t_{0}^{-2 / \beta}=\beta^{2}\left(1-\frac{2}{\beta}\right) .
$$

We conclude that, as $c \rightarrow c_{0}^{+},\left(h_{c}\right)^{+}$behaves as a quadratic polynomial of opening $-4 \beta$, and thus $I_{c}$ converges to $\pi / \sqrt{2 \beta}$. Hence

$$
\left(\frac{\pi}{\sqrt{2 \beta}}, \frac{\pi}{2}\right) \subset\left\{I_{c}: c>c_{0}\right\} \quad \text { and also } \quad\left\{I_{c}: c>c_{0}\right\} \subset\left(\frac{\pi}{\beta}, \frac{\pi}{2}\right)
$$

by (5.5) and (5.6).
Summarizing the discussion above yields the following:
Proposition 5.1 Homogenous solutions to (1.1) are periodic on the unit circle.
(i) If $-2<\alpha<0$, then the only homogeneous solution is the radial one.
(ii) If $\alpha>0$, then there exists a homogeneous solution of principal period $\frac{2 \pi}{k}$ if and only if

$$
\frac{\pi}{k} \in\left\{I_{c}, c>c_{0}(\alpha)\right\} .
$$

In addition,

$$
\left(\frac{\pi}{\sqrt{2 \beta}}, \frac{\pi}{2}\right) \subset\left\{I_{c}, c>c_{0}\right\} \subset\left(\frac{\pi}{\beta}, \frac{\pi}{2}\right) \quad \text { with } \beta=2+\frac{\alpha}{2} .
$$

Using the proposition above, we will now prove Theorem 1.4. We begin with two useful remarks.

Remark 5.2. From (5.2) we see that any point in the positive quadrant can be written as $\left(t, \sqrt{2 h_{c}}\right)$ for a suitable $c$. Hence, given any point $x_{0} \in \partial B_{1}$ and any positive symmetric unimodular matrix $A$, there exists a homogeneous solution $w$ in a neighborhood of $x_{0}$ such that $D^{2} w\left(x_{0}\right)=A$.

Remark 5.3. Equation (5.2) gives

$$
\left[\left(h^{\prime}+\beta t\right)+\beta(\beta-1) t\right] t^{\frac{2}{\beta}-1}=c\left(1-\frac{1}{\beta}\right)
$$

hence $\Delta w\left[r^{2} w_{r r}\right]^{2 / \beta-1}$ is constant for any local homogeneous solution $w$. This quantity will play a crucial role in the proof of Theorem 1.4.
Definition 5.4 For any solution $u$ of equation (1.1), we define

$$
J_{u}(x):=(\Delta u)\left(r^{2} u_{r r}\right)^{\gamma}, \quad \gamma:=\frac{2}{\beta}-1 .
$$

Remark 5.5. The quantity $J_{u}(x)$ remains invariant under the homogeneous scaling

$$
v(x)=r^{-\beta} u(r x), \quad J_{v}(x)=J_{u}(r x) .
$$

We denote by $J_{0}$ the constant obtained when we evaluate $J$ on the radial solution $u_{0}$ of (1.1).

Proposition 5.6 The function $\left|J_{u}-J_{0}\right|$ cannot have an interior maximum in $\Omega \backslash\{0\}$ unless it is constant.

Proof: We compute the linearized operator $u^{i j} M_{i j}$ for

$$
M=\log J_{u}=\log (\Delta u)+\gamma \log \left(x_{i} x_{j} u_{i j}\right)
$$

at a point $x \in \Omega \backslash\{0\}$ where $J_{u}(x) \neq J_{0}$.
By choosing an appropriate system of coordinates and by rescaling, we can assume that $|x|=1$ and $D^{2} u$ is diagonal. By differentiating the equation (1.1) twice, we obtain

$$
\begin{equation*}
u^{i i} u_{k i i}=\alpha x_{k} \tag{5.7}
\end{equation*}
$$

and

$$
u^{i i} u_{k l i i}=u^{i i} u^{j j} u_{k i j} u_{l i j}+\alpha\left(\delta_{k}^{l}-2 x_{k} x_{l}\right) .
$$

Since the linearized equation of each second derivative of $u$ depends on $D^{3} u, D^{2} u$, and $x$, we see that

$$
\begin{equation*}
u^{i j} M_{i j}=H\left(D^{3} u, D^{2} u, x\right) \tag{5.8}
\end{equation*}
$$

where $H$ is a quadratic polynomial in $D^{3} u$ for fixed $D^{2} u>0$ and $x$.
Let $w$ denote the (local) homogeneous solution for which $D^{2} u(x)=D^{2} w(x)$. Since $M_{w}=\log J_{w}$ is constant, we have

$$
H\left(D^{3} w, D^{2} w, \cdot\right)=0
$$

in a neighborhood of $x$.

Claim. We have

$$
\left\|D^{3} u(x)-D^{3} w(x)\right\| \leq C|\nabla M|
$$

with the constant $C$ depending on $D^{2} u$ and $x$.
To prove this claim, we obtain from (5.7) and the equalities

$$
M_{k}=\frac{u_{i i k}}{\Delta u}+\gamma \frac{x_{i} x_{j} u_{i j k}+2 x_{i} u_{i k}}{x_{i} x_{j} u_{i j}}
$$

the following system for the third derivatives of $u$ :

$$
\left(\begin{array}{cccc}
1 & 0 & 1 & 0 \\
0 & 1 & 0 & 1 \\
b_{1} & d_{1} & b_{2} & 0 \\
0 & b_{1} & d_{2} & b_{2}
\end{array}\right)\left(\begin{array}{c}
\frac{u_{111}}{u_{11}} \\
\frac{u_{112}}{u_{11}} \\
\frac{u_{221}}{u_{22}} \\
\frac{u_{22}}{u_{22}}
\end{array}\right)=\left(\begin{array}{c}
\alpha x_{1} \\
\alpha x_{2} \\
M_{1}-2 \gamma \frac{x_{1} u_{11}}{u_{r r}} \\
M_{2}-2 \gamma \frac{x_{2} u_{22}}{u_{r r}}
\end{array}\right)
$$

and

$$
b_{i}=\frac{u_{i i}}{\Delta u}+\gamma \frac{x_{i}^{2} u_{i i}}{u_{r r}}, \quad d_{i}=2 \gamma u_{i i} \frac{x_{1} x_{2}}{u_{r r}} .
$$

The third-order derivatives of $w$ solve the same system but with no dependence on $M$ in the right-hand side vector (since the corresponding $M$ for $w$ is constant).

It is enough to show that the determinant of the third-order derivatives coefficient matrix above is positive. This determinant is equal to

$$
d_{1} d_{2}+\left(b_{1}-b_{2}\right)^{2}=4 \gamma^{2}\left(\frac{x_{1} x_{2}}{u_{r r}}\right)^{2}+\left(b_{1}-b_{2}\right)^{2}
$$

and can vanish only if one of the coordinates, say, $x_{2}=0$, and $b_{1}=b_{2}$, i.e.,

$$
u_{11}^{2}=\frac{1-\gamma}{1+\gamma}=\beta-1
$$

This implies that $J(x)=J_{0}$, which is a contradiction. Thus, the determinant is positive and the claim is proved.

Since $H$ depends quadratically on $D^{3} u$ and $D^{2} u=D^{2} w$ at $x$, the claim above implies that

$$
\begin{aligned}
\left|H\left(D^{3} u, D^{2} u, x\right)\right| & =\left|H\left(D^{3} u, D^{2} u, x\right)-H\left(D^{3} w, D^{2} u, x\right)\right| \\
& \leq C\left(x, D^{2} u\right)\left(|\nabla M|+|\nabla M|^{2}\right) .
\end{aligned}
$$

Hence (5.8) implies that on the set where $J(x) \neq J_{0}$, there exists a smooth function $C(x)$ depending on $u$ such that

$$
\left|u^{i j} M_{i j}\right| \leq C(x)\left(|\nabla M|+|\nabla M|^{2}\right) .
$$

From the strong maximum principle, we conclude that $M$ cannot have a local maximum or minimum in this set unless it is constant. With this the proposition is proved.

Theorem 1.4 will follow from the proposition below.

Proposition 5.7 Suppose that $u$ is a solution $u$ of (1.1), with $\alpha>-2$, which satisfies

$$
\begin{equation*}
c|x|^{\beta} \leq u(x) \leq C|x|^{\beta}, \quad \beta=2+\frac{\alpha}{2} \tag{5.9}
\end{equation*}
$$

Then the limit

$$
J_{u}(0):=\lim _{x \rightarrow 0} J_{u}(x)
$$

exists. Moreover, iffor a sequence of $r_{k} \rightarrow 0$ the blowup solutions

$$
v_{r_{k}}:=r_{k}^{-\beta} u\left(r_{k} x\right)
$$

converge uniformly on compact sets to the solution $w$, then $w$ is homogeneous of degree $\beta$ with $J_{w}=J_{u}(0)$.

Proof: From (5.9) we find that as $x \rightarrow 0, J_{u}(x)$ is bounded away from 0 and $\infty$ by constants depending on $c$ and $C$. We will first show that $\lim _{x \rightarrow 0} J_{u}(x)=$ $J(0)$ exists.

We may assume, without loss of generality, that

$$
\limsup _{x \rightarrow 0} J_{u}(x):=k>J_{0}
$$

Let $x_{i}$ be a sequence of points for which limsup is achieved. The blowup solutions $v_{r_{i}}, r_{i}=\left|x_{i}\right|$, have a subsequence that converges uniformly on compact sets of $\mathbb{R}^{2}$ to a solution $v$. Moreover, there exists a point $y$ on the unit circle for which

$$
J_{v}(y)=k \geq \limsup _{x \rightarrow 0} J_{v}
$$

hence, by Proposition 5.6, $J_{v}$ is constant.
This argument also shows that if

$$
J_{u}(z) \leq k-\varepsilon \quad \text { then } \quad J_{u}(x) \leq k-\delta(\varepsilon) \text { on the circle }|x|=|z|
$$

Thus, if there exists a sequence of points $y_{j} \rightarrow 0$ with

$$
\lim _{y_{j} \rightarrow 0} J_{u}\left(y_{j}\right)<k
$$

then $J_{u}$ would have an interior maximum in the annulus $\left\{x:\left|y_{j}\right| \leq|x| \leq\left|y_{j^{\prime}}\right|\right\}$ that contains one of the points $x_{i}$ given above, a contradiction. This shows that $\lim _{x \rightarrow 0} J_{u}(x)$ exists.

It remains to prove that if $J_{v}$ is constant, then $v$ is homogeneous. It suffices to show that $D^{2} v$ is homogeneous of degree $\beta-2$, or more precisely that for each second derivative $v_{i j}$, we have

$$
\begin{equation*}
x \cdot \nabla v_{i j}=(\beta-2) v_{i j} \tag{5.10}
\end{equation*}
$$

To this end, for a fixed point $x$, we consider the homogeneous solution $w$ with $D^{2} w(x)=D^{2} v(x)$. Since

$$
\nabla J_{v}(x)=\nabla J_{w}(x)=0
$$

the third derivatives of $v$ and $w$ solve the same system. We have seen in the proof of Proposition 5.6 that this system is solvable provided $J_{v} \neq J_{0}$. Thus $D^{3} v(x)=$ $D^{3} w(x)$ if $J_{v} \neq J_{0}$. Since (5.10) is obviously true for $w$, this implies that the equality holds for $u$ as well.

If $J_{v}=J_{0}$, we denote by $\Gamma$ the set where $D^{2} u(x)$ does not coincide with the Hessian of the radial solution. From the proof of Proposition 5.6, we still obtain $D^{3} v(x)=D^{3} w(x)$ if $x \in \Gamma$, and by continuity (5.10) holds for $x \in \bar{\Gamma}$. If $x$ is in the open set $\bar{\Gamma}^{c}$, then $D^{2} v$ coincides with $D^{2} u_{0}$ and (5.10) is again satisfied. This finishes the proof of the proposition.

Proof of Theorem 1.4: The proof of the theorem readily follows from Propositions 4.1, 5.1, and 5.7.

## 6 Proof of Theorem 1.3

We consider the Dirichlet problem

$$
\begin{cases}\operatorname{det} D^{2} u=|x|^{\alpha} & \text { in } B_{1}  \tag{6.1}\\ u=u_{0}-\varepsilon \cos (2 \theta) & \text { on } \partial B_{1}\end{cases}
$$

in the range of exponents $\alpha>0$. Here

$$
u_{0}(x)=c_{\alpha}|x|^{\beta}, \quad \beta=2+\frac{\alpha}{2}
$$

denotes the radial solution of the equation, i.e., det $D^{2} u_{0}=|x|^{\alpha}$. We write the solution as

$$
\begin{equation*}
u=u_{0}-\varepsilon v . \tag{6.2}
\end{equation*}
$$

Heuristically, if $\varepsilon$ is small, $v$ satisfies the linearized equation at $u_{0}$, namely,

$$
\left(D^{2} u_{0}\right)^{-1}: D^{2} v=0
$$

where we use the notation $A: B=\sum_{i j} a_{i j} b_{i j}$ for the Frobenius inner product between two $n \times n$ matrices $A$ and $B$.

At any point $x_{0} \in B_{1}$, we denote by $v$ and $\tau$ the unit normal (radial) and unit tangential direction, respectively, to the circle $|x|=\left|x_{0}\right|$ at $x_{0}$. In ( $\left.\nu, \tau\right)$ coordinates,

$$
D^{2} u_{0}=c_{\alpha} r^{\beta-2}\left(\begin{array}{cc}
\beta(\beta-1) & 0 \\
0 & \beta
\end{array}\right)
$$

hence $v$ satisfies the equation

$$
v_{\nu \nu}+(\beta-1) v_{\tau \tau}=0
$$

Solving this equation with boundary data $v=\cos (2 \theta)$, we obtain the solution

$$
v=r^{\rho} \cos (2 \theta)
$$

with

$$
\rho(\rho-1)+(\beta-1)(\rho-4)=0 .
$$

Solving the quadratic equation with respect to $\rho$ gives

$$
\rho=\frac{2-\beta \pm \sqrt{\beta^{2}+12 \beta-12}}{2}
$$

Since $\beta:=2+\frac{\alpha}{2}>2$, the only acceptable solution is

$$
\rho=\frac{2-\beta+\sqrt{\beta^{2}+12 \beta-12}}{2}
$$

and it satisfies

$$
\begin{equation*}
2<\rho<\beta \tag{6.3}
\end{equation*}
$$

which suggests that close to the origin the perturbation term $\varepsilon v$ dominates $u_{0}$.
We wish to show that the solution $u$ of the Dirichlet problem (6.1) admits at the origin the nonradial behavior (1.6) if $\varepsilon \leq \varepsilon_{0}$ with $\varepsilon_{0}$ sufficiently small. We will argue by contradiction.

Assume that $u$ has the radial behavior

$$
c_{0}|x|^{\beta} \leq u(x) \leq C_{0}|x|^{\beta}
$$

with $c_{0}$ and $C_{0}$ universal constants. By rescaling, we deduce that

$$
c I \leq|x|^{2-\beta} D^{2} u(x) \leq C I
$$

with $I$ denoting the identity matrix.
The function $v$ that is defined by (6.2) satisfies

$$
|v| \leq 1, \quad v=\cos (2 \theta) \text { on } \partial B_{1}
$$

and solves the equation $a^{i j} v_{i j}=0$ with

$$
\begin{aligned}
A=\left(a^{i j}\right) & =\int_{0}^{1}\left(t D^{2} u_{0}+(1-t) D^{2} u\right)^{-1} d t \\
& =\int_{0}^{1}\left(D^{2} u_{0}+\varepsilon(t-1) D^{2} v\right)^{-1} d t
\end{aligned}
$$

Hence

$$
\begin{equation*}
c I \leq r^{\beta-2} A \leq C I \tag{6.4}
\end{equation*}
$$

The solution $u$ has bounded third-order derivatives in $B_{1} \backslash B_{1 / 2}$; thus

$$
\left|D^{2} v(x)\right| \leq C\|v\|_{L^{\infty}} \leq C \quad \text { in } B_{1} \backslash B_{1 / 2}
$$

By rescaling we obtain the bound

$$
\left|D^{2} v(x)\right| \leq C|x|^{-2}
$$

From this we find that

$$
r^{\beta-2}\left|A-D^{2} u_{0}^{-1}\right| \leq C \varepsilon r^{-\beta}
$$

hence $v$ satisfies the Dirichlet problem

$$
\begin{cases}f^{i j} v_{i j}=0 & \text { in } B_{1}  \tag{6.5}\\ v=\cos (2 \theta) & \text { on } \partial B_{1}\end{cases}
$$

with

$$
F:=c r^{\beta-2} A
$$

Hence by (6.4),

$$
c I \leq F \leq C I \quad \text { and } \quad\left|F-F_{0}\right| \leq C \varepsilon r^{-\beta}
$$

with

$$
F_{0}:=v \otimes v+(\beta-1) \tau \otimes \tau
$$

(As before, we denote by $\nu$ and $\tau$ the unit normal (radial) and unit tangential directions to the circle $|x|=\left|x_{0}\right|$ at each point $x_{0} \in B_{1}$.)

Also,

$$
|v| \leq 1 \quad \text { on } B_{1} .
$$

From the definitions of $A$ and $F$ we also obtain

$$
\begin{equation*}
\left\|\nabla\left(F-F_{0}\right)\right\| \leq C\left(r_{0}\right) \varepsilon \quad \text { for }|x| \geq r_{0} . \tag{6.6}
\end{equation*}
$$

Set

$$
w:=r^{\rho} \cos (2 \theta) .
$$

Then $w$ satisfies the equation

$$
F_{0}: D^{2} w=0
$$

thus we have

$$
\left|f^{i j} w_{i j}\right| \leq C r^{\rho-2} \min \left\{\varepsilon r^{-\beta}, 1\right\} .
$$

Applying the Aleksandrov maximum principle on $v-w$ (see [4, theorem 9.1]), we find that

$$
|v-w| \leq C \varepsilon^{\delta}
$$

and therefore (see (6.6))

$$
\begin{equation*}
\left|D^{2} v-D^{2} w\right| \leq C^{\prime}\left(r_{0}\right) \varepsilon^{\delta} \quad \text { for }|x| \geq r_{0} \tag{6.7}
\end{equation*}
$$

We next compute

$$
M_{u}(x):=\log (\Delta u)+\gamma \log \left(r^{2} u_{r r}\right), \quad \gamma:=\frac{2}{\beta}-1,
$$

in terms of $M_{u_{0}}$ for $|x| \geq r_{0}$ with $r_{0}$ small and fixed. We recall that $M_{u_{0}}$ is constant in $x$. Since $u=u_{0}-\varepsilon v$, we find that

$$
\begin{aligned}
M_{u}(x)= & M_{u_{0}}-\varepsilon\left(\frac{\Delta v}{\Delta u_{0}}+\gamma \frac{v_{r r}}{u_{0, r r}}\right) \\
& -\frac{\varepsilon^{2}}{2}\left(\left(\frac{\Delta v}{\Delta u_{0}}\right)^{2}+\gamma\left(\frac{v_{r r}}{u_{0, r r}}\right)^{2}\right)+O\left(\varepsilon^{3}\right) .
\end{aligned}
$$

Because

$$
\operatorname{det} D^{2} u=\operatorname{det} D^{2} u_{0}
$$

the function $v$ satisfies the equation

$$
-u_{0, r r} v_{\tau \tau}-u_{0, \tau \tau} v_{r r}+\varepsilon \operatorname{det} D^{2} v=0
$$

or, equivalently (since $\left.u_{0}(r)=c_{\alpha} r^{\beta}\right)$

$$
v_{r r}+(\beta-1) v_{\tau \tau}=\varepsilon \frac{r^{2-\beta}}{c_{\alpha} \beta} \operatorname{det} D^{2} v
$$

The last equality implies that

$$
\frac{\Delta v}{\Delta u_{0}}+\gamma \frac{v_{r r}}{u_{0, r r}}=\varepsilon \frac{r^{2(2-\beta)}}{c_{\alpha}^{2} \beta^{3}(\beta-1)} \operatorname{det} D^{2} v,
$$

and also that

$$
\begin{aligned}
\left(\frac{\Delta v}{\Delta u_{0}}\right)^{2}+\gamma\left(\frac{v_{r r}}{u_{0, r r}}\right)^{2} & =\left(1+\frac{1}{\gamma}\right)\left(\frac{\Delta v}{\Delta u_{0}}\right)^{2}+O(\varepsilon) \\
& =-\frac{2 r^{2(2-\beta)}}{c_{\alpha}^{2} \beta^{2}(\beta-2)}(\Delta v)^{2}+O(\varepsilon)
\end{aligned}
$$

From (6.7) and the above we conclude that

$$
M_{u}(x)=M_{u_{0}}+\varepsilon^{2} r^{2(2-\beta)}\left[a_{1}\left(-\operatorname{det} D^{2} w\right)+a_{2}(\Delta w)^{2}\right]+O\left(\varepsilon^{2+\delta}\right)
$$

for $|x| \geq r_{0}$, with $O\left(\varepsilon^{2+\delta}\right)$ depending on $r_{0}$. The constants $a_{1}$ and $a_{2}$ are given by

$$
a_{1}=\frac{1}{c_{\alpha}^{2} \beta^{3}(\beta-1)} \quad \text { and } \quad a_{2}=\frac{2}{c_{\alpha}^{2} \beta^{2}(\beta-2)} .
$$

We recall that $w(r, \theta)=r^{\rho} \cos (2 \theta)$. Then, a direct computation shows that each term in the square brackets above is positive. Thus the $\varepsilon^{2}$-term is positive and homogeneous of degree $2(\rho-\beta)$ with $\rho<\beta$ (as shown in (6.3)). We conclude from Proposition 5.6 that

$$
\lim _{x \rightarrow 0} M_{u}(x)>M_{u_{0}}
$$

Hence, from Proposition 5.7, the blowup limit of $u$ at the origin cannot be $u_{0}$. On the other hand, from the symmetry of the boundary data for $u$, we conclude that the function $v-v(0)$ has exactly two disconnected components where it is positive (or negative). Thus the blowup limit at the origin for $u$ has period $\pi$ on the unit circle, which contradicts Proposition 5.1.

## 7 Proof of Theorem 1.1

In this final section we will present the last steps of the proof of Theorem 1.1. We distinguish the two different cases of behavior at the origin, (1.5) and (1.6).

Case 1: Radial Behavior. We will show that solutions of (1.1) with the radial behavior (1.5) are $C^{2, \alpha / 2}$.

We begin by observing that solutions of (1.1) satisfy, in $B_{1} \backslash B_{1 / 2}$, the estimate

$$
\begin{equation*}
\left\|D^{2} u\right\|_{C^{0,1}\left(B_{1} \backslash B_{1 / 2}\right)} \leq C(\alpha) \tag{7.1}
\end{equation*}
$$

provided that

$$
\begin{equation*}
c(\alpha)|x|^{2+\frac{\alpha}{2}} \leq u(x) \leq C(\alpha)|x|^{2+\frac{\alpha}{2}} \tag{7.2}
\end{equation*}
$$

For any $r>0$, the rescaled functions

$$
\begin{equation*}
u^{r}(x):=r^{-2-\frac{\alpha}{2}} u(r x) \tag{7.3}
\end{equation*}
$$

solve the equation (1.1). Since $u$ has the radial behavior (1.5) at the origin, each function $u^{r}$ satisfies (7.2). Hence, applying (7.1) to $u^{r}$, we obtain for $x, y \in$ $B_{1} \backslash B_{1 / 2}$ the estimates

$$
\left|D^{2} u(r x)-D^{2} u(r y)\right| \leq r^{\frac{\alpha}{2}}|x-y|, \quad\left|D^{2} u(r x)\right| \leq C r^{\frac{\alpha}{2}}
$$

The above estimates readily imply that $u \in C^{2, \alpha / 2}$.
Case 2: Nonradial Behavior. In the rest of the section we will show that solutions of (1.1) that satisfy the nonradial behavior (1.6) are also of class $C^{2, \delta}$ for some $\delta>0$. The idea is simple: we approximate $u$ with quadratic polynomials in the $x_{2}$-direction. However, the proof is quite technical.

In order to simplify the constants, we assume that $u$ solves the equation

$$
\begin{equation*}
\operatorname{det} D^{2} u=2(2+\alpha)(1+\alpha)|x|^{\alpha} \tag{7.4}
\end{equation*}
$$

instead of (1.1) and (after rescaling) that

$$
\begin{equation*}
u(x)=\left|x_{1}\right|^{2+\alpha}+x_{2}^{2}+O\left(\left(\left|x_{1}\right|^{2+\alpha}+x_{2}^{2}\right)^{1+\delta}\right) \quad \text { as }|x| \rightarrow 0 \tag{7.5}
\end{equation*}
$$

From now on, we will denote points in $\mathbb{R}^{2}$ with capital letters, $X=\left(x_{1}, x_{2}\right)$.
The Hölder continuity of the second-order derivatives of $u$ follows easily from the next proposition.

Proposition 7.1 Let $\lambda>0$ be small and

$$
Y \in \Omega_{\lambda}:=\left\{\lambda \leq\left|x_{1}\right|^{2+\alpha}+x_{2}^{2} \leq 2 \lambda\right\}
$$

Then there exist universal constants $C$ and $\mu$ such that in $B:=B\left(Y, \lambda^{1+\alpha}\right)$, we have

$$
\left\|D^{2} u\right\|_{C^{\mu}(B)} \leq C \quad \text { and } \quad\left\|D^{2} u-D^{2} u(0)\right\|_{L^{\infty}(B)} \leq \lambda^{\mu}
$$

We will show that in the sections

$$
S_{X_{0}, t}:=\left\{X: u(X)<u\left(X_{0}\right)+\nabla u\left(X_{0}\right) \cdot\left(X-X_{0}\right)+t\right\}
$$

of $u$ at the point

$$
X_{0}=\left(0, x_{0}\right), \quad\left|x_{0}\right| \leq 2 \lambda^{\frac{1}{2}}
$$

we can approximate $u$ by quadratic polynomials of opening 2 on vertical segments. We begin by making the following definition:

DEFINITION 7.2 We say that $u \in Q(e, \varepsilon, \Omega)$ if for any vertical segment $l \subset \Omega$ of length less than $e$, there exists a quadratic polynomial $P_{x_{1}, l}\left(x_{2}\right)$ of opening 2 , namely

$$
P_{x_{1}, l}\left(x_{2}\right)=x_{2}^{2}+p\left(x_{1}, l\right) x_{2}+r\left(x_{1}, l\right)
$$

such that

$$
\left|u\left(x_{1}, x_{2}\right)-P_{x_{1}, l}\left(x_{2}\right)\right| \leq \varepsilon e^{2} \quad \text { on } l
$$

Notice that for $c<1$ we have

$$
Q(e, \varepsilon, \Omega) \subset Q\left(c e, c^{-2} \varepsilon, \Omega\right)
$$

The plan of the proof of Theorem 1.1 is as follows: We prove Proposition 7.1 for points $Y \in S_{X_{0}, t}$ with $t \leq \lambda$. We first show that $u$ belongs to some appropriate $Q$ classes and distinguish two cases one when $t \geq \lambda^{\alpha / 2+1-\delta_{1}}$ for some fixed $\delta_{1}>0$, and the other when $t=\lambda^{\alpha / 2+1-\delta_{1}}$. In the first case we use the same method as in Lemma 3.1 and approximate the right-hand side $|f(X)|^{\alpha / 2}$ of the rescaled Monge-Ampère equation with $\left|x_{1}\right|^{\alpha}$ (see Lemma 7.3). In the second case we approximate $f(X)$ with a more general polynomial $x_{1}^{2}+p x_{1}+q$ and obtain a better approximation ( $Q$ class) for $u$ (Lemma 7.4).

The Hölder estimates for points $Y \in S_{X_{0}, t},\left|x_{0}\right| \geq \lambda^{1 / 2}$, are obtained in appropriate sections $S_{Y, \sigma}$ in which all the values of $|x|$ are comparable. In these sections the Monge-Ampère equation is nondegenerate and the classical estimates apply. To obtain the appropriate section $S_{Y, \sigma}$, we distinguish two cases, depending on the distance from $Y$ to the $x_{2}$-axis. If $\left|y_{1}\right| \geq \lambda^{1 / 2}$, then we take $\sigma$ so that $S_{Y, \sigma}$ is at a distance greater than $\left|y_{1}\right| / 2$ from the $x_{2}$-axis (Lemma 7.5). If $\left|y_{1}\right| \leq \lambda^{1 / 2}$, then we take $\sigma=\lambda^{(2+\alpha) / 2}$ and $S_{Y, \sigma}$ is close enough to the $x_{2}$-axis so that all its points are at a distance comparable to $\lambda^{1 / 2}$ from the origin (Lemma 7.6).

In what follows we will denote by $A_{t}$ and $D_{t}$ the matrices

$$
A_{t}=\left(\begin{array}{cc}
a_{11} & 0 \\
a_{21} & a_{22}
\end{array}\right), \quad D_{t}=\left(\begin{array}{cc}
t^{1 /(2+\alpha)} & 0 \\
0 & t^{1 / 2}
\end{array}\right)
$$

LEMMA 7.3 Let $X_{0}=\left(0, x_{0}\right)$ with $\left|x_{0}\right| \leq 2 \lambda^{1 / 2}, 0<\lambda<1$. Then, for any $\delta_{1}>0$ and

$$
\begin{equation*}
\lambda^{\frac{\alpha}{2}+1-\delta_{1}} \leq t \leq \lambda \tag{7.6}
\end{equation*}
$$

there exists a small $\delta_{2}>0$, depending on $\delta_{1}$, such that

$$
\begin{equation*}
S_{X_{0}, t}-X_{0} \in A_{t} D_{t}\left(\Gamma \pm t^{\delta_{2}}\right) \tag{7.7}
\end{equation*}
$$

with

$$
\begin{equation*}
\left|A_{t}-I\right| \leq t^{\delta_{2}} \tag{7.8}
\end{equation*}
$$

Moreover,

$$
u \in Q\left(t^{\frac{1}{2}}, \lambda^{\delta_{2}}, S_{X_{0}, t}\right)
$$

Proof: We begin by observing that if $t=\lambda$, then the conclusion of the lemma follows from the expansion (7.5) with matrix $A_{t}=I$. We will show by induction, using at each step the approximation Lemma 3.1, that (7.7) and (7.8) hold for every $t=\lambda t_{0}^{k}, k \in \mathbb{N}$, which satisfies (7.6).

Assume that (7.7) and (7.8) hold for some $t=\lambda t_{0}^{k}$ satisfying (7.6), with $A_{t}$ bounded and $a_{t, 11}$ bounded from below. Consider the rescaling

$$
\begin{equation*}
v(X):=\frac{1}{t}\left(u\left(X_{0}+A_{t} D_{t} X\right)-u\left(X_{0}\right)-\nabla u\left(X_{0}\right)\left(A_{t} D_{t} X\right)\right) . \tag{7.9}
\end{equation*}
$$

Since $u$ satisfies (7.4), the function $v$ satisfies the equation

$$
\begin{equation*}
\operatorname{det} D^{2} v=2(2+\alpha)(1+\alpha) a_{11}^{2} a_{22}^{2} t^{-\frac{\alpha}{2+\alpha}}\left|X_{0}+A_{t} D_{t} X\right|^{\alpha} . \tag{7.10}
\end{equation*}
$$

Since

$$
\begin{equation*}
\left|X_{0}+A_{t} D_{t} X\right|^{2}=\left(t^{\frac{1}{2+\alpha}} a_{11} x_{1}\right)^{2}+\left(t^{\frac{1}{2+\alpha}} a_{12} x_{1}+t^{\frac{1}{2}} a_{22} x_{2}+x_{0}\right)^{2} \tag{7.11}
\end{equation*}
$$

and $\left|x_{0}\right| \leq 2 \lambda^{1 / 2}$, we conclude from the above that $v$ satisfies

$$
\begin{equation*}
\operatorname{det} D^{2} v=c|f(X)|^{\frac{\alpha}{2}}, \quad S_{0,1}^{v} \in \Gamma \pm t^{\delta_{2}} \tag{7.12}
\end{equation*}
$$

with

$$
\left|f(X)-x_{1}^{2}\right| \leq C\left(\lambda^{\frac{1}{2}} t^{-\frac{1}{2+\alpha}}+t^{\frac{\alpha}{2(2+\alpha)}}\right) \leq t^{\frac{\delta_{1}}{2(2+\alpha)}} .
$$

Notice that the last inequality holds if (7.6) is satisfied.
Lemma 3.1 with $\varepsilon=t^{\delta^{\prime}}$ and $\delta^{\prime}\left(\delta_{1}, \alpha\right)>0$ small yields

$$
S_{X_{0}, t_{0} t}^{u}-X_{0} \in A_{t_{0} t} D_{t_{0} t}\left(\Gamma \pm\left(t_{0} t\right)^{\delta_{2}}\right)
$$

with

$$
A_{t_{0} t}=A_{t} E_{t}, \quad\left|E_{t}-I\right| \leq C t^{\delta_{2}}
$$

Thus, (7.7) and (7.8) hold for $t^{\prime}=t t_{0}$. If $t^{\prime} \leq \lambda^{\alpha / 2+1-\delta_{1}}$, we stop; otherwise we continue the induction.

From (7.12) we find that

$$
\begin{equation*}
\left|v-\left(\left|x_{1}\right|^{2+\alpha}+x_{2}^{2}\right)\right| \leq C t^{\delta_{2}} \quad \text { in } S_{0,1}^{v} \tag{7.13}
\end{equation*}
$$

which, together with (7.9) and (7.8), yields

$$
u \in Q\left(t^{\frac{1}{2}}, C \lambda^{\delta_{2}}, S_{X_{0}, t}^{u}\right) .
$$

The lemma is proved by replacing $\delta_{2}$ with $\delta_{2} / 2$.

We next examine the borderline case $t=\lambda^{\alpha / 2+1-\delta_{1}}$ and show the better approximation (7.15) of $u$ by quadratic polynomials in the $x_{2}$-variable.

We begin by observing that the conclusion of the previous lemma implies that

$$
S_{X_{0}, t}-X_{0} \in A_{t} D_{t}\left(\Gamma \pm \lambda^{\delta_{2}}\right), \quad\left|A_{t}-I\right| \leq \lambda^{\delta_{2}}
$$

for all $\lambda^{\alpha / 2+1-\delta_{1}} \leq t \leq \lambda$.
LEMMA 7.4 Assume that, for $t=\lambda^{\alpha / 2+1-\delta_{1}}$ and $\delta_{2} \ll \delta_{1}$, we have

$$
\begin{equation*}
S_{X_{0}, t}-X_{0} \in A_{t} D_{t}\left(\Gamma \pm \lambda^{\delta_{2}}\right), \quad\left|A_{t}-I\right| \leq \lambda^{\delta_{2}} \tag{7.14}
\end{equation*}
$$

Then if $\delta_{1}$ is small and universal, we have

$$
\begin{equation*}
u \in Q\left(e, C \lambda^{\delta_{2}}, S_{X_{0}, \frac{t}{2}}\right) \quad \text { for all } e \text { with } \lambda^{\frac{2+\alpha}{4}} \leq e \leq t^{\frac{1}{2}} \tag{7.15}
\end{equation*}
$$

Proof: Let $v$ be the rescaling defined in (7.9). It follows from (7.10), (7.11), and (7.14) that $v$ satisfies

$$
\operatorname{det} D^{2} v=c f(X)^{\frac{\alpha}{2}}, \quad S_{0,1}^{v} \in \Gamma \pm \lambda^{\delta_{2}}
$$

with

$$
\left|f(X)-x_{1}^{2}-p x_{1}-q\right| \leq t^{\frac{\alpha}{2(2+\alpha)}}, \quad|p|,|q| \leq \lambda^{\frac{\delta_{1}}{2+\alpha}}
$$

thus

$$
\left|f(X)^{\frac{\alpha}{2}}-\left(x_{1}^{2}+p x_{1}+q\right)^{\frac{\alpha}{2}}\right| \leq \varepsilon:=t^{\delta_{0}(\alpha)}, \quad \delta_{0}(\alpha)=\frac{\alpha \min \{\alpha, 2\}}{4(2+\alpha)}
$$

Similarly, as in the proof of Lemma3.1, we define the function $w$ as the solution to

$$
\operatorname{det} D^{2} w=c\left(x_{1}^{2}+p x_{1}+q\right)^{\frac{\alpha}{2}}, \quad w=1 \text { on } \partial S_{0,1}^{v}
$$

and obtain (see (7.13)) that

$$
|v-w| \leq C \varepsilon^{\frac{1}{2}}=C t^{\frac{\delta_{0}(\alpha)}{2}} \quad \text { and } \quad\left|w-\left(\left|x_{1}\right|^{2+\alpha}+x_{2}^{2}\right)\right| \leq C \lambda^{\delta_{2}}
$$

By considering the partial Legendre transform $w^{*}$, one can deduce from the last inequality, the bounds on $|p|$ and $|q|$, and Lemma 2.4 that

$$
\left|w_{22}-2\right| \leq C \lambda^{\delta_{2}} \quad \text { in } S_{0,1 / 2}^{v}
$$

This implies that

$$
w \in Q\left(e, C \lambda^{\delta_{2}}, S_{0,1 / 2}^{v}\right) \quad \text { for any } e
$$

hence

$$
v \in Q\left(e, C \lambda^{\delta_{2}}, S_{0,1 / 2}^{v}\right) \quad \text { for } e \geq t^{\frac{\delta_{0}(\alpha)}{8}}
$$

Then, as at the end of the proof of the previous lemma, we obtain that

$$
u \in Q\left(t^{\frac{1}{2}} e, C \lambda^{\delta_{2}}, S_{0, t / 2}^{u}\right) \quad \text { for } e \geq t^{\frac{\delta_{0}(\alpha)}{8}},
$$

from which the lemma follows, since

$$
t^{\frac{1}{2}} t^{\frac{\delta_{0}(\alpha)}{8}} \leq \lambda^{\frac{2+\alpha}{4}}
$$

for $\delta_{1}$ small and universal (depending only on $\alpha$ ).
The next lemma proves Proposition 7.1 for a point $Y \in S_{X_{0}, \lambda}$ at distance greater than $\lambda^{1 / 2}$ from the $x_{2}$-axis, assuming the conclusions of Lemmas 7.3 and 7.4.

Lemma 7.5 Assume that for $\lambda^{\alpha / 2+1-\delta_{1}} \leq t \leq \lambda$, we have

$$
\begin{equation*}
S_{X_{0}, t}-X_{0} \in A_{t} D_{t}\left(\Gamma \pm \lambda^{\delta_{2}}\right), \quad\left|A_{t}-I\right| \leq \lambda^{\delta_{2}} \tag{7.16}
\end{equation*}
$$

and

$$
u \in Q\left(e, C \lambda^{\delta_{2}}, S_{X_{0}, \frac{t}{2}}\right) \text { for some } e, \quad \lambda^{\frac{2+\alpha}{4}} \leq e \leq t^{\frac{1}{2}}
$$

If

$$
Y=\left(y_{1}, y_{2}\right) \in S_{X_{0}, \frac{t}{3}}, \quad 1 \leq\left|y_{1}\right| e^{-\frac{2}{2+\alpha}} \leq 2,
$$

then $D^{2} u$ is Hölder-continuous in the ball $B:=B\left(Y, \lambda^{1+\alpha}\right)$, and for some constant $0<\beta<1$, it satisfies

$$
\begin{equation*}
\left\|D^{2} u\right\|_{C^{0, \beta}(B)} \leq C \quad \text { and } \quad\left|D^{2} u(Y)-D^{2} u(0)\right| \leq C \lambda^{\beta} . \tag{7.17}
\end{equation*}
$$

Proof: Consider the section $S_{Y, c e^{2}}^{u}$ for a small constant $c$. By Theorem 2.2 there exists a matrix

$$
F:=\left(\begin{array}{ll}
a & 0 \\
d & b
\end{array}\right), \quad a, b>0,
$$

such that

$$
\begin{equation*}
F B_{1 / C_{0}} \subset S_{Y, c e^{2}}^{u}-Y \subset F B_{1}, \quad C_{0}(\alpha)>0 \text { universal. } \tag{7.18}
\end{equation*}
$$

Using the assumptions of the lemma and (7.18), we derive bounds on the coefficients of the matrix $F$. Clearly

$$
v:=\frac{c^{1 / 2} e}{b}
$$

satisfies the bound

$$
\begin{equation*}
\frac{1}{2 C_{0}} \leq \frac{c^{1 / 2} e}{b} \leq 2 \tag{7.19}
\end{equation*}
$$

Since $e \leq t^{1 / 2}$, the corresponding section for the rescaling $v$ (see (7.9), (7.13)) satisfies

$$
S_{\tilde{Y}, c e^{2 / t}}^{v} \subset\left\{\left|x_{1}\right|^{2+\alpha}+x_{2}^{2} \leq \frac{3}{4}\right\},
$$

or, more precisely,

$$
S_{\tilde{Y}, c e^{2} / t}^{v}-\tilde{Y} \subset\left(\left(\frac{e^{2}}{t}\right)^{\frac{1}{2+\alpha}}+\lambda^{\frac{\delta_{2}}{2(2+\alpha)}}\right) B_{1}
$$

Thus,

$$
D_{t}^{-1} A_{t}^{-1} F B_{1 / C_{0}} \subset\left(e^{\frac{2}{2+\alpha}} t^{-\frac{1}{2+\alpha}}+\lambda^{\delta_{3}}\right) B_{1} .
$$

The last inclusion implies the estimate

$$
\begin{equation*}
|d| \leq 2 C_{0}\left(e^{\frac{2}{2+\alpha}} t^{\frac{\alpha}{2(2+\alpha)}}+\lambda^{\delta_{3}} t^{\frac{1}{2}}+a \lambda^{\delta_{3}}\right) \leq 4 C_{0}\left(e^{\frac{2}{2+\alpha}}+a\right) \lambda^{\delta_{3}} \tag{7.20}
\end{equation*}
$$

The rescaling

$$
w(x):=\frac{1}{b^{2}} u(Y+F x)
$$

satisfies

$$
\begin{equation*}
\operatorname{det} D^{2} w=\frac{a^{2}}{b^{2}} f(x)^{\frac{\alpha}{2}}, \quad B_{1 / C_{0}} \subset S_{0, \nu^{2}}^{w} \subset B_{1}, \tag{7.21}
\end{equation*}
$$

with

$$
\begin{equation*}
f(x):=\left(y_{1}+a x_{1}\right)^{2}+\left(y_{2}+d x_{1}+b x_{2}\right)^{2} \tag{7.22}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|w-P_{x_{1}}^{\prime}\left(x_{2}\right)\right| \leq \lambda^{\delta_{3}} \quad \text { in } S_{0, \nu^{2}}^{w} . \tag{7.23}
\end{equation*}
$$

We claim that if $c$ is chosen small and universal, then

$$
\begin{equation*}
2 a \leq e^{\frac{2}{2+\alpha}} \leq\left|y_{1}\right| . \tag{7.24}
\end{equation*}
$$

Indeed, otherwise from (7.21), we would deduce that

$$
\operatorname{det} D^{2} w \geq a^{2+\alpha} b^{-2}\left(x_{1}+\frac{y_{1}}{a}\right)^{\alpha}
$$

with

$$
(2 a)^{2+\alpha} b^{-2} \geq e^{2} b^{-2} \geq c^{-1} v^{2}
$$

and for small $c$ we would contradict $B_{1 / C_{0}} \subset S_{0, v^{2}}^{w}$, since $v$ is bounded.

From (7.19), (7.20), (7.24), and $\left|y_{2}\right| \leq 4 \lambda^{1 / 2}$ we obtain that $f(x) / y_{1}^{2}$ is bounded away from 0 and $\infty$ by universal constants, and also its derivatives are bounded by universal constants. From (7.21) we find that

$$
c_{1} \leq \frac{a^{2}\left|y_{1}\right|^{\alpha}}{b^{2}} \leq C_{1},
$$

which implies that $a^{2+\alpha},\left|y_{1}\right|^{2+\alpha}, b^{2}$, and $e^{2}$ are all comparable. Moreover, using also (7.23), we have

$$
\begin{equation*}
\left\|D^{2} w\right\|_{C^{0,1}} \leq C, \quad\left|w_{22}-2\right| \leq \lambda^{\delta_{4}} \text { in } \frac{S_{0, v^{2}}}{2} . \tag{7.25}
\end{equation*}
$$

Hence

$$
\begin{equation*}
\left|w_{22}(x)-w_{22}(y)\right| \leq C \lambda^{\frac{\delta_{4}}{2}}|x-y|^{\frac{1}{2}} \quad \text { for } x, y \in \frac{S_{0, \nu^{2}}}{2} \tag{7.26}
\end{equation*}
$$

Also, we have

$$
D^{2} u(Y+F x)=b^{2}\left(F^{-1}\right)^{T} D^{2} w(x) F^{-1}
$$

with

$$
b F^{-1}=\left(\begin{array}{cc}
\frac{b}{a} & 0 \\
-\frac{d}{a} & 1
\end{array}\right)=\left(\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right)+O\left(\lambda^{\delta_{3}}\right),
$$

which together with (7.25) implies the second part of the conclusion (7.17).
Finally, since

$$
|F x| \geq \frac{b|x|}{2} \geq \lambda^{1+\alpha}|x|,
$$

we obtain from (7.25) and (7.26) the estimate

$$
\left|D^{2} u(Y+F x)-D^{2} u(Y+F y)\right| \leq C \lambda^{\frac{\delta_{4}}{2}}|x-y|^{\frac{1}{2}} \leq C|F x-F y|^{\beta} .
$$

This finishes the proof of the lemma.
The next lemma proves Hölder continuity when $Y$ is $\lambda^{1 / 2}$ close to the $x_{2}$-axis.
Lemma 7.6 Assume that (7.16) holds for $t=\lambda^{\alpha / 2+1-\delta_{1}}$,

$$
u \in Q\left(e, \lambda^{\delta_{2}}, S_{X_{0}, \frac{t}{2}}\right) \quad \text { for } e=\lambda^{\frac{2+\alpha}{4}}
$$

and

$$
\left|x_{0}\right| \geq \frac{\lambda^{1 / 2}}{2}, \quad Y \in S_{X_{0}, \frac{t}{3}}, \quad\left|y_{1}\right| \leq e^{\frac{2}{2+\alpha}}
$$

Then the conclusion of Lemma 7.5 still holds.
Proof: The proof is very similar to that of Lemma 7.5. The only difference is that now the second term of $f$ in (7.22) dominates the sum.

Indeed, since $\lambda^{1 / 2} \geq\left|y_{1}\right|$ and $\left|y_{2}\right| \geq \lambda^{1 / 2} / 4$, the function $f(x) / y_{2}^{2}$ is bounded away from 0 and $\infty$ by universal constants, and also its derivatives are bounded by universal constants. Hence $a^{2+\alpha}, y_{2}^{2+\alpha}, b^{2}$, and $e^{2}$ are all comparable and the rest of the proof is the same.

Proof of Proposition 7.1: For $Y \in \Omega_{\lambda}$ we consider the section $S_{Y, \sigma}^{u}$ that becomes tangent to the $x_{2}$-axis at $X_{0}=\left(0, x_{0}\right)$. Since $|x|^{\alpha} d x$ is doubling, there exists a universal constant $C_{1}$ such that

$$
Y \in S_{X_{0}, t / 3}^{u}, \quad\left|x_{0}\right| \leq 2 \lambda^{\frac{1}{2}}, \quad t:=C_{1} \sigma \leq C_{2} \lambda .
$$

We distinguish the following three cases:
(1) If $t \geq t_{0}:=\lambda^{\alpha / 2+1-\delta_{1}}$, then the proposition follows from Lemmas 7.3 and 7.5 with

$$
e=\left|y_{1}\right|^{\frac{2+\alpha}{2}} \geq c_{1} t^{\frac{1}{2}}
$$

(2) If $t \leq t_{0}$ and $\left|y_{1}\right| \geq \lambda^{1 / 2}$, then we apply Lemmas 7.4 and 7.5 for $S_{X_{0}, t_{0}}$ with $e$ defined as above.
(3) If $t \leq t_{0}$ and $\left|y_{1}\right| \leq \lambda^{1 / 2}$, then we apply Lemma 7.4 and Lemma 7.6. We remark that the hypothesis $\left|x_{0}\right| \geq \lambda^{1 / 2} / 2$ is satisfied because $Y \in \Omega_{\lambda}$.

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